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# First principle study of half metallic ferromagnetism and transport properties of spinel's $\text{ZnFe}_2(\text{S/Se})_4$ for spintronic

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## PAPER

First principle study of half metallic ferromagnetism and transport properties of spinel's  $\text{ZnFe}_2(\text{S/Se})_4$  for spintronicAbeer A AlObaid<sup>1</sup>, Tahani I Al-Muhimeed<sup>1</sup>, Abdur Rahim<sup>2</sup>, Ghazanfar Nazir<sup>3,4</sup>, S Bouzgarrou<sup>5,6</sup>, Abeer Mera<sup>7,8</sup>, A I Aljameel<sup>9</sup>, H H Hegazy<sup>10,11</sup>, G Murtaza<sup>2,12</sup>  and Qasim Mahmood<sup>13</sup> <sup>1</sup> Department of Chemistry, College of Science, King Saud University, PO Box 22452, Riyadh 11451, Saudi Arabia<sup>2</sup> Materials Modeling Lab, Department of Physics, Islamia College Peshawar 25120, Khyber Pakhtunkhwa, Pakistan<sup>3</sup> Department of Chemistry, Inha University, 100 Inharo, Incheon, 22212, Republic of Korea<sup>4</sup> Department of Physics, Korea Advanced Institute of Science and Technology (KAIST), Daejeon 34141, Republic of Korea<sup>5</sup> Laboratoire de Microélectronique et Instrumentation (UR 03/13-04), Faculté des Sciences de Monastir, Avenue de l'Environnement 5000 Monastir, Tunisia.<sup>6</sup> Department of Physics, College of Science, Qassim University, P. O. 64, Buraidah, Saudi Arabia<sup>7</sup> Department of Physics, College of Arts and Science, Prince Sattam Bin Abdulaziz University, Wadi Addawasir, 11991, Kingdom of Saudi Arabia, Saudi Arabia<sup>8</sup> Department of Physics, Faculty of Science, Kafrelsheikh University, Kafrelsheikh 33516; Egypt<sup>9</sup> Department of Physics, College of Science, Imam Mohammad Ibn Saud Islamic University (IMSIU), Riyadh 11623, Saudi Arabia<sup>10</sup> Department of Physics, Faculty of Science, King Khalid University, PO Box 9004, Abha, Saudi Arabia<sup>11</sup> Department of Physics, Faculty of Science, Al-Azhar University; 71524, Assiut; Egypt<sup>12</sup> Department of Mathematics & Natural sciences, Prince Mohammad Bin Fahd University, P. O. Box 1664, Alkhobar 31952, Saudi Arabia<sup>13</sup> Basic & Applied Scientific Research Center, and Department of Physics, Imam Abdulrahman Bin Faisal University, PO Box 1982, 31441 Dammam, Saudi ArabiaE-mail: [zafarforall2004@gmail.com](mailto:zafarforall2004@gmail.com)**Keywords:** spinels, DFT, Half metallicity, spintronics, thermoelectric efficiency**Abstract**

Half metallic ferromagnetism received remarkable attention due to its immense technological applications in spintronic. In this perspective, thiospinels  $\text{ZnFe}_2\text{S/Se}_4$  is addressed by employing density functional theory (DFT) for spintronics and thermoelectric devices. The optimized energy versus volume leads to the confirmation of the Ferromagnetic (FM) states stability of both spinels. The formation energy confirms the thermodynamic stability. TDOS and PDOS are determined to confirm spin polarization and half-metallic ferromagnetism. Ferromagnetism is explored by exchange energies and magnetic moments. In addition, thermoelectric characteristics are explored by electrical and thermal conductivities, Seebeck coefficient, and power factor to evaluate their potential in thermoelectric applications.

**1. Introduction**

Spintronic is an emerging field in the scientific research community and is defined as a novel type of electronics that manipulates electrons by charge and spin of the electrons [1–3]. Its recent development begins since in 1988 with the invention of GMR. [4]. Half metal can be defined in terms of spin polarization (SP). Spin-polarized leads to resolving the heat problems that occur in IC and various electronic devices. Materials with high spin polarization, high magnetic phase transition (close to room temperature), and high magnetic moment are required to utilize in spintronic technology. Materials which do not exhibit spin polarization are diamagnetic because of unchanged DOS in both channels. However, asymmetric exist in ferromagnetic materials because of the exchange mechanism which limits SP less than unity. On the other hand, half-metallic contains states at the Fermi level in one channel and insulators in the second channel [5]. Thus, the materials persist 100% polarization and are highly recommended materials for spintronics. The discovery of half metals made a great contribution to the field of spintronics due to their technological implications in efficient spintronic devices including magnetic recording, electronic memory, and effective magnetic sensor [6–8].

