

# Chapter #2: X-Ray Diffraction and the Reciprocal Lattice

## Lecture 3: Reciprocal Lattice – Formal Definition, Properties, and Diffraction Condition

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### 1. Introduction

In the previous lectures, we derived Bragg's law and established the condition for constructive interference in X-ray diffraction. The set of points in reciprocal space that satisfy Bragg's condition form what is known as the reciprocal lattice.

Therefore, reciprocal space provides a simpler geometric interpretation of diffraction phenomena.

#### Why do we need the reciprocal lattice?

We need the reciprocal lattice because waves (like X-rays, electrons, neutrons) interact with crystals in wave-vector (momentum) space, not real space. The reciprocal lattice is the natural language for describing diffraction, Bragg conditions, and electronic structure.

### 2. Definition of Reciprocal Lattice Vectors

If the primitive vectors of the direct lattice are  $\mathbf{a}_1$ ,  $\mathbf{a}_2$ , and  $\mathbf{a}_3$ , we define three new vectors  $\mathbf{b}_1$ ,  $\mathbf{b}_2$ , and  $\mathbf{b}_3$  such that:

$$\begin{array}{l} 1- \quad e^{i\Delta\vec{k}\cdot\vec{R}} = 1 \\ 2- \quad \mathbf{a}_i \cdot \mathbf{b}_j = 2\pi\delta_{ij} \end{array}$$

where  $\delta_{ij}$  is the Kronecker delta symbol:

$$\delta_{ij} = 1 \quad \text{if } i = j$$

$$\delta_{ij} = 0 \quad \text{if } i \neq j$$

$$\mathbf{a}_1 \cdot \mathbf{b}_1 = 2\pi$$

$$\mathbf{a}_1 \cdot \mathbf{b}_2 = 0$$

The vectors  $\mathbf{b}_1$ ,  $\mathbf{b}_2$ , and  $\mathbf{b}_3$  are called the reciprocal lattice primitive vectors.

### 3. Dimensional Analysis

The dimensions of direct lattice vectors are length (L), while the dimensions of reciprocal lattice vectors are inverse length (L<sup>-1</sup>). Therefore, the wave vector **k** has the same physical dimension as a reciprocal lattice vector.

$$a_i \sim \text{length}$$

$$b_i \sim \frac{1}{\text{length}}$$

This confirms that reciprocal space is naturally associated with momentum space.

**Remember that:**

$$k = \frac{2\pi}{\lambda}$$

$$p = \hbar k$$

### 4. Explicit Form of Reciprocal Lattice Vectors

The reciprocal lattice vectors can be written as:

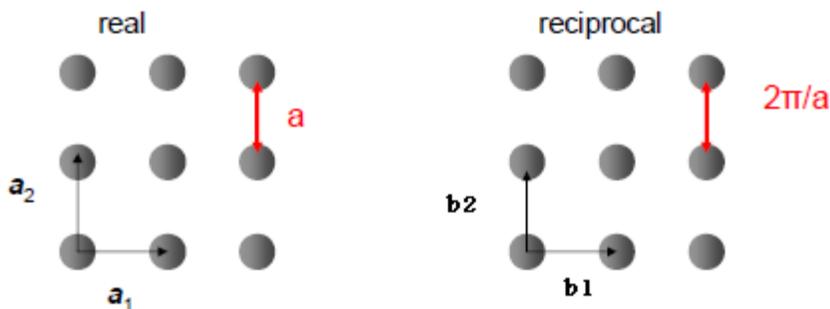
$$\vec{b}_1 = 2\pi \frac{\vec{a}_2 \times \vec{a}_3}{V} = 2\pi \frac{\vec{a}_2 \times \vec{a}_3}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)}$$

$$\vec{b}_2 = 2\pi \frac{\vec{a}_3 \times \vec{a}_1}{V}$$

$$\vec{b}_3 = 2\pi \frac{\vec{a}_1 \times \vec{a}_2}{V}$$

where  $V = \vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)$  is the volume of the primitive cell.

In crystallography, the factor  $2\pi$  is sometimes omitted, while in solid-state physics it is typically retained.



## 5. Orthogonality Properties

Let's check the validation of equations (1) and (2).

$$\vec{a}_1 \cdot \vec{b}_1 = \vec{a}_1 \cdot 2\pi \frac{\vec{a}_2 \times \vec{a}_3}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = 2\pi \frac{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = 2\pi$$

$$\vec{a}_1 \cdot \vec{b}_2 = \vec{a}_1 \cdot 2\pi \frac{\vec{a}_3 \times \vec{a}_1}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = 2\pi a^3 \frac{\hat{x} \cdot (\hat{z} \times \hat{x})}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = 2\pi a^3 \frac{\hat{x} \cdot \hat{y}}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = 0$$

- Each reciprocal vector is perpendicular to the plane formed by the other two direct lattice vectors. For example,  $\vec{b}_1$  is perpendicular to the plane formed by  $\vec{a}_2$  and  $\vec{a}_3$ .
- If the direct lattice vectors  $\vec{a}_1, \vec{a}_2, \text{ and } \vec{a}_3$  are mutually orthogonal, then the reciprocal lattice vectors also form an orthogonal set. In this special case, each reciprocal vector  $\vec{b}_i$  is parallel to the corresponding direct vector  $\vec{a}_i$ .
- However, in non-orthogonal lattices (i.e., Monoclinic, Triclinic), the reciprocal lattice vectors are not necessarily parallel to the direct lattice vectors.

