

### Final Exam

Tuesday 08/Rabee'a'l Thani/1439 (26/12/2017)	PHYS 505	Academic year 1438-39H
8:00 – 11:00 am	Advanced Quantum Mechanics	First Semester

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*Please answer all questions*

1. You are given that the state  $|j_1 = 1, j_2 = 1, J = 2, M = 0 \rangle$  for two  $p$  electrons is given by:

$$|j_1 = 1, j_2 = 1, J = 2, M = -1 \rangle = \frac{1}{\sqrt{2}} \left\{ |j_1 = 1, m_1 = 0 \rangle |j_2 = 1, m_2 = -1 \rangle + |j_1 = 1, m_1 = -1 \rangle |j_2 = 1, m_2 = 0 \rangle \right\}$$

Find the state  $|j_1 = 1, j_2 = 1, J = 2, M = -2 \rangle$ .

You are given the lowering operator for all different types of angular momentum ( $K = J, j, l, s$ ):  $K_- |k, m_k \rangle = \hbar \sqrt{k(k+1) - m_k(m_k-1)} |k, m_k - 1 \rangle$ . Also for two particles the lowering operator is  $K_- = K_1^1 + K_2^2$

(10 marks)

**Solution:**

$$\begin{aligned} l_- |j_1 = 1, j_2 = 1, J = 2, M = -1 \rangle &= \hbar \sqrt{J(J+1) - M(M-1)} |j_1 = 1, j_2 = 1, J = 2, M = -2 \rangle = \\ &\hbar \sqrt{2(2+1) - (-1)(-1-1)} |j_1 = 1, j_2 = 1, J = 2, M = -2 \rangle = 2\hbar |j_1 = 1, j_2 = 1, J = 2, M = -2 \rangle \end{aligned} \quad (a)$$

$$\left(l_-^1 + l_-^2\right) \frac{1}{\sqrt{2}} \left\{ \left| j_1 = 1, m_1 = 0 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle + \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = 0 \right\rangle \right\} =$$

$$\frac{1}{\sqrt{2}} \left\{ \left( l_-^1 + l_-^2 \right) \left| j_1 = 1, m_1 = 0 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle + \left( l_-^1 + l_-^2 \right) \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = 0 \right\rangle \right\} =$$

$$\frac{1}{\sqrt{2}} \left\{ l_-^1 \left| j_1 = 1, m_1 = 0 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle + l_-^2 \left| j_1 = 1, m_1 = 0 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle + \right.$$

$$l_-^1 \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = 0 \right\rangle + l_-^2 \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = 0 \right\rangle \right\} =$$

(b)

$$\frac{1}{\sqrt{2}} \left\{ l_-^1 \left| j_1 = 1, m_1 = 0 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle + \left| j_1 = 1, m_1 = 0 \right\rangle \left( l_-^2 \left| j_2 = 1, m_2 = -1 \right\rangle \right) + \right.$$

$$l_-^1 \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = 0 \right\rangle + \left| j_1 = 1, m_1 = -1 \right\rangle \left( l_-^2 \left| j_2 = 1, m_2 = 0 \right\rangle \right) \right\} =$$

$$\frac{1}{\sqrt{2}} \left\{ \hbar \sqrt{1(1+1) - 0(0-1)} \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle + 0 + \right.$$

$$0 + \hbar \sqrt{1(1+1) - 0(0-1)} \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle \right\} =$$

$$\frac{1}{\sqrt{2}} \left\{ \hbar \sqrt{2} \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle + \hbar \sqrt{2} \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle \right\} =$$

$$2\hbar \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle$$

(b)

Comparing (a) and (b) we must have (a)=(b) thus,

$$2\hbar \left| j_1 = 1, j_2 = 1, J = 2, M = -2 \right\rangle = 2\hbar \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle \Rightarrow$$

$$\left| j_1 = 1, j_2 = 1, J = 2, M = -2 \right\rangle = \left| j_1 = 1, m_1 = -1 \right\rangle \left| j_2 = 1, m_2 = -1 \right\rangle$$

2. Three particles with spin  $s_1 = s_2 = s_3 = 1/2$  make up a system. A) Calculate the total spin of the system. (Hint: consider that the three particles are combined as follows: first you combine the two of them and then this system with the third one). B) If the three particles interact with the Hamiltonian  $H = A[\mathbf{s}_1 \cdot \mathbf{s}_2 + \mathbf{s}_1 \cdot \mathbf{s}_3 + \mathbf{s}_2 \cdot \mathbf{s}_3]$ , where  $A$  is a given constant. Find the possible energy eigenvalues. (Hint: start with the total angular momentum  $\mathbf{S}^2 = (\mathbf{s}_1 + \mathbf{s}_2 + \mathbf{s}_3)^2$ . You are given that for any type of angular momentum ( $K = J, j, l, s$ ):  $K^2 |k, m\rangle = \hbar^2 k(k+1) |k, m\rangle$ .