Over the last few decades, researchers have been focused on many compounds, including (half, full and quaternary) Heusler alloys [9, 10], dilute semiconductor [11], transition metal oxides [12], Spinel FeO<sub>4</sub> [13], perovskites BaBkO<sub>3</sub> [14], double perovskites [15] and halide double perovskites [16]. Looking for half metal's materials with high Curie temperature at RT has been focused on for more than 15 years. A special class of materials called spinels with cubic structures generalized by the formula AB<sub>2</sub>O<sub>4</sub> with unique properties has attracted more attention [17, 18]. To tune the physical properties of such complex crystals their present multiple degrees of freedom. This class of materials has diverse properties including its potential application in photocatalytic degradation [19–21], charge storing devices [22], thermoelectricity [23] energy conversion applications [24]. Thiospinels have several desirable optical properties, including transparency across at large photosensitivity [25], and nonlinear optical susceptibility [26] all of which make them capable candidates for optoelectronic devices. Thiospinels (Zn/Hg) In<sub>2</sub>S<sub>4</sub> are used in terms of transport properties for energy conversion devices [27]. Moreover, XIn<sub>2</sub>S<sub>4</sub> (X = Cd, Mg) has the potential ability for optoelectronic devices [28].

Ferromagnetism materials with semiconducting nature are the best option for spintronic applications. The ferromagnetic spinel semiconductors AB<sub>2</sub>X<sub>4</sub> (A = Cu, Hg, Cd, B = Cr, and X = S, Se) are well-known magnetic semiconductors, for instance, CuCr<sub>2</sub>Se<sub>4</sub> [29–31]. Further, the Curie temperature range 84 K–430 K, and ferromagnetic semiconductor behavior make them appealing for spintronics [32]. Spinel ferromagnetic materials play a key role regarding ferroelectric, memory, and spintronic applications [33, 34]. Owing to Colossal Magnetoresistance in the transition metal ferro-spinel materials offers a fascinating platform for researchers [35, 36]. The ferromagnetic characteristics of MgCr<sub>2</sub>S<sub>4</sub> and MgMn<sub>2</sub>S<sub>4</sub> are considered excellent FM [37, 38]. Q. Mahmood *et al* worked on ferromagnetism in MgCr<sub>2</sub>X<sub>4</sub> (X = S, Se) [39]. Recently, Mehmood *et al* demonstrated the magnetic behavior of MgPr<sub>2</sub>(S/Se)<sub>2</sub> for spintronic applications by DFT calculation [40]. Moreover, various properties of Spinel Ca(V/Mn)<sub>2</sub>S<sub>4</sub> have been investigated computationally including structural, magnetic, and thermoelectric [41].

The above vibrant review ensures that there is no literature available on Spinel's ZnFe<sub>2</sub>S/Se<sub>4</sub> regarding their half-metallic ferromagnetism and transport properties. In this article, we used TB-mBJ to determine electronic band structures, the partial and total density of states, and the half-metallic character of Zn-based thiospinels. The ferromagnetism due to exchange mechanism and spin polarization without clustering effect is the primary goal of our research. Thermoelectric parameters are investigated by employing BoltzTrap code. The study of the effect of temperature, and thermal conductivity on the spin of electrons large effect the reliability of the devices. Therefore, the effects of these parameters are also discussed in detail. The theoretical data presented in this work can cover the lack of physical properties information on the ZnFe<sub>2</sub>S/Se<sub>4</sub> compounds.

## 2. Computational details

We employed the full-potential linearized augmented plane wave method FP-APW realized in the WIEN2K package within the framework of DFT to calculate various properties including structural, electronic, magnetic, and thermoelectric [42–44]. For reducing interatomic forces in electronic structures, optimization has been done through PBEsol approximation. The structural characteristics are calculated by Perdew–Burke–Ernzerhof PBEsol [45, 46] and modified Becke Johnson potential of Tran and Blaha (TB-mBJ) [47] because PBEsol was analyzed the ground state properties more accurately but underestimate the electronic bandgap. Therefore, to improve the bandgap accurately, the TB-mBJ potential has been implemented over the PBEsol approximation. The solution of the electronic system inside the muffin-tin region is taken spherically harmonic. The k-mesh of the order 12 × 12 × 12 has been selected as the point at which the energy is released by the system [48]. The convergence parameters are adjusted as  $K_{\max} \times R_{\text{MT}} = 8.0$ , muffin radius ( $R_{\text{MT}}$ ), and  $K_{\max}$  wave vector in the reciprocal lattice, Gaussian factor  $G_{\max} = 18$ , and angular momentum  $\ell_{\max} = 10$ . The change was converged up to 0.001 m Ry. Furthermore, the TB-mBJ converged energy and the optimized electronic structures were used to calculate the thermoelectric behavior through classical Boltzmann transport theory by using BoltzTraP Package [49].

## 3. Result and discussion

### 3.1. Structural analysis

To optimize the cubic crystal structure of thiospinels ZnFe<sub>2</sub>(S/Se)<sub>4</sub> with Fd-3m space group we have used PBEsol approximation. The crystal structure of the Zn-based thiospinels is presented in figure 1. The relaxation process is used for atomic positions with Zn (0.125, 0.125, and 0.125), Fe (0.5, 0.5, and 0.5) and O (0.25, 0.25, and 0.25) which are in accordance with the space group Fd-3m.

