#### Thermodynamics: An Engineering Approach 8th Edition Yunus A. Çengel, Michael A. Boles McGraw-Hill, 2015

### CHAPTER 8 EXERGY

Lecture slides by Mehmet Kanoglu

Copyright © The McGraw-Hill Education. Permission required for reproduction or display.

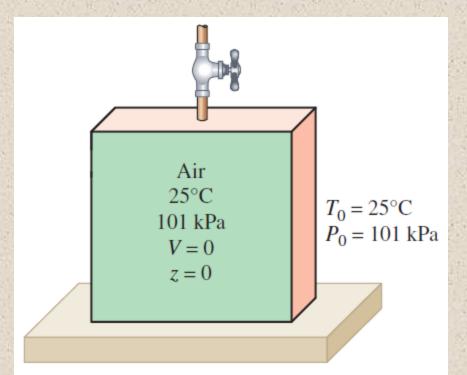
# **Objectives**

- Examine the performance of engineering devices in light of the second law of thermodynamics.
- Define *exergy*, which is the maximum useful work that could be obtained from the system at a given state in a specified environment.
- Define reversible work, which is the maximum useful work that can be obtained as a system undergoes a process between two specified states.
- Define the exergy destruction, which is the wasted work potential during a process as a result of irreversibilities.
- Define the second-law efficiency.
- Develop the exergy balance relation.
- Apply exergy balance to closed systems and control volumes.

# **EXERGY: WORK POTENTIAL OF ENERGY**

The useful work potential of a given amount of energy at some specified state is called *exergy*, which is also called the *availability* or *available energy*.

A system is said to be in the dead state when it is in thermodynamic equilibrium with the environment it is in.



#### FIGURE 8-1

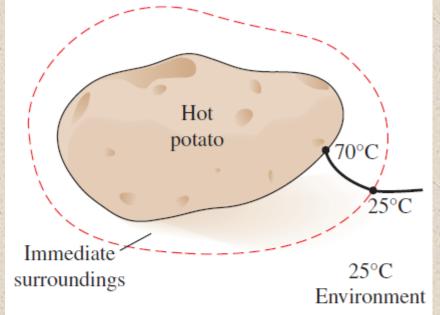
A system that is in equilibrium with its environment is said to be at the dead state. A system delivers the maximum possible work as it undergoes a reversible process from the specified initial state to the state of its environment, that is, the dead state. This represents the useful work potential of the system at the specified state and is called exergy.

Exergy represents the upper limit on the amount of work a device can deliver without violating any thermodynamic laws.



### FIGURE 8–3

The atmosphere contains a tremendous amount of energy, but no exergy.



### FIGURE 8–2

The immediate surroundings of a hot potato are simply the temperature gradient zone of the air next to the potato.

### **Exergy (Work Potential) Associated** with Kinetic and Potential Energy

(kJ/kg)

Exergy of kinetic energy:

 $x_{\rm pe} = {\rm pe} = gz$  (kJ/kg)

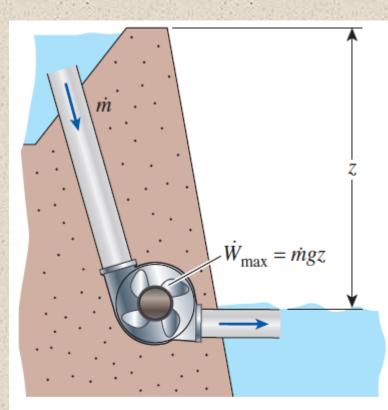
Exergy of potential energy:

 $x_{\rm ke} = {\rm ke} = \frac{V^2}{2}$ Unavailable energy Total energy Exergy

#### FIGURE 8–6

Unavailable energy is the portion of energy that cannot be converted to work by even a reversible heat engine.

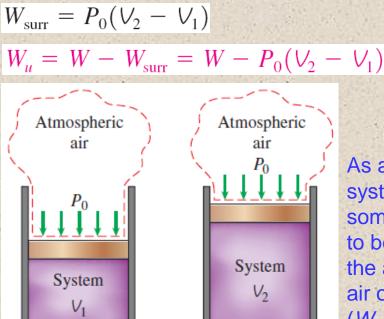
The exergies of kinetic and potential energies are equal to themselves, and they are entirely available for work.



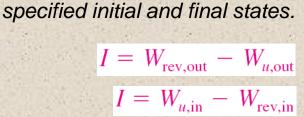
### FIGURE 8–4

The *work potential* or *exergy* of potential energy is equal to the potential energy itself.

## **REVERSIBLE WORK AND IRREVERSIBILITY**



As a closed system expands, some work needs to be done to push the atmospheric air out of the way  $(W_{surr})$ .

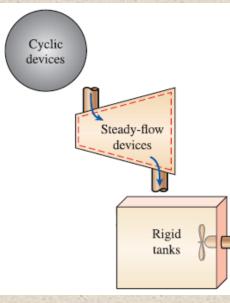


**Reversible work W**<sub>rev</sub>: The maximum amount of

minimum work that needs to be supplied) as a

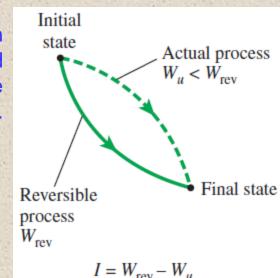
useful work that can be produced (or the

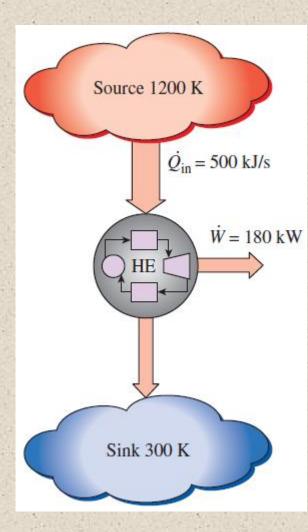
system undergoes a process between the



The difference between reversible work and actual useful work is the irreversibility.

For constant-volume systems, the total actual and useful works are identical  $(W_u = W)$ .





# The Rate of Irreversibility of a Heat Engine

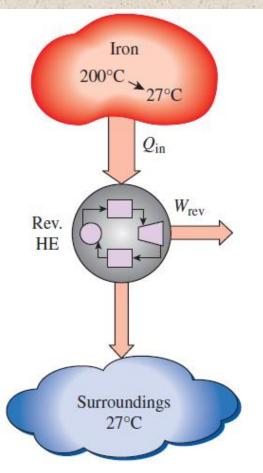
$$\dot{W}_{\text{rev,out}} = \eta_{\text{th,rev}} \dot{Q}_{\text{in}} = \left(1 - \frac{T_{\text{sink}}}{T_{\text{source}}}\right) \dot{Q}_{\text{in}}$$
$$= \left(1 - \frac{300 \text{ K}}{1200 \text{ K}}\right) (500 \text{ kW}) = 375 \text{ kW}$$

80 G S

18 G & X

$$\dot{I} = \dot{W}_{\text{rev,out}} - \dot{W}_{u,\text{out}} = 375 - 180 = 195 \text{ kW}$$

# Irreversibility during the Cooling of an Iron Block



#### FIGURE 8–12

An irreversible heat transfer process can be made reversible by the use of a reversible heat engine.