## 6. General Reciprocal Lattice Vector

A general reciprocal lattice vector can be written as:

$$\vec{G} = m_1 \vec{b}_1 + m_2 \vec{b}_2 + m_3 \vec{b}_3$$

where  $m_1, m_2, m_3$  are integers.

We can also check the validations of equations (1) and (2).

$$e^{i\vec{G} \cdot \vec{R}} = e^{i(m_1 \vec{b}_1 + m_2 \vec{b}_2 + m_3 \vec{b}_3) \cdot (n_1 \vec{a}_1 + n_2 \vec{a}_2 + n_3 \vec{a}_3)}$$

$$\vec{R} = n_1 \vec{a}_1 + n_2 \vec{a}_2 + n_3 \vec{a}_3$$

Since;  $\vec{a}_i \cdot \vec{b}_j = 2\pi \delta_{ij}$

Then,  $e^{i2\pi(m_1 n_1 + m_2 n_2 + m_3 n_3)} = e^{i2\pi n} = 1$

Where  $n$  is an integer.

## 7. Diffraction Condition in Reciprocal Space

**Theory:** the set of reciprocal lattice vectors  $\vec{G}$  determines the possible X-ray diffractions.

From the diffraction condition derived in the previous lecture:

$$(\vec{k}' - \vec{k}) \cdot \vec{R} = \Delta \vec{k} \cdot \vec{R} = 2\pi n$$

Therefore, we can say that the scattering vector  $\Delta\vec{k}$  is equivalent to the reciprocal lattice vector  $\vec{G}$ .

$$(\Delta\vec{k} = \vec{G})$$

This means that diffraction occurs when the change in wave vector equals a reciprocal lattice vector. This condition is completely equivalent to Bragg's law.

Therefore, the set of reciprocal lattice vectors determines all possible X-ray reflections.

- This is the Laue condition for diffraction.

See some visualizations of the reciprocal lattices

[http://www.msm.cam.ac.uk/doitpoms/tlplib/reciprocal\\_lattice/reciprocal\\_lattice.php](http://www.msm.cam.ac.uk/doitpoms/tlplib/reciprocal_lattice/reciprocal_lattice.php)

[http://www.matter.org.uk/diffraction/geometry/lattice\\_vectors.htm](http://www.matter.org.uk/diffraction/geometry/lattice_vectors.htm)

## □ Reciprocal Lattices to SC, FCC, and BCC

|            | <u>Direct lattice</u>  | <u>Reciprocal lattice</u>  | <u>Volume</u> |
|------------|--|--|---------------|
| <u>SC</u>  | $\begin{cases} \mathbf{a}_1 = a\mathbf{x} \\ \mathbf{a}_2 = a\mathbf{y} \\ \mathbf{a}_3 = a\mathbf{z} \end{cases}$   | $\begin{cases} \mathbf{b}_1 = (2\pi/a)\mathbf{x} \\ \mathbf{b}_2 = (2\pi/a)\mathbf{y} \\ \mathbf{b}_3 = (2\pi/a)\mathbf{z} \end{cases}$  | $(2\pi/a)^3$  |
| <u>FCC</u> | $\begin{cases} \mathbf{a}_1 = \frac{1}{2}a(\mathbf{x} + \mathbf{y}) \\ \mathbf{a}_2 = \frac{1}{2}a(\mathbf{y} + \mathbf{z}) \\ \mathbf{a}_3 = \frac{1}{2}a(\mathbf{z} + \mathbf{x}) \end{cases}$   | $\begin{cases} \mathbf{b}_1 = \frac{2\pi}{a}(-\mathbf{x} + \mathbf{y} - \mathbf{z}) \\ \mathbf{b}_2 = \frac{2\pi}{a}(\mathbf{x} - \mathbf{y} + \mathbf{z}) \\ \mathbf{b}_3 = \frac{2\pi}{a}(\mathbf{x} + \mathbf{y} - \mathbf{z}) \end{cases}$ | $2(2\pi/a)^3$ |
| <u>BCC</u> | $\begin{cases} \mathbf{a}_1 = \frac{1}{2}a(\mathbf{x} + \mathbf{y} - \mathbf{z}) \\ \mathbf{a}_2 = \frac{1}{2}a(-\mathbf{x} + \mathbf{y} + \mathbf{z}) \\ \mathbf{a}_3 = \frac{1}{2}a(\mathbf{x} - \mathbf{y} + \mathbf{z}) \end{cases}$ | $\begin{cases} \mathbf{b}_1 = \frac{2\pi}{a}(\mathbf{y} + \mathbf{z}) \\ \mathbf{b}_2 = \frac{2\pi}{a}(\mathbf{x} + \mathbf{z}) \\ \mathbf{b}_3 = \frac{2\pi}{a}(\mathbf{x} + \mathbf{y}) \end{cases}$   | $4(2\pi/a)^3$ |

## 8. Exercises

1. Show that the reciprocal of the reciprocal lattice of a simple cubic lattice is again simple cubic.
2. Determine the Bravais lattice of the reciprocal lattice of BCC.
3. Determine the Bravais lattice of the reciprocal lattice of FCC.