The optimized ZnFe<sub>2</sub>(S/Se)<sub>4</sub> for FM and AFM states are presented in figure 2. The planned analogy of FM and AFM revealed that ZnFe<sub>2</sub>(S/Se)<sub>4</sub> possess ferromagnetic nature owing to their lower energy. Similar results are

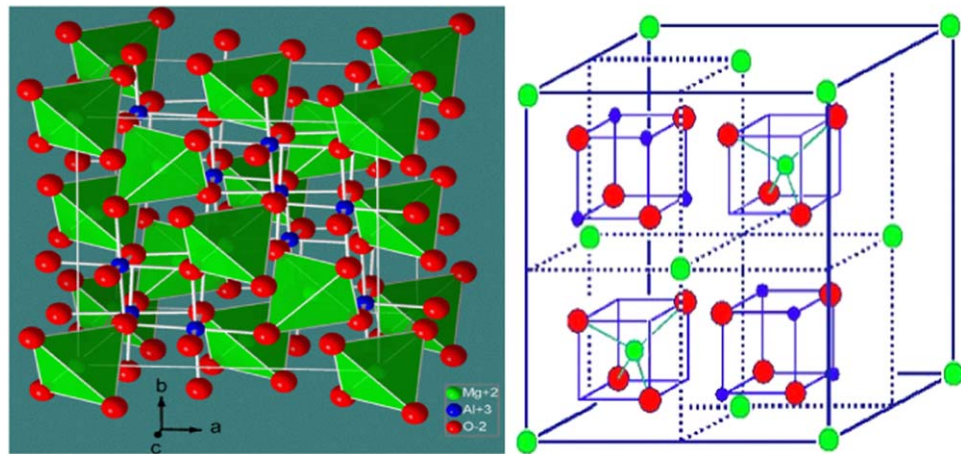


Figure 1. the crystal structure in atomic form and polyhedral form of cubic Spinels  $\text{ZnFe}_2\text{S}/\text{Se}_4$ .

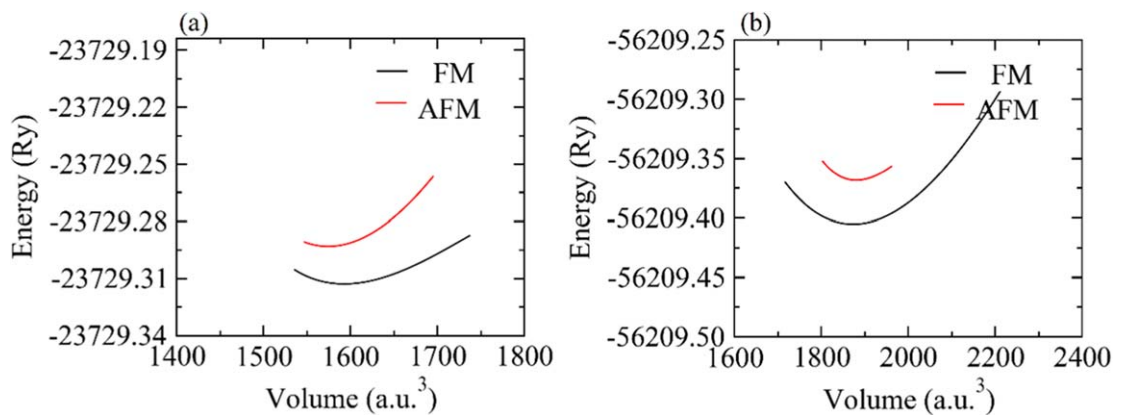


Figure 2 . Optimized plots in FM (black color) and AFM (red color) states of (a)  $\text{ZnFe}_2\text{S}_4$  and (b)  $\text{ZnFe}_2\text{Se}_4$ .

evident from the literature of  $\text{XCr}_2\text{O}_4$  ( $\text{X} = \text{Zn}, \text{Cd}$ ), and  $\text{AV}_2\text{O}_4$  ( $\text{A} = \text{Zn}, \text{Cd}, \text{Hg}$ ) which ensure the lowest energy in FM states confirm its stability. Therefore, the consistency of calculated results with the existing literature is proof of the reliability of the study [50, 51].