W<sub>rev</sub>

$$\delta W_{\text{rev}} = \eta_{\text{th,rev}} \, \delta Q_{\text{in}} = \left(1 - \frac{T_{\text{sink}}}{T_{\text{source}}}\right) \delta Q_{\text{in}} = \left(1 - \frac{T_0}{T}\right) \delta Q_{\text{in}}$$

$$W_{\text{rev}} = \int \left(1 - \frac{T_0}{T}\right) \delta Q_{\text{in}}$$

$$\delta E_{\text{in}} - \delta E_{\text{out}} = dE_{\text{system}}$$

$$\delta E_{\text{in}} - \delta E_{\text{out}} = dE_{\text{system}}$$

$$C_{\text{hange in internal, kinetic, potential, etc., energies}$$

$$-\delta Q_{\text{out}} = dU = mc_{\text{avg}}dT$$

$$\delta Q_{\text{in,heat engine}} = \delta Q_{\text{out,system}} = -mc_{\text{avg}}dT$$

$$= \int_{T_1}^{T_0} \left(1 - \frac{T_0}{T}\right) (-mc_{\text{avg}}dT) = mc_{\text{avg}}(T_1 - T_0) - mc_{\text{avg}}T_0 \ln \frac{T_1}{T_0}$$

$$= (500 \text{ kg})(0.45 \text{ kJ/kg·K}) \left[ (473 - 300) \text{ K} - (300 \text{ K}) \ln \frac{473 \text{ K}}{300 \text{ K}} \right]$$

$$= 8191 \text{ kJ}$$
Surrounding air

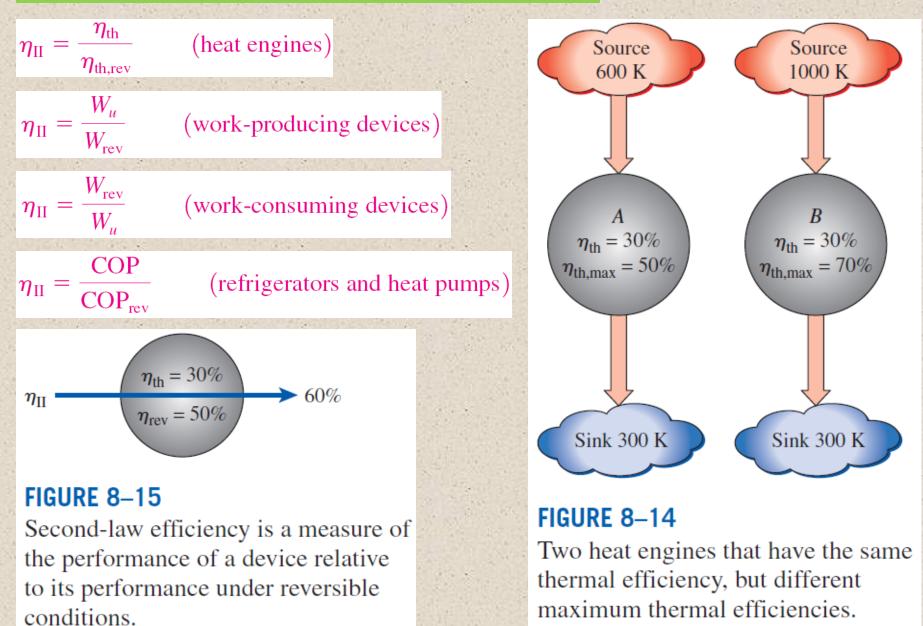
 $T_0 = 27^{\circ} \text{C}$ 

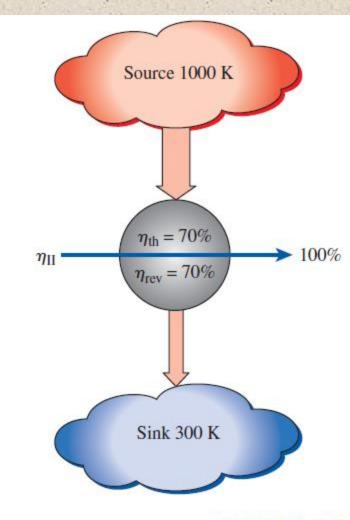
Iron

200°C ★ 27°C



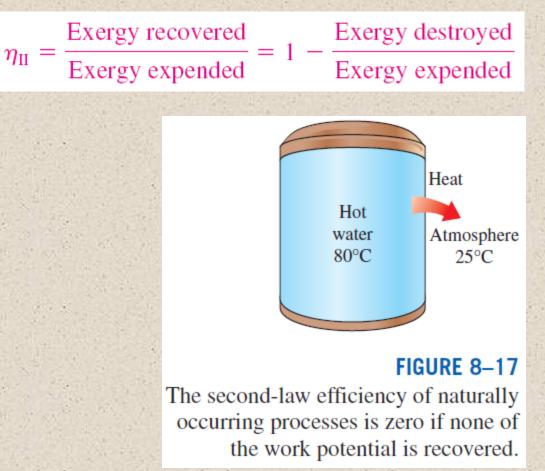
# **SECOND-LAW EFFICIENCY**



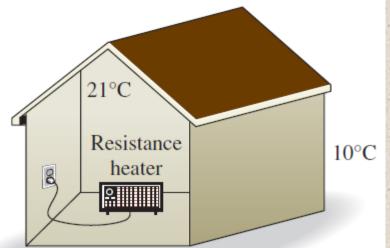


#### FIGURE 8–16 Second-law efficiency of all reversible devices is 100 percent.

General definition of exergy efficiency:



### **Second-Law Efficiency of Resistance Heaters**



$$\eta_{\text{II,electric heater}} = \frac{\dot{X}_{\text{recovered}}}{\dot{X}_{\text{expended}}} = \frac{\dot{X}_{\text{heat}}}{\dot{W}_{e}}$$
$$\dot{Q}_{e} = \dot{W}_{e} = \frac{\dot{Q}_{e}(1 - T_{0}/T_{H})}{\dot{W}_{e}} = 1 - \frac{T_{0}}{T_{H}}$$

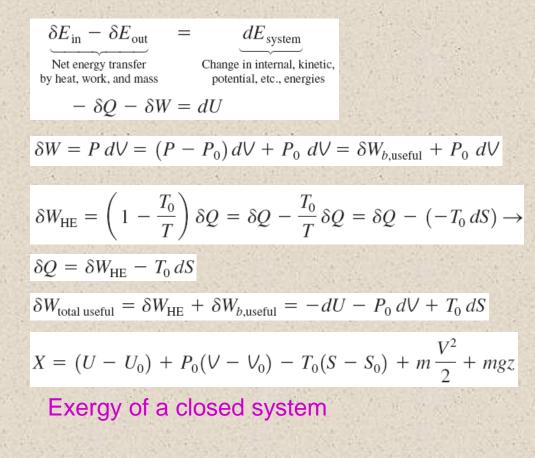
$$\text{COP}_{\text{HP,rev}} = \frac{1}{1 - T_L/T_H} = \frac{1}{1 - (10 + 273 \text{ K})/(21 + 273 \text{ K})} = 26.7$$

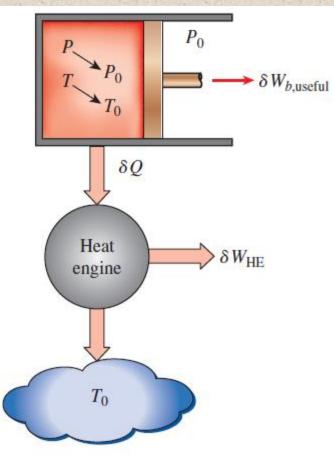
$$\eta_{\text{II}} = \frac{\text{COP}}{\text{COP}_{\text{rev}}} = \frac{1.0}{26.7} = 0.037 \text{ or } 3.7\%$$

$$\eta_{\text{II,electric heater}} = 1 - \frac{T_0}{T_H} = 1 - \frac{(10 + 273)\text{K}}{(21 + 273)\text{K}} = 0.037 \text{ or } 3.7\%$$