(10 marks)

### Solution:

A) We combine first the two spins  $s_1, s_2$  and we get:

$$s_{1,2} = |s_1 - s_2| \dots |s_1 + s_2| = \left| \frac{1}{2} - \frac{1}{2} \right| \dots \left| \frac{1}{2} + \frac{1}{2} \right| = 0, 1.$$

Then we combine this with the third one:

- i) For  $s_{1,2} = 0$  we have  $S = |s_{1,2} - s_3| \dots |s_{1,2} + s_3| = \left| 0 - \frac{1}{2} \right| \dots \left| 0 + \frac{1}{2} \right| = \frac{1}{2}$ .
- ii) For  $s_{1,2} = 1$  we have  $S = |s_{1,2} - s_3| \dots |s_{1,2} + s_3| = \left| 1 - \frac{1}{2} \right| \dots \left| 1 + \frac{1}{2} \right| = \frac{1}{2}, \frac{3}{2}$ .

B) For the calculation of the energy we have that:

$$\mathbf{S}^2 = (\mathbf{s}_1 + \mathbf{s}_2 + \mathbf{s}_3)^2 \Rightarrow \mathbf{S}^2 = \mathbf{s}_1^2 + \mathbf{s}_2^2 + \mathbf{s}_3^2 + 2\mathbf{s}_1 \cdot \mathbf{s}_2 + 2\mathbf{s}_2 \cdot \mathbf{s}_3 + 2\mathbf{s}_1 \cdot \mathbf{s}_3 \Rightarrow$$

$$2\mathbf{s}_1 \cdot \mathbf{s}_2 + 2\mathbf{s}_2 \cdot \mathbf{s}_3 + 2\mathbf{s}_1 \cdot \mathbf{s}_3 = \mathbf{S}^2 - \mathbf{s}_1^2 - \mathbf{s}_2^2 - \mathbf{s}_3^2 \Rightarrow$$

$$\mathbf{s}_1 \cdot \mathbf{s}_2 + \mathbf{s}_2 \cdot \mathbf{s}_3 + \mathbf{s}_1 \cdot \mathbf{s}_3 = \frac{1}{2}(\mathbf{S}^2 - \mathbf{s}_1^2 - \mathbf{s}_2^2 - \mathbf{s}_3^2)$$

Thus the energy eigenvalues are:

$$H = \frac{A\hbar^2}{2} [S(S+1) - s_1(s_1+1) - s_2(s_2+1) - s_3(s_3+1)]. \text{ So we have:}$$

- i) For  $S = 1/2$  we have

$$H = \frac{A\hbar^2}{2} \left[ \frac{1}{2} \left( \frac{1}{2} + 1 \right) - \frac{1}{2} \left( \frac{1}{2} + 1 \right) - \frac{1}{2} \left( \frac{1}{2} + 1 \right) - \frac{1}{2} \left( \frac{1}{2} + 1 \right) \right] = \frac{A\hbar^2}{2} \left( -\frac{3}{2} \right) = -\frac{3A\hbar^2}{4}.$$

- ii) For  $S = 3/2$  we have

$$H = \frac{A\hbar^2}{2} \left[ \frac{3}{2} \left( \frac{3}{2} + 1 \right) - \frac{1}{2} \left( \frac{1}{2} + 1 \right) - \frac{1}{2} \left( \frac{1}{2} + 1 \right) - \frac{1}{2} \left( \frac{1}{2} + 1 \right) \right] = \frac{A\hbar^2}{2} \left( \frac{15}{4} - \frac{3}{4} - \frac{3}{4} - \frac{3}{4} \right) = \frac{3A\hbar^2}{4}$$

3. We know that for a simple harmonic oscillator the unperturbed eigenfunctions and eigenenergies of the system are:  $\psi_n^{(0)} = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n \left( \sqrt{\frac{m\omega}{\hbar}} x \right) e^{-m\omega x^2/2\hbar}$  and  $E_n^{(0)} = \left( n + \frac{1}{2} \right) \hbar\omega$ .

