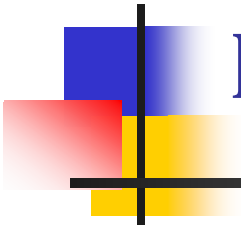


# PHYS-454

## 8 - The continuous spectrum of the potential wells.





# Momentum eigenstates of a free particle.

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- We can easily prove that the wave function

$$\psi = Ae^{ikx}$$

is an eigenfunction of the momentum operator

$$\hat{p} = -i\hbar d / dx$$

and describes particles which move with a specific momentum

$$p = \hbar k$$



## *The particle flux density-a*

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- In physics we know that the flux density of any physical quantity (mass, charge energy etc.) which is distribute in space with density  $\rho$  and moves with a velocity  $\mathbf{v}$  is given by:

$$\mathbf{J} = \rho \mathbf{v}$$



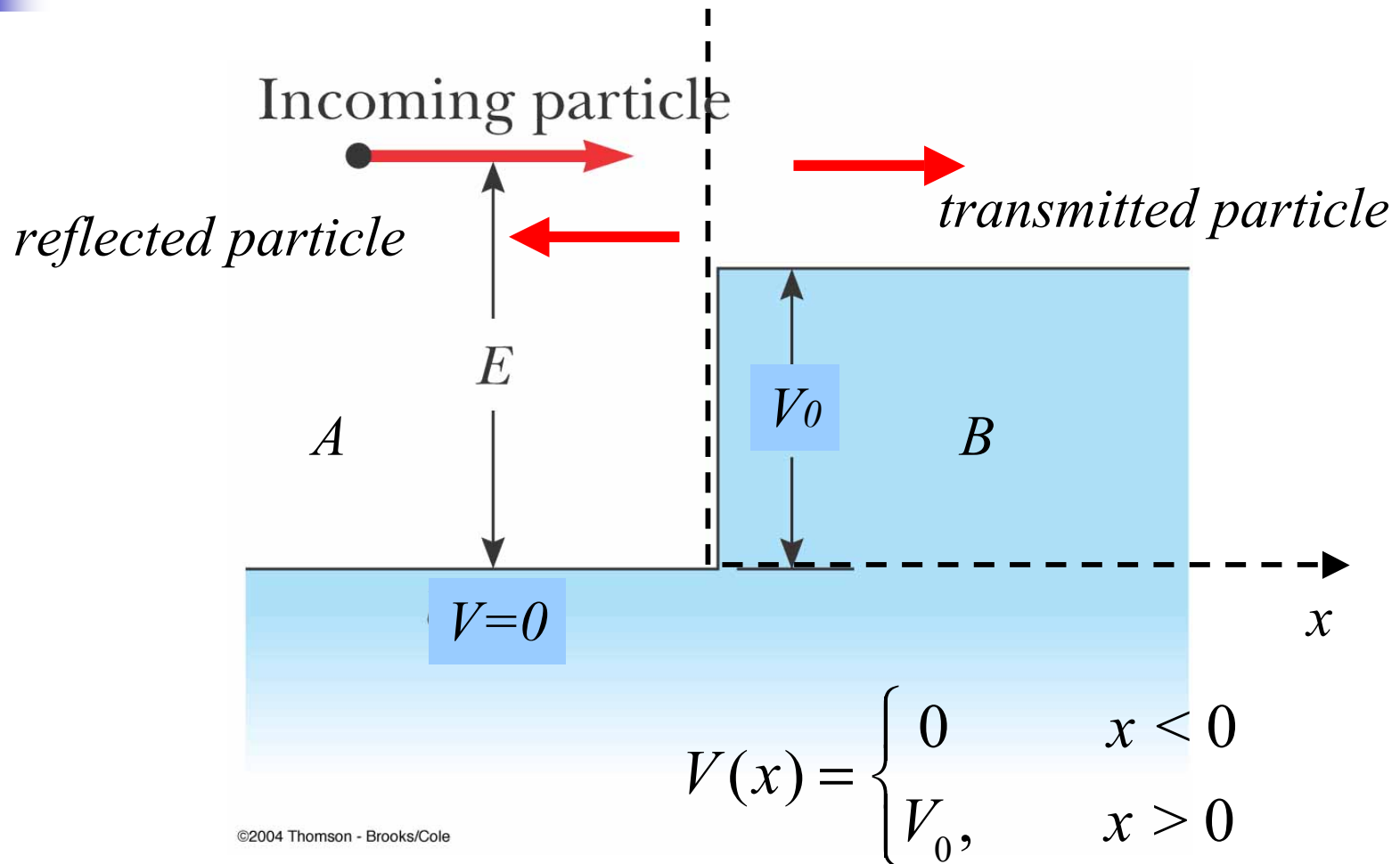
# *The particle flux density- $b$*

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- In an experiment in quantum mechanics we have a flux of particles which make up a *beam*. Then the particle flux density (i.e. the number of particles that cross per second a unit surface perpendicular to the flux) will be given by:

$$J = |\psi|^2 v = |A|^2 \frac{\hbar k}{m}$$

# The square potential step: The effect of reflection-a



# *The square potential step: The effect of reflection-b*

## ■ A) The case where $E > V_0$

The Schrödinger equation takes the form in the two regions A and B:

$$A: \quad \psi_A'' + \varepsilon \psi_A = \psi_A'' + k^2 \psi_A = 0$$

$$B: \quad \psi_B'' + (\varepsilon - U_0) \psi_B = \psi_B'' + k'^2 \psi_B = 0$$

$$\varepsilon = \frac{2mE}{\hbar^2}, \quad U_0 = \frac{2mV_0}{\hbar^2}$$

$$k^2 = \varepsilon, \quad k'^2 = \varepsilon - U_0 > 0$$



## *The square potential step: The effect of reflection-c*

In this case the general solutions of Schrödinger equations are:

$$\psi_A = A_+ e^{ikx} + A_- e^{-ikx}$$

$$\psi_B = B_+ e^{ik'x} + B_- e^{-ik'x}$$

The index + is for the constant of the wave  $\exp(ikx)$  which travel to the right and the index - for the constant of the wave  $\exp(-ikx)$  which goes to the left.



## *The square potential step: The effect of reflection-d*

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- Here there is an important point: The wave functions of the free particle are everywhere finite **but not** square integrable since they are not zero at infinity!
- Their physical meaning is that they are proper of describing *scattering* problems.
- The problem we study here is a scattering problem.
- This means that in the region B we must have only particles that go to the right .





## *The square potential step: The effect of reflection-e*

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- The coefficients  $A$  determine the densities of the incident and reflected beams respectively.
- Since the incident beam is regulated by us we can chose  $A_+ = 1$
- Thus we can chose the following conditions of scattering from the left

$$A_+ = 1 \quad B_- = 0$$



# *The square potential step: The effect of reflection-f*

- The general solution takes the form:

$$\psi_A = e^{ikx} + Ae^{-ikx}, \quad \psi_B = Be^{ik'x}$$

- In an experiment we care about the coefficients of reflection  $R$  and transmission  $T$ , given by

$$R = \frac{J_R}{J_I}, \quad T = \frac{J_T}{J_I}$$

- where

$J_I$  = the flux of incident beam

$J_R$  = the flux of reflected beam

$J_T$  = the flux of transmitted beam



## *The square potential step: The effect of reflection-g*

- The coefficients of reflection and transmission satisfy the relation

$$R + T = 1$$

- We can prove that the coefficients take the form

$$R = \left( \frac{k - k'}{k + k'} \right)^2 = \left( \frac{\sqrt{E} - \sqrt{E - V_0}}{\sqrt{E} + \sqrt{E - V_0}} \right)^2$$

$$T = \frac{4kk'}{(k + k')^2} = \frac{4\sqrt{E(E - V_0)}}{(\sqrt{E} + \sqrt{E - V_0})^2}$$



# *Graphical solutions of the eigenvalues-d*

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- **B) The case where  $E < V_0$**

The Schrödinger equation takes the form in the two regions A and B:

$$A: \quad \psi_A'' + \varepsilon \psi_A = \psi_A'' + k^2 \psi_A = 0$$

$$B: \quad \psi_B'' + (\varepsilon - U_0) \psi_B = \psi_B'' - \gamma^2 \psi_B = 0$$

$$\varepsilon = \frac{2mE}{\hbar^2}, \quad U_0 = \frac{2mV_0}{\hbar^2}$$

$$k^2 = \varepsilon, \quad \gamma^2 = U_0 - \varepsilon > 0$$



# *Graphical solutions of the eigenvalues-e*

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In this case the general solutions of  
Schrödinger equations are:

$$\psi_A = e^{ikx} + Ae^{-ikx}$$

$$\psi_B = Be^{-\gamma x}$$

$$A = \frac{k - i\gamma}{k + i\gamma}, \quad B = \frac{2k}{k + i\gamma}$$



# *Graphical solutions of the eigenvalues-f*

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- In this case the coefficients of reflection and transmission are:

$$R = |A|^2 = \left| \frac{k - i\gamma}{k + i\gamma} \right|^2 = 1 \quad T = 0$$

- This is not a surprising result. In region B the wave function is a decreasing exponential so the particle has no probability to reach the infinity and to be detected there.



## *Discussion-a*

### *The effect of reflection*

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- The most exciting aspect of the solution is the possibility of a particle to be reflected even if it has the required energy to go ahead. This is a consequence of the discontinuity of the potential at  $x=0$ .
- It is the analogue of the discontinuity of the refractive index when an e/m wave passes at the boundary of two transparent media.



## *Discussion-b*

### *The transmission coefficient at high energies*

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- When the energy  $E$  is very large the coefficient of transmission becomes equal to 1. This is not a surprise. At high energies the wavelength of a particle is very small, so the classical behavior prevails.





## *Discussion-c*

### *A paradox!*

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- A very strange result is that the coefficients of reflection **do not depend** on Planck's constant. This is a paradox since it is a quantum mechanical result.
- The origin of this paradox is the non-natural discontinuity of the potential at  $x=0$ . The wave behavior of matter is dominant when the wavelength is comparable in size with the discontinuity. But since the discontinuity has a zero "length" it will be always far smaller of the particle's wavelength.