# **EXERGY CHANGE OF A SYSTEM**

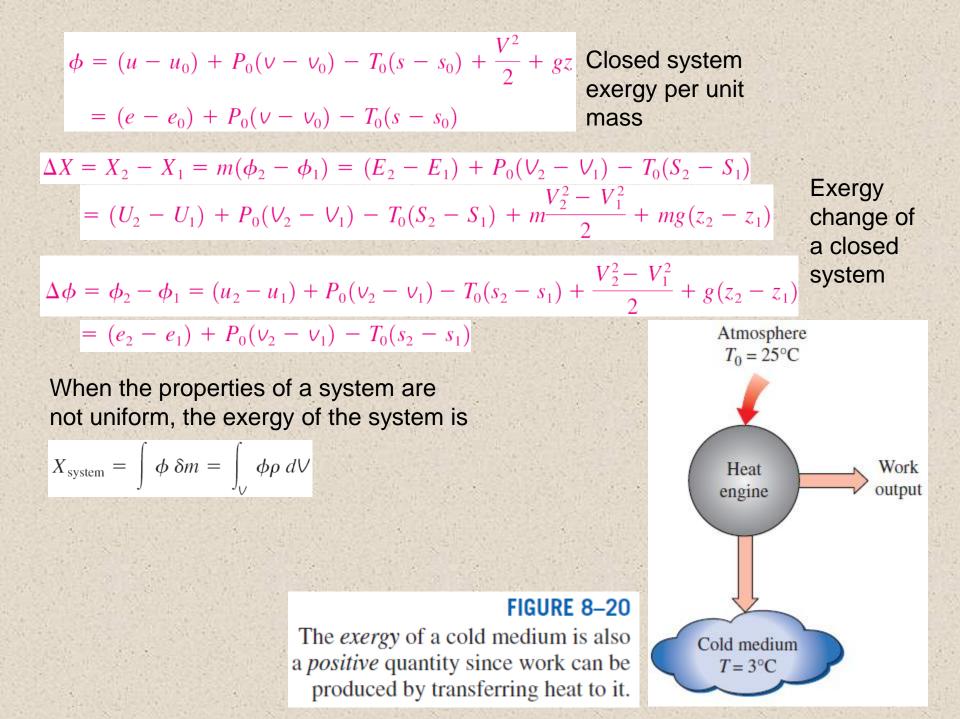
### Exergy of a Fixed Mass: Nonflow (or Closed System) Exergy





#### FIGURE 8–19

The *exergy* of a specified mass at a specified state is the useful work that can be produced as the mass undergoes a reversible process to the state of the environment.



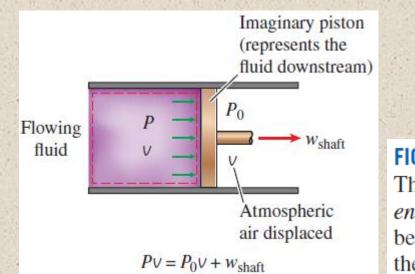
### Exergy of a Flow Stream: Flow (or Stream) Exergy

$$\begin{aligned} x_{\text{flowing fluid}} &= x_{\text{nonflowing fluid}} + x_{\text{flow}} \\ &= (u - u_0) + P_0(v - v_0) - T_0(s - s_0) + \frac{V^2}{2} + gz + (P - P_0)v \\ &= (u + Pv) - (u_0 + P_0v_0) - T_0(s - s_0) + \frac{V^2}{2} + gz \\ &= (h - h_0) - T_0(s - s_0) + \frac{V^2}{2} + gz \end{aligned}$$

$$\begin{aligned} & \text{Exergy of flow energy} \\ x_{\text{flow}} &= Pv - P_0v = (P - P_0)v \\ &\text{Kow} = Pv - P_0v = (P - P_0)v \end{aligned}$$

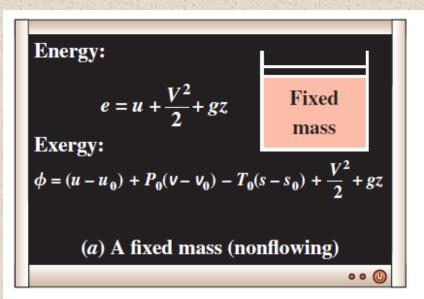
$$\begin{aligned} & \text{Flow} \\ &\text{exergy} \end{aligned}$$

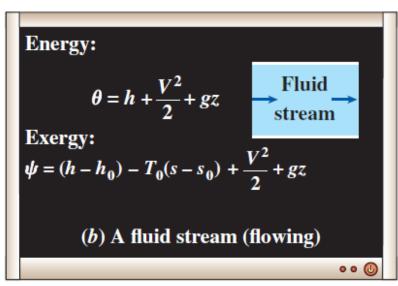
$$\begin{aligned} &\psi = (h - h_0) - T_0(s - s_0) + \frac{V^2}{2} + gz \\ &\text{Exergy change of flow} \end{aligned}$$



#### FIGURE 8–21

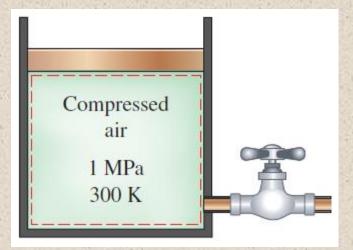
The *exergy* associated with *flow energy* is the useful work that would be delivered by an imaginary piston in the flow section.





#### FIGURE 8–22

The *energy* and *exergy* contents of (*a*) a fixed mass and (*b*) a fluid stream.



### Work Potential of Compressed Air in a Tank

 $X_1 = m\phi_1$ 

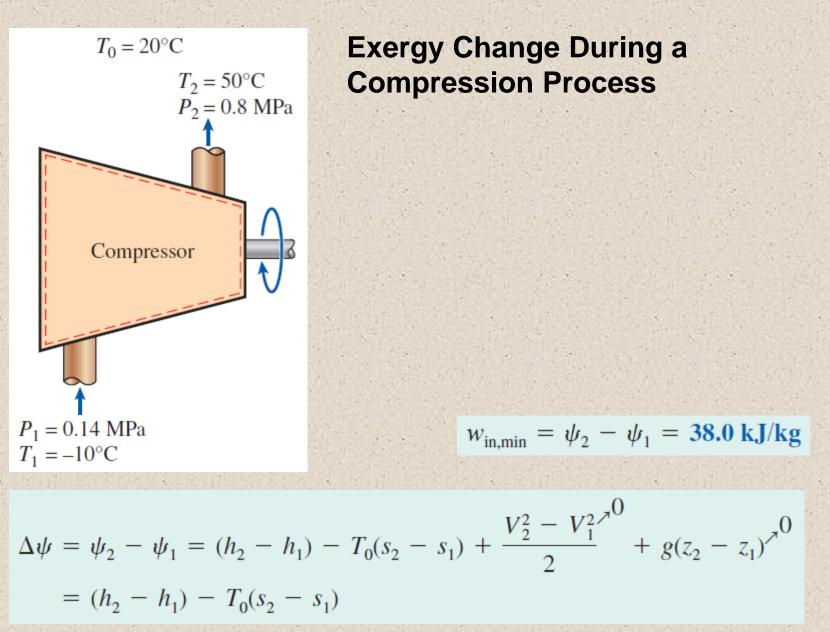
$$= m \left[ (u_1 - u_0)^{1/2} + P_0(v_1 - v_0) - T_0(s_1 - s_0) + \frac{V_1^{2/2}}{2} + gz_1^{1/2} \right]$$