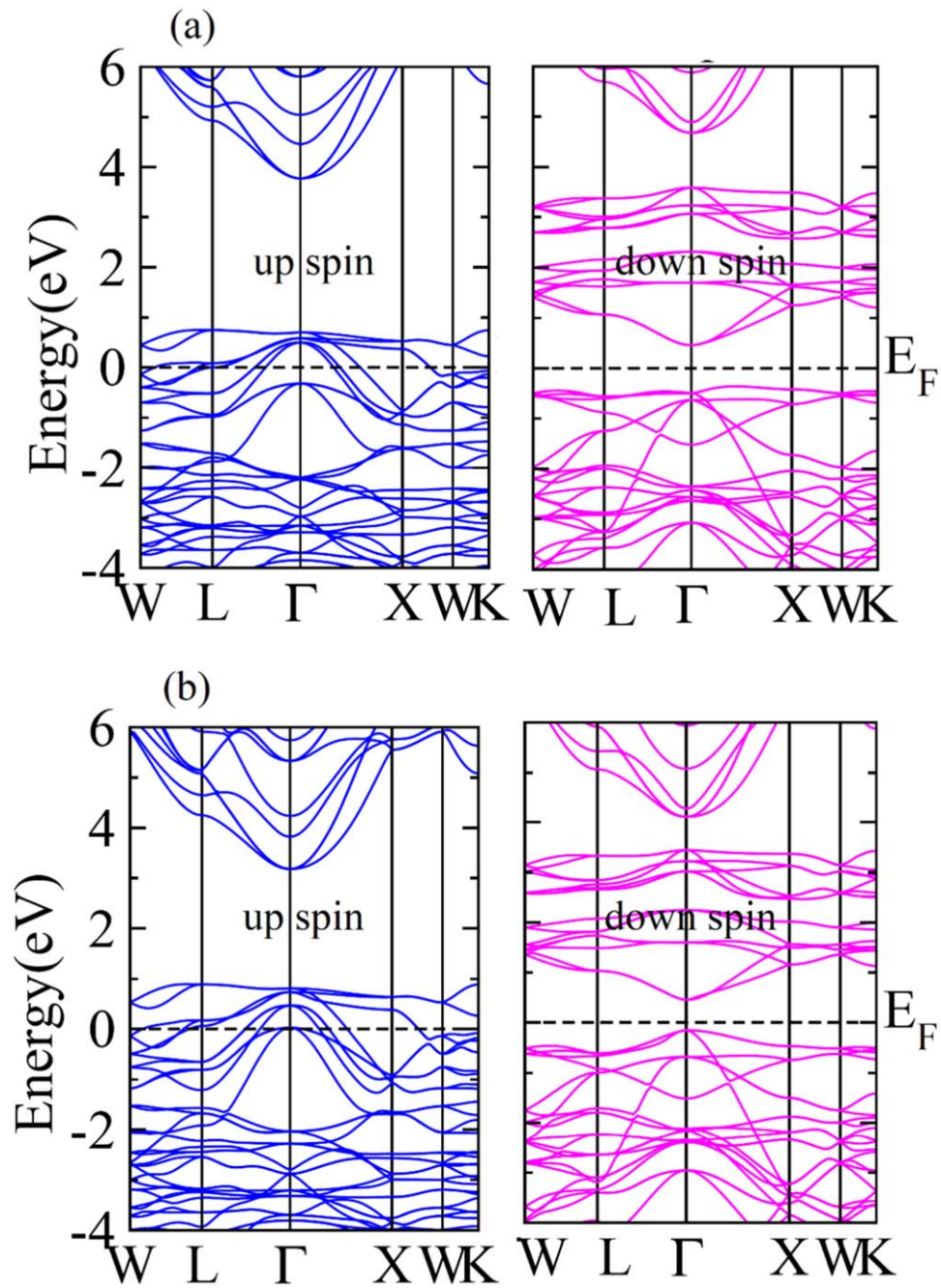
The formation energy of studied compounds have been calculated by the relation

$$\Delta H_f = E_{\text{Total}}(\text{Zn}_1\text{Fe}_m\text{S}/\text{Se}_n) - lE_{\text{Zn}} - mE_{\text{Fe}} - nE_{\text{S/Se}} \quad (1)$$

Where  $E_{\text{Total}}(\text{Zn}_1\text{Fe}_m\text{S}/\text{Se}_n)$ ,  $E_{\text{Zn}}$ ,  $E_{\text{Fe}}$  and  $E_{\text{S/Se}}$  are the energies of  $\text{ZnFe}_2\text{S}/\text{Se}_4$ , Zn, Fe, and S/Se, respectively. The computed values are  $-3.4$  eV for  $\text{ZnFe}_2\text{S}_4$ , and  $-3.1$  eV for  $\text{ZnFe}_2\text{Se}_4$  which confirm that studied materials are thermodynamically favorable [52]. Furthermore, we have computed the Curie temperature by the Classical Heisenberg model whose mathematical form is  $T_c = \Delta E / 3K_B$ , where  $\Delta E$  is the energy difference between paramagnetic and ferromagnetic states,  $K_B$  is Boltzmann constant [53]. The computed values are 315 K, and 305 K which show the room temperature ferromagnetism.

### 3.2. Electronic bandstructure

Band structure analysis is an important step in the understanding of material nature and its suitability for device applications. Due to the ferromagnetic nature of these materials, we have calculated spin-up and spin-down band structures presented in figures 3(a), (b). The valence band maxima (VBM) lie at the K-symmetry point while conduction band minima (CBM) at Gamma point in the up-spin channel of  $\text{ZnFe}_2\text{S}_4$  shown in figure 3(a). In addition, the VBM lies between Gamma-X points while CBM lies at Gamma points in down-spin channels showing a semiconductor nature. Similarly, the up spin and down spin channels are plotted for  $\text{ZnFe}_2\text{Se}_4$  schemed in figure 3(b). In the up-spin channel, the VBM stays on the K-symmetry point crossing the Fermi level and CBM stays on Gamma-symmetry. Whereas, in the down spin channel both the VBM and CBM lie at Gamma symmetry point having a narrow gap about Fermi level leading to semiconducting nature. Therefore, by



**Figure 3** . band structures of (a)  $\text{ZnFe}_2\text{S}_4$  and (b)  $\text{ZnFe}_2\text{Se}_4$  in up spin ( $\uparrow$ ) and down spin ( $\downarrow$ ) channels.

combining both insulating and semiconducting behavior of up spin and down spin channels form ferromagnetic materials.