We add a small perturbation to the system given by  $V(x) = \lambda^2 \delta(x)$ . Find the first order corrections of the energy eigenvalues.

You are given:

$$E_n^{(1)} = \left\langle \psi_n^{(0)} \middle| V \psi_n^{(0)} \right\rangle = \int_{-\infty}^{+\infty} \left( \psi_n^{(0)} \right)^* V \psi_n^{(0)} dx, \quad \int_{-\infty}^{+\infty} f(x) \delta(x) dx = f(0).$$

$$H_{2n}(0) = (-1)^n \frac{(2n)!}{n!}, \quad H_{2n+1}(0) = 0.$$

(10 marks)

**Solution:**

$$E_n^{(1)} = \int_{-\infty}^{+\infty} \left( \psi_n^{(0)}(x) \right)^* V \psi_n^{(0)}(x) dx = \int_{-\infty}^{+\infty} \left( \psi_n^{(0)}(x) \right)^* \lambda^2 \delta(x) \psi_n^{(0)}(x) dx = \int_{-\infty}^{+\infty} \lambda^2 \delta(x) \left| \psi_n^{(0)}(x) \right|^2 dx = \lambda^2 \int_{-\infty}^{+\infty} \lambda^2 \delta(x) \left| \psi_n^{(0)}(x) \right|^2 dx = \lambda^2 \left| \psi_n^{(0)}(0) \right|^2$$

But for the wave-function we have:

$$\psi_n^{(0)}(x) = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n \left( \sqrt{\frac{m\omega}{\hbar}} x \right) e^{-m\omega x^2/2\hbar} \Rightarrow \psi_n^{(0)}(0) = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n(0) \Rightarrow$$

$$\left| \psi_n^{(0)}(0) \right|^2 = \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \frac{1}{2^n n!} [H_n(0)]^2$$

Thus:

$$E_n^{(1)} = \lambda^2 \left| \psi_n^{(0)}(0) \right|^2 = \lambda^2 \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \frac{1}{2^n n!} [H_n(0)]^2.$$

- i) If  $n$  is even (let  $n = 2m, m = 0, 1, 2, \dots$ ) the energy shift is given by:

$$E_{2m}^{(1)} = \lambda^2 \left| \psi_n^{(0)}(0) \right|^2 = \lambda^2 \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \frac{1}{2^n n!} [H_n(0)]^2 =$$

$$\lambda^2 \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \frac{1}{2^{2m} (2m)!} \left[ (-1)^m \frac{(2m)!}{m!} \right]^2 = \lambda^2 \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \frac{(2m)!}{2^{2m} (m!)^2}$$

ii) If  $n$  is odd the energy shift is zero.

Thus only even states experience a first order correction to the energy.

4. A) For the scattering in the spherically symmetric potential  $V(\mathbf{r}) = A\delta(r - R)$ , (where  $A > 0$  a positive constant and  $\delta(r - R)$  the  $\delta$  function, calculate in the Born approximation the quantities  $f_B(\theta)$  and  $d\sigma / d\Omega$ .

(8 marks)

You are given that:

$$\int_0^\infty f(r)\delta(r - R)dr = f(R)$$

$$\frac{d\sigma}{d\Omega} = |f(\theta)|^2, \quad f(\theta) = -\frac{2m}{q\hbar^2} \int_0^\infty rV(r)\sin(qr)dr, \quad q = 2k\sin(\theta/2)$$

(10 marks)

**Solution:**

$$f(\theta) = -\frac{2m}{q\hbar^2} \int_0^\infty rV(r)\sin(qr)dr = -\frac{2m}{q\hbar^2} \int_0^\infty rA\delta(\mathbf{r} - \mathbf{R})\sin(qr)dr =$$

$$-\frac{2mA}{q\hbar^2} \int_0^\infty r\delta(\mathbf{r} - \mathbf{R})\sin(qr)dr = -\frac{2mA}{q\hbar^2} R\sin(qR)$$

Thus

$$\frac{d\sigma}{d\Omega} = |f(\theta)|^2 = -\left( \frac{2mA}{q\hbar^2} R\sin(qR) \right)^2 = \frac{4m^2 A^2}{q^2 \hbar^4} R^2 \sin^2(qR) =$$

$$\frac{4m^2 A^2}{4k^2 \hbar^4 \sin^2(\theta/2)} R^2 \sin^2(2kR\sin(\theta/2))$$