$$= m[P_0(v_1 - v_0) - T_0(s_1 - s_0)]$$

$$P_0(v_1 - v_0) = P_0\left(\frac{RT_1}{P_1} - \frac{RT_0}{P_0}\right) = RT_0\left(\frac{P_0}{P_1} - 1\right)$$
 (since  $T_1 = T_0$ )

$$T_0(s_1 - s_0) = T_0 \left( c_p \ln \frac{T_1}{T_0} - R \ln \frac{P_1}{P_0} \right) = -RT_0 \ln \frac{P_1}{P_0} \quad \text{(since } T_1 = T_0)$$

$$\phi_1 = RT_0 \left(\frac{P_0}{P_1} - 1\right) + RT_0 \ln \frac{P_1}{P_0} = RT_0 \left(\ln \frac{P_1}{P_0} + \frac{P_0}{P_1} - 1\right) \qquad X_1 = m_1 \phi_1$$



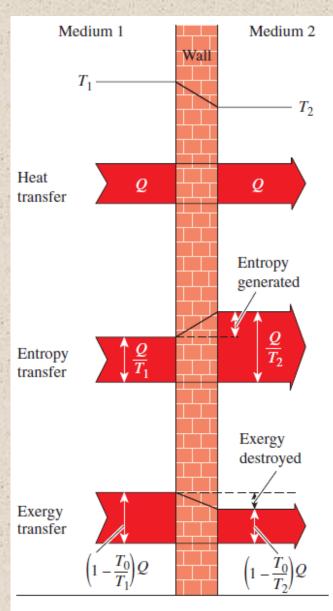
### EXERGY TRANSFER BY **HEAT, WORK, AND MASS**

### **Exergy by Heat Transfer, Q**

 $X_{\text{heat}} = \left(1 - \frac{T_0}{T}\right)Q$  Exergy transfer by heat

 $T_0$ 

 $X_{\text{heat}} = \int \left(1 - \frac{T_0}{T}\right) \delta Q$  When temperature is not constant



#### FIGURE 8–26

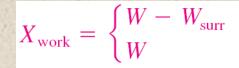
The transfer and destruction of exergy during a heat transfer process through a finite temperature difference.

#### **FIGURE 8–25**

The Carnot efficiency  $\eta_{\rm C} = 1 - T_0/T$ represents the fraction of the energy transferred from a heat source at temperature T that can be converted to work in an environment at temperature  $T_0$ .

Heat Source Temperature: T Energy transferred: E Exergy =  $\left(1 - \frac{T_0}{T}\right)E$ 

### Exergy Transfer by Work, W



(for boundary work)
(for other forms of work)

$$W_{\rm surr} = P_0(V_2 - V_1)$$

### Exergy Transfer by Mass, m

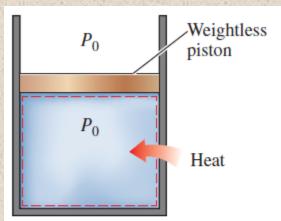
$$X_{\text{mass}} = m\psi \qquad \psi = (h - h_0) - T_0(s - s_0) + \frac{V^2}{2} + gz$$

$$\dot{X}_{\text{mass}} = \int_{A_c} \psi \rho V_n \, dA_c \qquad X_{\text{mass}} = \int \psi \, \delta m = \int_{\Delta t} \dot{X}_{\text{mass}} \, dt$$

h s m ms ψ mψ

#### FIGURE 8–28

Mass contains energy, entropy, and exergy, and thus mass flow into or out of a system is accompanied by energy, entropy, and exergy transfer.



#### FIGURE 8–27

There is no useful work transfer associated with boundary work when the pressure of the system is maintained constant at atmospheric pressure.

### THE DECREASE OF EXERGY PRINCIPLE AND EXERGY DESTRUCTION

Energy balance:

$$E_{\rm in}^{\mathcal{A}^0} - E_{\rm out}^{\mathcal{A}^0} = \Delta E_{\rm system} \rightarrow 0 = E_2 - E_1$$

Entropy balance:

$$S_{\text{in}}^{\mathcal{I}0} - S_{\text{out}}^{\mathcal{I}0} + S_{\text{gen}} = \Delta S_{\text{system}} \rightarrow S_{\text{gen}} = S_2 - S_1$$

Multiplying the second relation by  $T_0$  and subtracting it from the first one gives

$$-T_0 S_{\text{gen}} = E_2 - E_1 - T_0 (S_2 - S_1)$$
(8-29)

From Eq. 8–17 we have

$$X_{2} - X_{1} = (E_{2} - E_{1}) + P_{0}(V_{2} - V_{1})^{\nearrow^{0}} - T_{0}(S_{2} - S_{1})$$

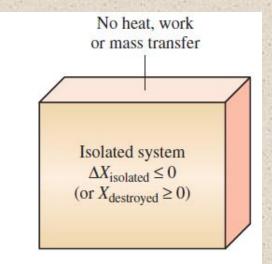
$$= (E_{2} - E_{1}) - T_{0}(S_{2} - S_{1})$$
(8-30)

since  $V_2 = V_1$  for an isolated system (it cannot involve any moving boundary and thus any boundary work). Combining Eqs. 8–29 and 8–30 gives

$$-T_0 S_{\text{gen}} = X_2 - X_1 \le 0 \tag{8-31}$$

since  $T_0$  is the thermodynamic temperature of the environment and thus a positive quantity,  $S_{\text{gen}} \ge 0$ , and thus  $T_0 S_{\text{gen}} \ge 0$ . Then we conclude that

 $\Delta X_{\text{isolated}} = (X_2 - X_1)_{\text{isolated}} \le 0$ 



#### FIGURE 8–29

The isolated system considered in the development of the decrease of exergy principle.

The exergy of an isolated system during a process always decreases or, in the limiting case of a reversible process, remains constant. In other words, it never increases and exergy is destroyed during an actual process. This is known as the decrease of exergy principle.

(8-32)

### **Exergy Destruction**

 $X_{\text{destroyed}} = T_0 S_{\text{gen}} \ge 0$ 

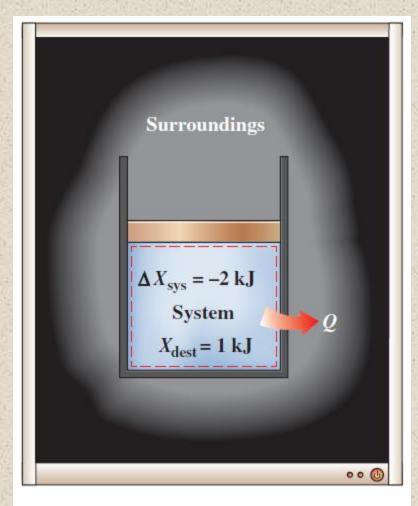
 $X_{\text{destroyed}} \begin{cases} > 0 & \text{Irreversible process} \\ = 0 & \text{Reversible process} \\ < 0 & \text{Impossible process} \end{cases}$ 

Exergy destroyed is a *positive quantity* for any actual process and becomes *zero* for a reversible process.

Exergy destroyed represents the lost work potential and is also called the *irreversibility* or *lost work*.

Can the exergy change of a system during a process be negative?