Materials with maximum spin polarizability are desirable for spintronic applications. The spin polarizability be calculated by the given mathematical relation [54],

$$P = \frac{N(\uparrow)E_F - N(\downarrow)E_F}{N(\uparrow)E_F + N(\downarrow)E_F} \times 100\% \quad (2)$$

Where  $N(\uparrow)$  and  $N(\downarrow)$  stand for the total density of states (TDOS). To take a clearer picture of half-metallicity and exchange mechanism in these thiospinels, we have plotted the total and partial density of states (PDOS) in figures 4(a)–(b). The Fermi level is present in the valence band for the up spin channel, and in the forbidden region for the down spin channel. The value of  $P$  is 100% for both compounds make them promise for spintronics. The distribution of valence and conduction states reveals that TDOS indicates half-metallic behavior in these spinels-like band structures.



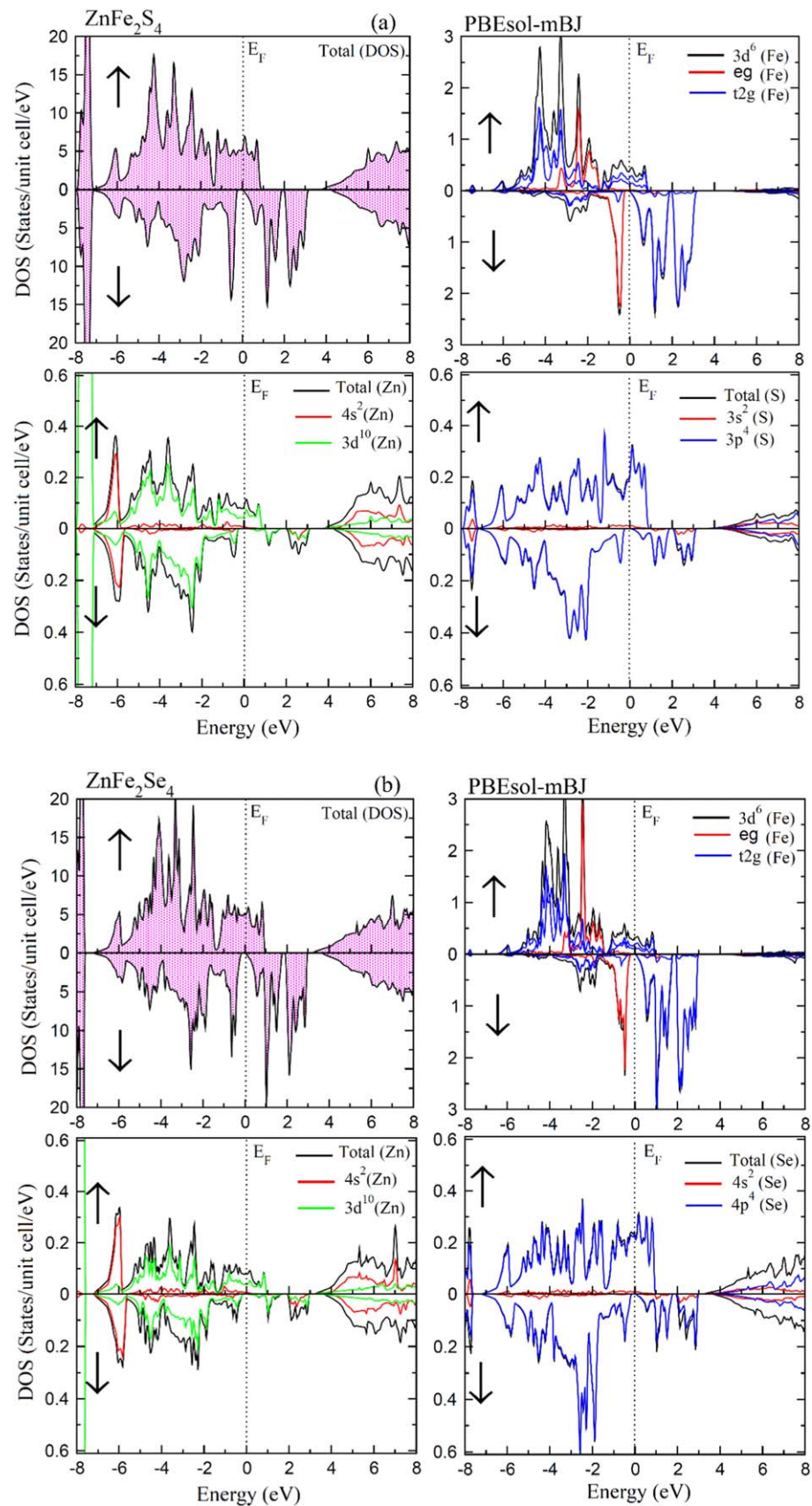


Figure 4. (a) DOS of  $\text{ZnFe}_2\text{S}_4$  in up spin (↑) and down spin (↓) channel. (b) DOS of  $\text{ZnFe}_2\text{Se}_4$  in up spin (↑) and down spin (↓) channels.