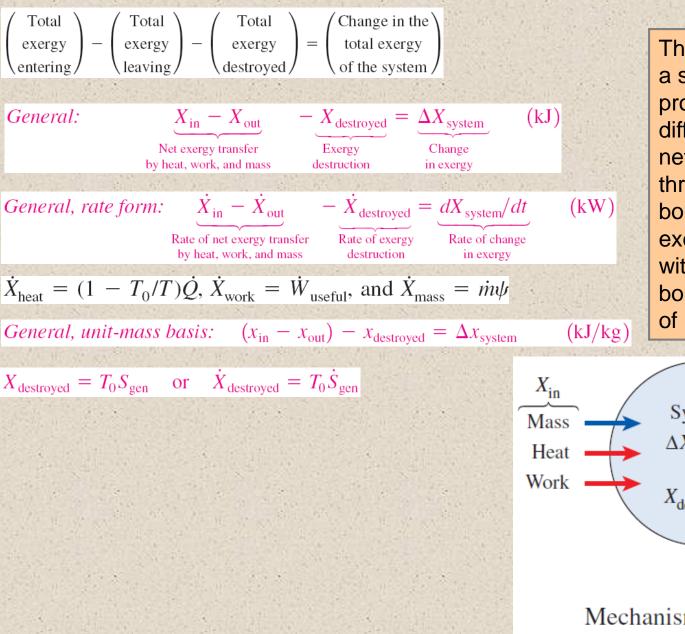
Consider heat transfer from a system to its surroundings. How do you compare exergy changes of the system and the surroundings?



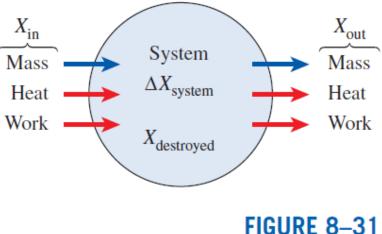
#### FIGURE 8-30

The exergy change of a system can be negative, but the exergy destruction cannot.

### **EXERGY BALANCE: CLOSED SYSTEMS**



The exergy change of a system during a process is equal to the difference between the net exergy transfer through the system boundary and the exergy destroyed within the system boundaries as a result of irreversibilities.



Mechanisms of exergy transfer.

Closed system:  

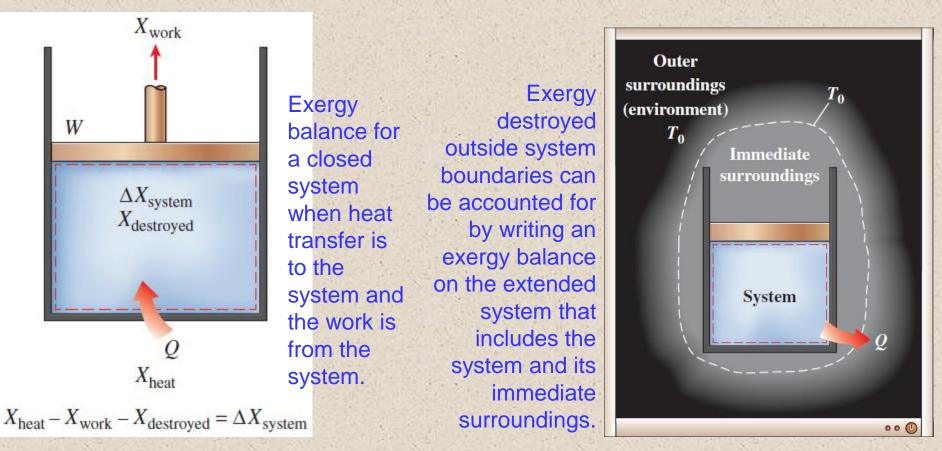
$$X_{\text{heat}} - X_{\text{work}} - X_{\text{destroyed}} = \Delta X_{\text{system}}$$
Closed system:  

$$\sum \left(1 - \frac{T_0}{T_k}\right) Q_k - \left[W - P_0(V_2 - V_1)\right] - T_0 S_{\text{gen}} = X_2 - X_2$$
Rate form:  

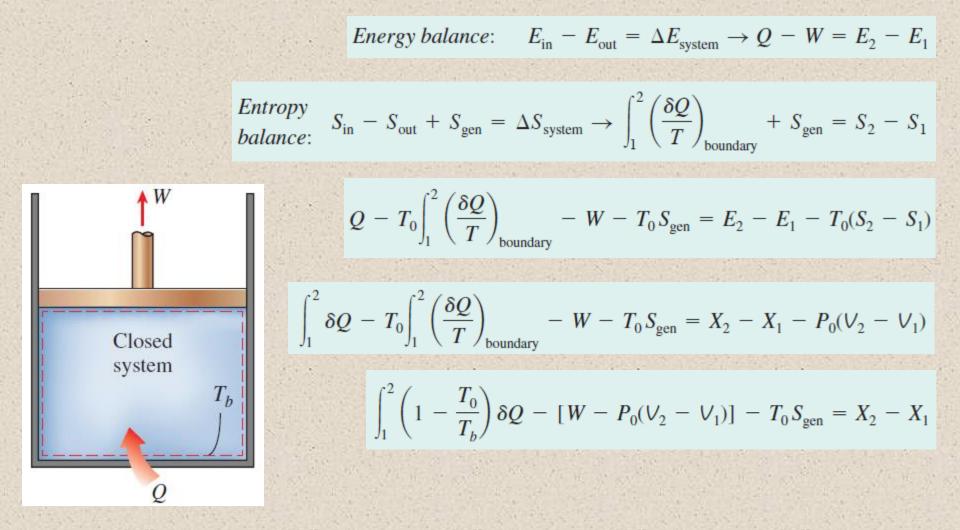
$$\sum \left(1 - \frac{T_0}{T_k}\right) \dot{Q}_k - \left(\dot{W} - P_0 \frac{dV_{\text{system}}}{dt}\right) - T_0 \dot{S}_{\text{gen}} = \frac{dX_{\text{system}}}{dt}$$

The heat transfer to a system and work done by the system are taken to be positive quantities.

 $Q_k$  is the heat transfer through the boundary at temperature  $T_k$  at location k.

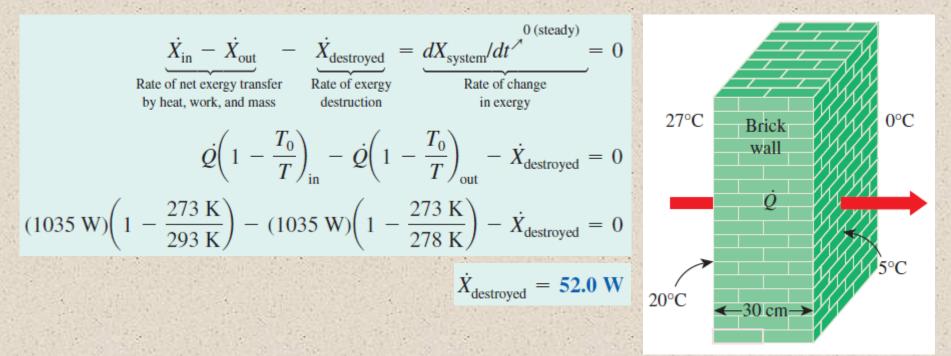


### **General Exergy Balance for Closed Systems**



24

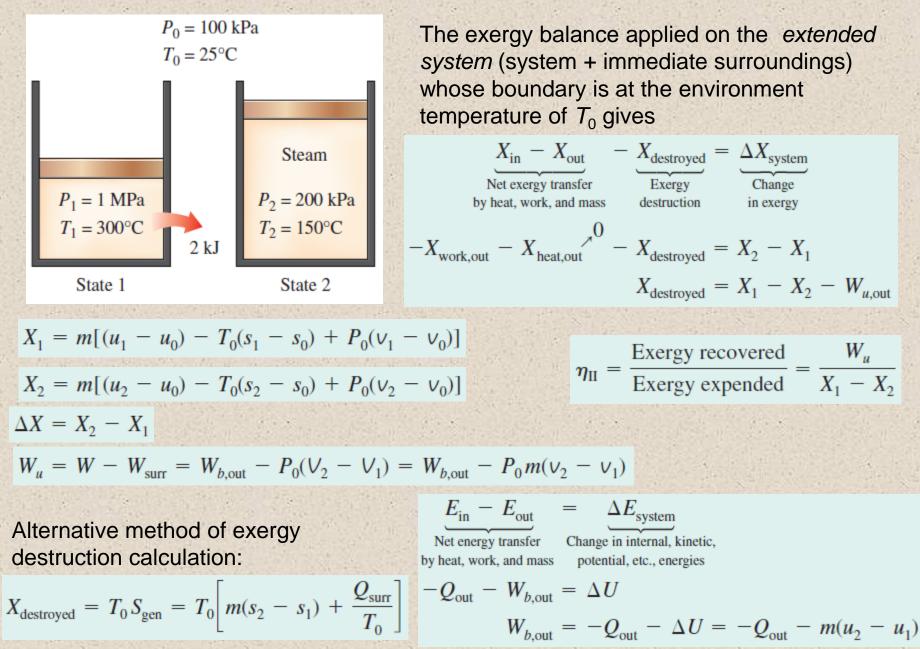
### **Exergy Destruction during Heat Conduction**



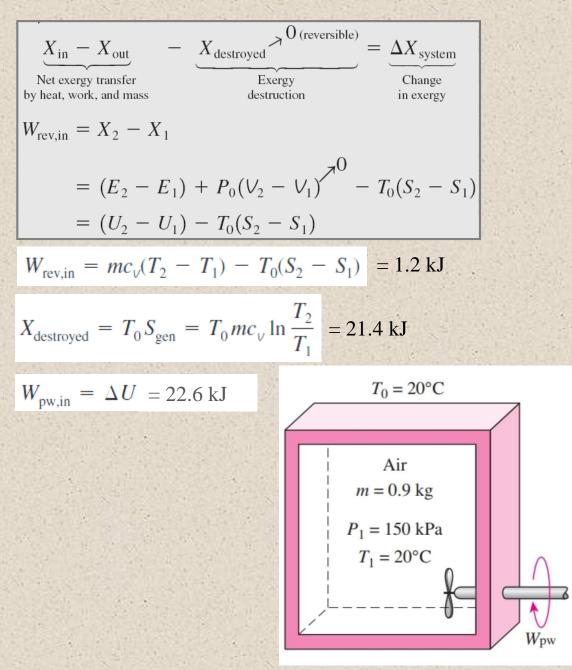
To determine the rate of total exergy destruction during this heat transfer process, we extend the system to include the regions on both sides of the wall that experience a temperature change. Then one side of the system boundary becomes room temperature while the other side, the temperature of the outdoors. The exergy balance for this *extended system* (system + immediate surroundings) is the same as that given above, except the two boundary temperatures are 300 and 273 K instead of 293 and 278 K, respectively. Then the rate of total exergy destruction becomes

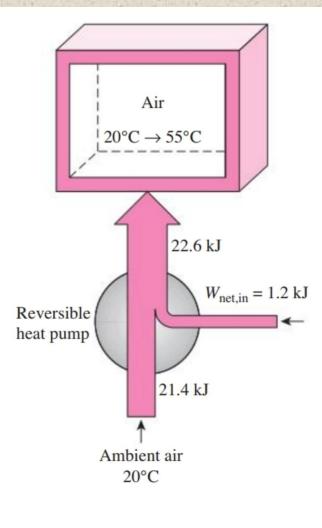
 $\dot{X}_{\text{destroyed,total}} = (1035 \text{ W}) \left( 1 - \frac{273 \text{ K}}{300 \text{ K}} \right) - (1035 \text{ W}) \left( 1 - \frac{273 \text{ K}}{273 \text{ K}} \right) = 93.2 \text{ W}$ 

### **Exergy Destruction During Expansion of Steam**



### **Exergy Destroyed During Stirring of a Gas**

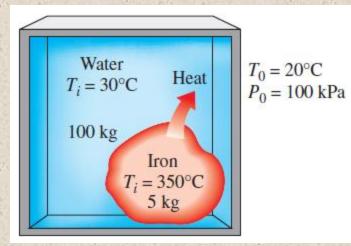




#### FIGURE 8–38

The same effect on the system can be accomplished by a reversible heat pump that consumes only 1.2 kJ of work.

### **Dropping a Hot Iron Block into Water**



$$\underbrace{\underbrace{E_{\text{in}} - E_{\text{out}}}_{\text{Net energy transfer}} = \underbrace{\Delta E_{\text{system}}}_{\text{Change in internal, kinetic, potential, etc., energies}} \\ 0 = \Delta U \\ 0 = (\Delta U)_{\text{iron}} + (\Delta U)_{\text{water}} \\ 0 = [mc(T_f - T_i)]_{\text{iron}} + [mc(T_f - T_i)]_{\text{water}} \\ 0 = [mc(T_f - T_i)]_{\text{iron}} + [mc(T_f - T_i)]_{\text{water}} \\ x = (U - U_0) - T_0(S - S_0) + P_0(V - V_0) \\ = mc(T - T_0) - T_0mc \ln \frac{T}{T_0} + 0 \\ = mc\left(T - T_0 - T_0 \ln \frac{T}{T_0}\right)$$

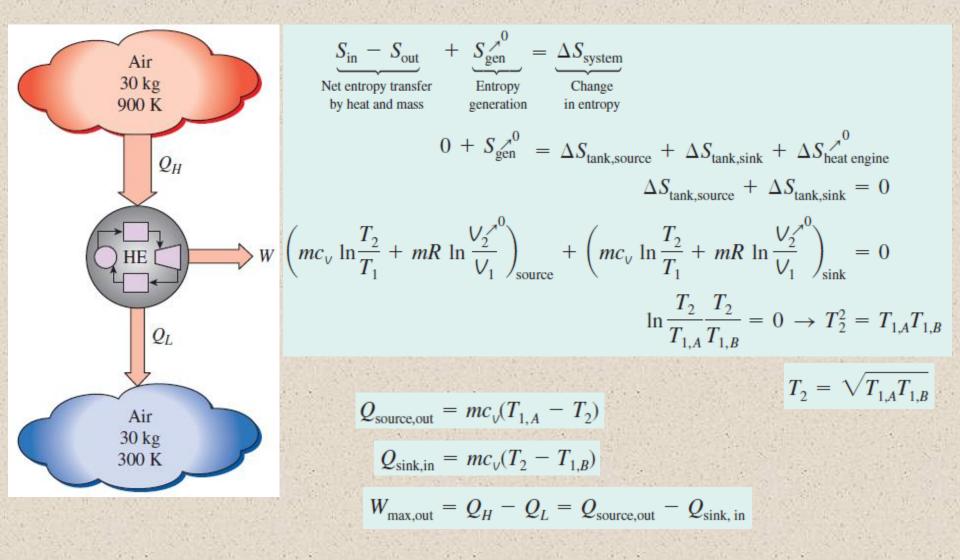
$$\underbrace{X_{\text{in}} - X_{\text{out}}}_{\text{Net exergy transfer}} - \underbrace{X_{\text{destroyed}}}_{\text{Exergy}} = \underbrace{\Delta X_{\text{system}}}_{\text{Change}}_{\text{in exergy}}$$