### 3.3. Magnetic properties

Magnetic properties of materials are crucial for determining their suitability to spintronic applications. To explain magnetism in materials two potential exchanges, play a vital role, e.g double-exchange and super

**Table 1.** The computed values of exchange energies ( $\Delta E_{\text{crystal}}$ ,  $\Delta_x(d)$ ,  $\Delta_x(pd)$ , and exchange constants ( $N_0\alpha$  and  $N_0\beta$ ) for  $\text{ZnFe}_2\text{S}/\text{Se}_4$ .

Compounds	$(\Delta E_{\text{crystal}})$	$\Delta_x(d)$	$\Delta_x(pd)$	$N_0\alpha$	$N_0\beta$
$\text{ZnFe}_2\text{S}_4$	-1.80	3.65	-0.027	-0.97	-0.34
$\text{ZnFe}_2\text{Se}_4$	-2.13	3.51	-0.046	-0.95	-0.37

exchange [55]. Super-exchange is responsible for anti-ferromagnetism while the double exchange generates ferromagnetism [56–58]. The FM nature of present thiospinels is confirmed through an optimization process. This FM nature arises due to the double exchange mechanism in these materials, where transition metal's d orbital split into degenerate triplet  $t_{2g}$  ( $d_{xy}$ ,  $d_{yz}$ ,  $d_{zx}$ ) states and doublet states  $e_g$  ( $d_z^2$  and  $d_x^2 - d_y^2$ ). The  $3d^6$  states of Fe atom contribute sharp peaks in the up and down spin channel ranges from (1 to -6) eV showing metallic behavior in up spin channel.

The splitting of  $e_g$  (Fe) and  $t_{2g}$  (Fe) in the down spin channel cause a narrow bandgap which shows the presence of semiconducting nature in spin-down configuration. The occurrence of  $e_g$  (Fe) and  $3d$  (Fe) states at separate energies ensures the magnetic moment of electrons exist which causes ferromagnetic character. From PDOS the  $3d^6$ ,  $e_g$ ,  $t_{2g}$  states of Fe are majorly responsible FM nature. In the individual PDOS of Zn, the valence band of total (Zn) states possesses high peaks crossing Fermi level, which confirm metallic nature in the up-spin configuration, while in the downward configuration it shows semiconducting behavior. However, the contribution of  $4s^2$  (Zn) is minimum in both configurations. The total and partial contribution of the S atom is displayed in figure 4, where  $3p^4$  state majorly presents high peaks in both channels.

By replacing S with Se, a similar pattern of results is obtained for  $\text{ZnFe}_2\text{Se}_4$  through the density of states. The TDOS holds metallic nature in the majority spin region while a semiconducting nature in the minority spin region due to hybridization among Fe, Zn, and Se. It is important to highlight the PDOS of  $\text{ZnFe}_2\text{Se}_4$ , which are sketched for Fe, Zn, and Se individually in figure 4(b). The first is one is for Fe where the  $3d^6$  state possesses high peaks in the low range of energy states (lower than -1) and crossing the Fermi level with lower intensity in the up-spin channel. It can be seen from figure 4 that narrow gap raised due to hybridization of doublet state of  $e_g$  (Fe) and triplet state of  $t_{2g}$  (Fe) in bonding and in an anti-bonding state in spin-down configuration near the  $E_F$ , respectively. Further, Zn total and  $3d^{10}$  (Zn) major peaks fluctuated in the lower energy range in up and down spin configuration. These peaks shifting beyond Fermi level with lower energy causes metallic nature in up spin configuration. The contribution due to  $4s^2$  (Zn) is lower near the Fermi level. Finally, the individual PDOS of Se demonstrates the Total (Se),  $4s^2$  (Se), and  $4p^4$  (Se) states. Total (Se) plays an important role in anti-bonding states in high energy states, while the  $4s^2$  (Se) state indicates a small contribution. The  $4p^4$  (Se) state is responsible for major peaks in both channels with metallic in up spin channel and semiconducting in a downward channel. Overall, this half-metallic character is due to the separating of d states orbitals of Fe and Zn in the existence of the external force of four Se atoms. As Fe-d states split into  $e_g$  ( $d_z^2$  and  $d_x^2 - d_y^2$ ) states  $t_{2g}$  ( $d_{xy}$ ,  $d_{yz}$ , and  $d_{zx}$ ) states [58]. Thus, crystal field energy  $\Delta CF$  arises from the splitting of  $e_g$  states and  $t_{2g}$  states and can be defined in terms of the difference between the two states i.e., ( $\Delta E = e_g - t_{2g}$ ) [59].

The behavior of ferromagnetism provokes by this crystal field energy and can be decreased by  $\Delta_x(d)$  is presented in table 1. The condition was evident for introducing ferromagnetism [60].