$$0 - X_{\text{destroyed}} = X_2 - X_1$$

$$X_{\text{destroyed}} = X_1 - X_2$$

28

### **Work Potential of Heat Transfer Between Two Tanks**



29

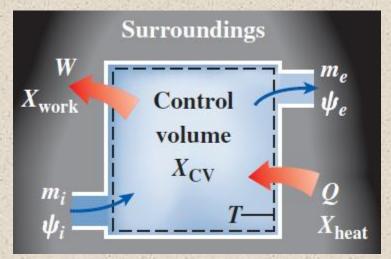
## **EXERGY BALANCE: CONTROL VOLUMES**

$$X_{\text{heat}} - X_{\text{work}} + X_{\text{mass,in}} - X_{\text{mass,out}} - X_{\text{destroyed}} = (X_2 - X_1)_{\text{CV}}$$

$$\sum \left(1 - \frac{T_0}{T_k}\right) Q_k - \left[W - P_0(V_2 - V_1)\right] + \sum_{\text{in}} m\psi - \sum_{\text{out}} m\psi - X_{\text{destroyed}} = (X_2 - X_1)_{\text{CV}}$$

$$\sum \left(1 - \frac{T_0}{T_k}\right) \dot{Q}_k - \left(\dot{W} - P_0 \frac{dV_{\rm CV}}{dt}\right) + \sum_{\rm in} \dot{m}\psi - \sum_{\rm out} \dot{m}\psi - \dot{X}_{\rm destroyed} = \frac{dX_{\rm CV}}{dt}$$

The rate of exergy change within the control volume during a process is equal to the rate of net exergy transfer through the control volume boundary by heat, work, and mass flow minus the rate of exergy destruction within the boundaries of the control volume.



#### FIGURE 8-41

Exergy is transferred into or out of a control volume by mass as well as heat and work transfer.

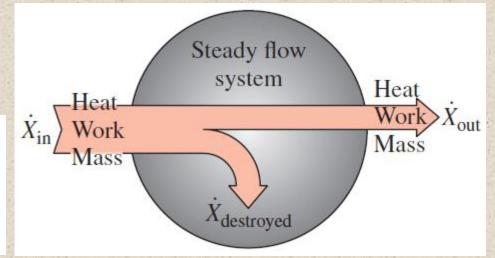
### **Exergy Balance for Steady-Flow Systems**

Most control volumes encountered in practice such as turbines, compressors, nozzles, diffusers, heat exchangers, pipes, and ducts operate steadily, and thus they experience no changes in their mass, energy, entropy, and exergy contents as well as their volumes. Therefore,  $dV_{CV}/dt = 0$  and  $dX_{CV}/dt = 0$  for such systems.

Steady-flow: 
$$\sum \left(1 - \frac{T_0}{T_k}\right)\dot{Q}_k - \dot{W} + \sum_{in}\dot{m}\psi - \sum_{out}\dot{m}\psi - \dot{X}_{destroyed} = 0$$
  
Single-stream:  $\sum \left(1 - \frac{T_0}{T_k}\right)\dot{Q}_k - \dot{W} + \dot{m}(\psi_1 - \psi_2) - \dot{X}_{destroyed} = 0$   
 $\psi_1 - \psi_2 = (h_1 - h_2) - T_0(s_1 - s_2) + \frac{V_1^2 - V_2^2}{2} + g(z_1 - z_2)$   
Per-unit  $\sum \left(1 - \frac{T_0}{T_k}\right)q_k - w + (\psi_1 - \psi_2) - x_{destroyed} = 0$ 

#### FIGURE 8–42

The exergy transfer to a steady-flow system is equal to the exergy transfer from it plus the exergy destruction within the system.



### **Reversible Work**

The exergy balance relations presented above can be used to determine the reversible work  $W_{rev}$  by setting the exergy destroyed equal to zero. The work W in that case becomes the reversible work.

General: 
$$W = W_{rev}$$
 when  $X_{destroyed} = 0$   
Single stream:  $\dot{W}_{rev} = \dot{m}(\psi_1 - \psi_2) + \sum \left(1 - \frac{T_0}{T_k}\right)\dot{Q}_k$  (kW)

Adiabatic, single stream:

$$\dot{W}_{\rm rev} = \dot{m}(\psi_1 - \psi_2)$$

The exergy destroyed is zero only for a reversible process, and reversible work represents the maximum work output for workproducing devices such as turbines and the minimum work input for work-consuming devices such as compressors.

### **Second-Law Efficiency of Steady-Flow Devices**

The second-law efficiency of various steady-flow devices can be determined from its general definition,  $\eta_{II} = (Exergy recovered)/(Exergy expended)$ . When the changes in kinetic and potential energies are negligible and the devices are adiabatic:

$$\eta_{\text{II,turb}} = \frac{w}{w_{\text{rev}}} = \frac{h_1 - h_2}{\psi_1 - \psi_2} \quad \text{or} \quad \eta_{\text{II,turb}} = 1 - \frac{T_0 s_{\text{gen}}}{\psi_1 - \psi_2} \quad \text{Turbine}$$

$$\eta_{\text{II,comp}} = \frac{w_{\text{rev,in}}}{w_{\text{in}}} = \frac{\psi_2 - \psi_1}{h_2 - h_1} \quad \text{or} \quad \eta_{\text{II,comp}} = 1 - \frac{T_0 s_{\text{gen}}}{h_2 - h_1} \quad \text{Compressor}$$

$$\eta_{\text{II,HX}} = \frac{\dot{m}_{\text{cold}}(\psi_4 - \psi_3)}{\dot{m}_{\text{hot}}(\psi_1 - \psi_2)} \quad \text{Heat}_{\text{exchanger}}$$

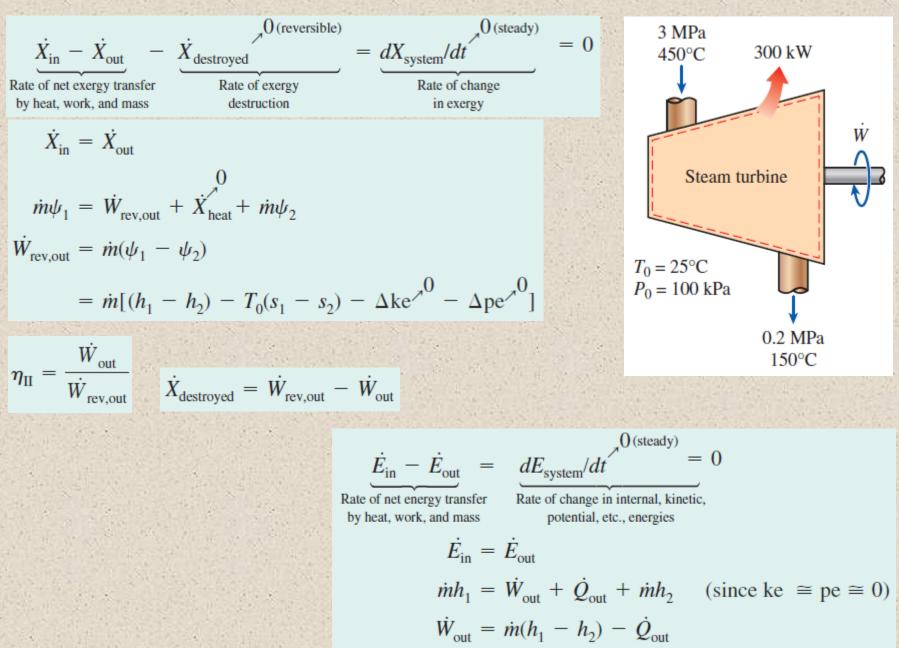
$$\eta_{\text{II,HX}} = 1 - \frac{T_0 \dot{S}_{\text{gen}}}{\dot{m}_{\text{hot}}(\psi_1 - \psi_2)}$$

$$\dot{S}_{\text{gen}} = \dot{m}_{\text{hot}}(s_2 - s_1) + \dot{m}_{\text{cold}}(s_4 - s_3)$$

$$\eta_{\text{II,mix}} = \frac{\dot{m}_3 \psi_3}{\dot{m}_1 \psi_1 + \dot{m}_2 \psi_2} \quad \text{Mixing}_{\text{chamber}}$$

$$\eta_{\text{II,mix}} = 1 - \frac{T_0 \dot{S}_{\text{gen}}}{\dot{m}_1 \psi_1 + \dot{m}_2 \psi_2} \quad \dot{S}_{\text{gen}} = \dot{m}_3 s_3 - \dot{m}_2 s_2 - \dot{m}_1 s_1$$

### Second-law analysis of a steam turbine



Exercy Destroyed During Mixing  

$$\frac{\dot{x}_{in} - \dot{x}_{out}}{\dot{x}_{in} - \dot{x}_{out}} + \frac{\dot{x}_{destroyed}}{(destroyed)} = dx_{system}/dt = 0$$

$$\frac{\dot{x}_{in} = \dot{x}_{out}}{\dot{n}_{1}\psi_{1} + \dot{n}_{2}\psi_{2}} = \dot{\psi}_{rev,out} + \dot{x}_{heat} + \dot{m}_{3}\psi_{3}$$

$$\frac{\dot{x}_{in} = \dot{x}_{out}}{\dot{n}_{rev,out} = \dot{m}_{1}\psi_{1} + \dot{m}_{2}\psi_{2} - \dot{m}_{3}\psi_{3}}$$

$$\frac{\dot{x}_{in} = \dot{x}_{out}}{\dot{n}_{rev,out} = \dot{m}_{1}\psi_{1} + \dot{m}_{2}\psi_{2} - \dot{m}_{3}\psi_{3}}$$

$$\frac{\dot{x}_{in} = \dot{x}_{out}}{\dot{n}_{rev,out} = \dot{m}_{1}\psi_{1} + \dot{m}_{2}\psi_{2} - \dot{m}_{3}\psi_{3}}$$

$$\frac{\dot{x}_{in} = \dot{x}_{out}}{\dot{n}_{rev,out} = \dot{m}_{1}\psi_{1} + \dot{m}_{2}\psi_{2} - \dot{m}_{3}\psi_{3}}$$

$$\frac{\dot{x}_{in} = \dot{x}_{out}}{\dot{n}_{rev,out} = \dot{m}_{1}\psi_{1} + \dot{m}_{2}\psi_{2} - \dot{m}_{3}\psi_{3}}$$

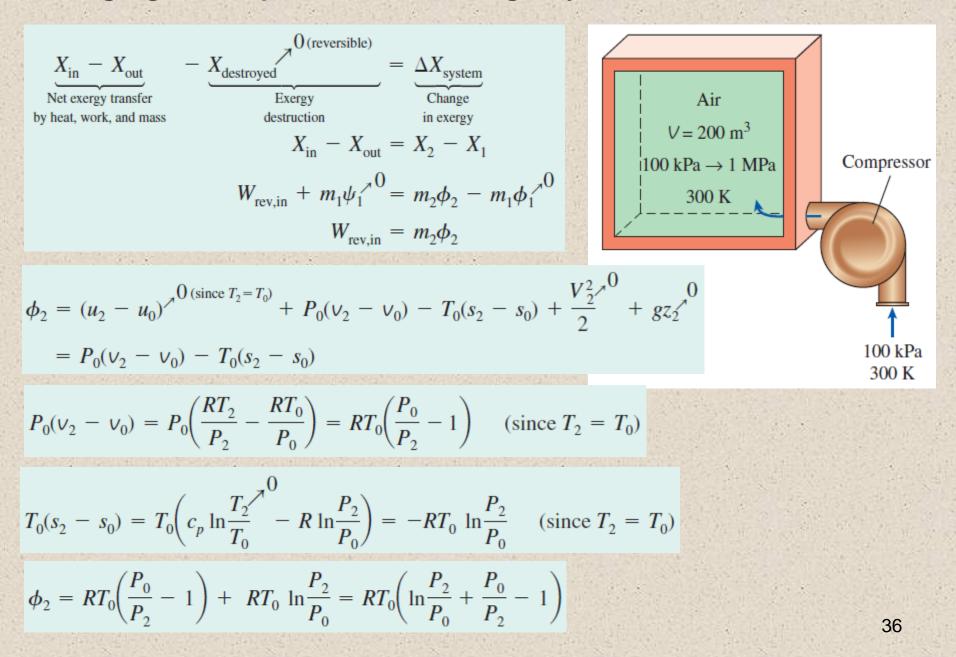
$$\frac{\dot{x}_{destroyed} = \dot{\psi}_{rev,out} - \dot{\psi}_{a} = T_{0}\dot{s}_{gen}}{\dot{x}_{destroyed} = \dot{\psi}_{rev,out}} - \dot{\psi}_{a} = T_{0}\dot{s}_{gen}}$$
Mass balance:
$$\frac{\dot{n}_{in} - \dot{n}_{out}}{\dot{n}_{in} - \dot{m}_{out}} = dR_{system}/dt = 0 \rightarrow \dot{m}_{1} + \dot{m}_{2} = \dot{m}_{3}}{destroyed} = \dot{\psi}_{rev,out}}$$

$$\frac{\dot{x}_{destroyed} = \dot{\psi}_{rev,out}}{\dot{x}_{destroyed} = \dot{\psi}_{rev,out}} + \dot{x}_{heat}, \dot{\mu}_{a} = 0$$

$$\frac{\dot{x}_{in} = \dot{x}_{out}}{\dot{n}_{in}h_{1} + \dot{m}_{2}h_{2} = \dot{m}_{3}h_{3} + \dot{Q}_{out}}(since \dot{W} = 0, ke = pe = 0)$$

$$\dot{Q}_{out} = \dot{m}_{ih}h_{1} + \dot{m}_{2}h_{2} - (\dot{m}_{1} + \dot{m}_{2})h_{3}$$
(35)

#### **Charging a Compressed Air Storage System**



# Summary

- Exergy: Work potential of energy
  - ✓ Exergy (work potential) associated with kinetic and potential energy
- Reversible work and irreversibility
- Second-law efficiency
- Exergy change of a system
  - Exergy of a fixed mass: Nonflow (or closed system) exergy
  - ✓ Exergy of a flow stream: Flow (or stream) exergy
- Exergy transfer by heat, work, and mass
- The decrease of exergy principle and exergy destruction
- Exergy balance: Closed systems
- Exergy balance: Control volumes
  - ✓ Exergy balance for steady-flow systems
  - ✓ Reversible work
  - ✓ Second-law efficiency of steady-flow devices