It is important to highlight the term exchange energy  $\Delta_x(pd)$  among the d states of Fe/Zn and 4p state of Se. The negative value confirms the occurrence of ferromagnetism. From table 1, it can be seen that  $\Delta CF$  decreases from  $\text{ZnFe}_2\text{S}_4$  to  $\text{ZnFe}_2\text{Se}_4$ , which indicates that Se-based spinels are more favorable for ferromagnetism. The exchange constants  $N_0\alpha$  and  $N_0\beta$  are calculated by [61].

$$N_0\alpha = \frac{\Delta E^c}{xS}, N_0\beta = \frac{\Delta E^v}{xS} \quad (3)$$

Where x and S stand for the concentration and magnetic moment of Fe atom. While  $\Delta E^c = E_{\downarrow}^c - E_{\uparrow}^c$  and  $\Delta E^v = E_{\downarrow}^v - E_{\uparrow}^v$  are the energies at VB and CB edges. In table 1. One can see the calculated values of  $N_0\alpha$  and  $N_0\beta$ . According to Zenger's exchange model (extensively used theory of the ferromagnetism in ferromagnetic semiconductors) [62, 63] the negative value of  $N_0\beta$  show the magnetic impurity through the energy gap in the down spin channel with lower energy. The total magnetic moment of these compounds was calculated to analyze ferromagnetic strength which arises due to MM of individual atoms and interstitial regions. From table 2. It is obvious that Fe is the most contributor to the TMM in both spinels while minor contribution comes from the Zn and interstitial regions.

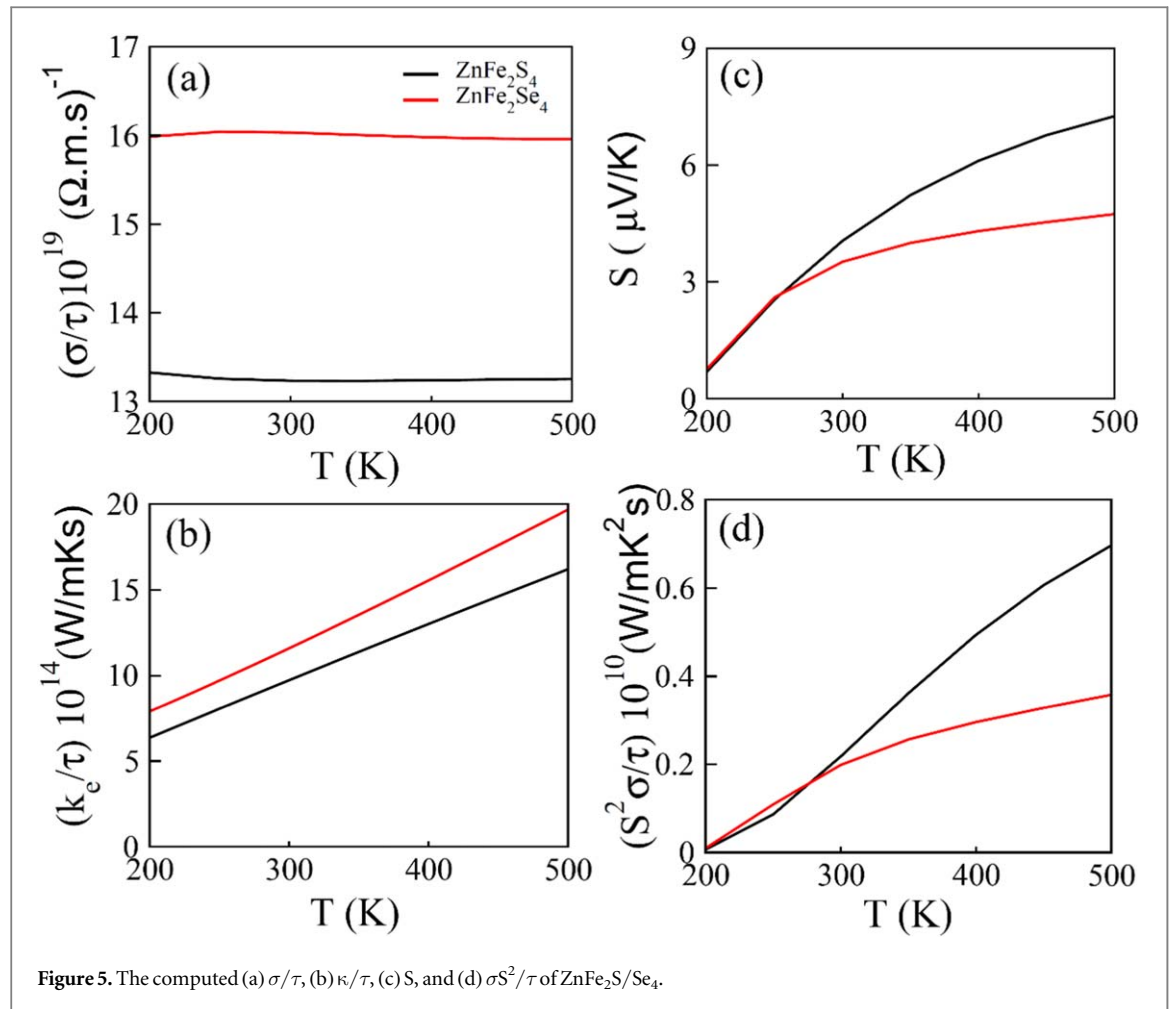


Figure 5. The computed (a)  $\sigma/\tau$ , (b)  $\kappa/\tau$ , (c)  $S$ , and (d)  $\sigma S^2/\tau$  of ZnFe<sub>2</sub>S/Se<sub>4</sub>.

Table 2. The magnetic moments for, Zn, Fe, S/Se. ZnFe<sub>2</sub>S<sub>4</sub> and ZnFe<sub>2</sub>Se<sub>4</sub>.

	Total ( $\mu_B$ )	Int ( $\mu_B$ )	Cs ( $\mu_B$ )	Fe ( $\mu_B$ )	(S/Se $\mu_B$ )
ZnFe <sub>2</sub> S <sub>4</sub>	4.00	0.057	0.001	3.03	0.023
ZnFe <sub>2</sub> Se <sub>4</sub>	4.00	0.331	0.003	2.80	0.027

### 3.4. Thermoelectric properties

To overcome the global energy crises, the transformation of heat energy into useful electrical energy is an effective way. Therefore, we investigated the thermoelectric behavior of spinels ZnFe<sub>2</sub>(S/Se)<sub>4</sub> through BoltzTrap code [49]. The thermoelectric parameters are plotted against temperature shown in figures 5(a)–(d). The plotted electrical conductivity against temperature demonstrates a small variation in the range of 200 K to 300 K and remains steady for high temperature in the case of both spinels. The thermal conductivity ( $\kappa$ ) of materials is the heat flow ( $q$ ) due to the temperature gradient following the Fourier law  $q = -k \nabla_x(T)$ .

Here, the computed thermal conductivity value increases linearly with the temperature range presented in figure 5(b) for both ZnFe<sub>2</sub>S<sub>4</sub> and ZnFe<sub>2</sub>Se<sub>4</sub> from  $6.5 \times 10^{14} (W/mKs)$  to  $16 \times 10^{14} (W/mKs)$  and  $7.5 \times 10^{14} (W/mKs)$  to  $20 \times 10^{14} (W/mKs)$ , respectively. Further, the Seebeck coefficient  $S$  explains voltage due to temperature gradient and can be calculated via the given equation.

$$S = \left( \frac{8}{3eh^2} \right) \pi^2 K_B^2 m^* T \left( \frac{\pi}{3n} \right)^{\frac{1}{2}} \quad (4)$$

Where  $e$ ,  $h$ ,  $K_B$ ,  $m^*$ ,  $T$ , and  $n$  are electronic charge, Planck constant, Boltzmann constant, effective mass, absolute temperature, and carrier concentration, respectively. Seebeck coefficient of ZnFe<sub>2</sub>(S/Se)<sub>4</sub> are plotted in figure 5(c). The computed value of  $S$  has the same value for both spinels up to 260 K. Next, the values of ZnFe<sub>2</sub>S<sub>4</sub> and ZnFe<sub>2</sub>Se<sub>4</sub> differed and reached at  $7.2 (\mu V/K)$  and  $4.5 (\mu V/K)$  at 500 K, respectively. Furthermore, the Power factor plays a key role in an understanding of thermoelectric performance with a mathematical expression of  $P = S^2 \sigma/\tau$ . The calculated value of  $P$  for ZnFe<sub>2</sub>S<sub>4</sub> started from  $0 W/mK^2s$  at 200 K and obtained a



maximum value of  $0.7 \text{ W} / \text{mK}^2 \text{s}$  at 500 K. For  $\text{ZnFe}_2\text{Se}_4$  the estimated values of  $P$  lies in range of  $(0 - 0.35) \text{ W} / \text{mK}^2 \text{s}$  in the whole temperature range. Hence, the figure of merit  $ZT$  is directly proportional to the power factor. Thus, materials with high power factors are suitable for energy conversion devices. In our case,  $\text{ZnFe}_2\text{S}_4$  is a potential candidate for such an application.

## 4. Conclusion

In this research article, we have investigated structural, electronic, magnetic, and thermoelectric properties of  $\text{ZnFe}_2(\text{S}/\text{Se})_4$  for spintronic and thermoelectric applications. The optimization of energy versus volume plots reveals that these spinels are favorable for ferromagnetism. The electronic band structures were calculated in terms of spin-up and down orientation with a narrow bandgap. The half-metallicity was confirmed by TDOS and PDOS with one channel having 100% spin polarization. The magnetic moment on nonmagnetic and interstitial sites has been reported. The electrical conductivity has not shown much variation with temperature. However, thermal conductivity, Seebeck coefficient, and the power factor increased with increasing temperature leading to high thermoelectric performance.

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## Data availability statement

The data that support the findings of this study are available upon reasonable request from the authors.

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