



# Ch.10: Rotation of a Rigid Object About a Fixed Axis

Physics 103: Classical Mechanics

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# 1. Rigid Body

## 2. Angular Position, Velocity, and Acceleration

## 3. Rotational Motion with Constant Angular Acceleration

## 4. Angular and Linear Quantities

## 5. Rotational Kinetic Energy

## 6. Calculation of Moments of Inertia

## 7. Torque

## 8. Relationship Between Torque and Angular Acceleration

## 9. Work, Power, and Energy in Rotational Motion

# 1.1 Definition

- A rigid body is an object that does not deform or change its shape during motion. The distance between any two points in a rigid body remains constant regardless of the motion of the body.



## 1.2 Point Objects to Rigid Objects

- So far we have been simplify objects as point masses to make calculations easier.
- However, when dealing with rigid bodies, we need to consider their size and shape, as well as how they rotate about an axis.
- Therefore, we can no longer treat them as point masses, and we need to use rotational dynamics to describe their motion.

1. Rigid Body

**2. Angular Position, Velocity, and Acceleration**

3. Rotational Motion with Constant Angular Acceleration

4. Angular and Linear Quantities

5. Rotational Kinetic Energy

6. Calculation of Moments of Inertia

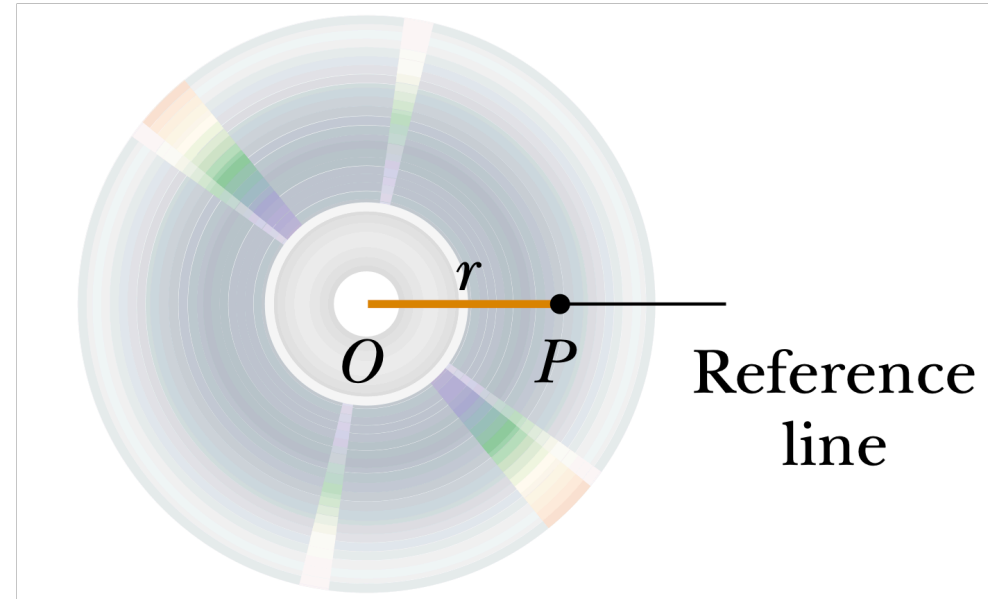
7. Torque

8. Relationship Between Torque and Angular Acceleration

9. Work, Power, and Energy in Rotational Motion

## 2.1 Angular Position

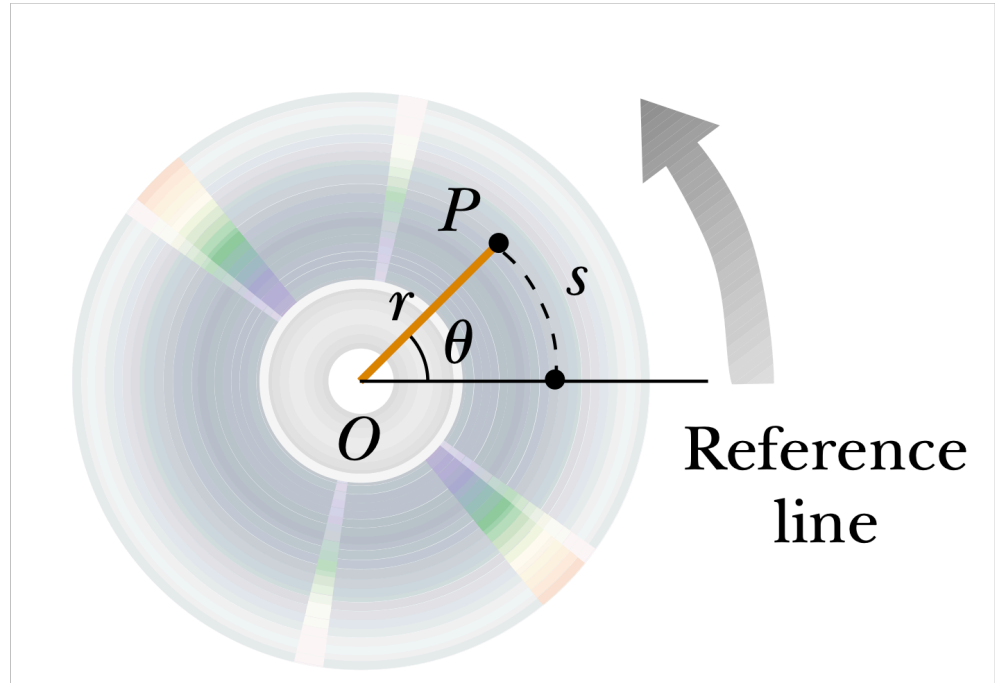
- The position of a point on a rigid body can be described using the angle  $\theta$  and a radius  $r$  from a reference point.



## 2.1 Angular Position

- The arc length  $s$  that the point travels is proportional to the angle  $\theta$  and the radius  $r$  by the equation:

$$s = r \theta$$



## 2.1 Angular Position

- Rearranging this equation gives us the angular position:

$$\theta = \frac{s}{r}$$

- Notice that theta has no unit, however, it is normally expressed in radians.
- 1 radians is equivalent to an arc length equal to the radius of the circle.
- Since the circumference of a circle is  $2\pi r$ , then one complete revolution is equal to  $\theta = 2\pi$  radians.

## 2.1 Angular Position

- To convert from degrees to radians, use the relation:

$$\theta(\text{rad}) = \frac{\pi}{180^\circ} \theta(\text{deg}) \approx \frac{3.14}{180} \theta(\text{deg})$$

To convert from revolutions to radians, use the relation:

$$\theta(\text{rad}) = 2\pi(\text{Number of rev})$$

To find the change in angular position, we use:

$$\Delta\theta = \theta_f - \theta_i$$

where  $\theta_f$  and  $\theta_i$  are the final and initial angular positions, respectively.

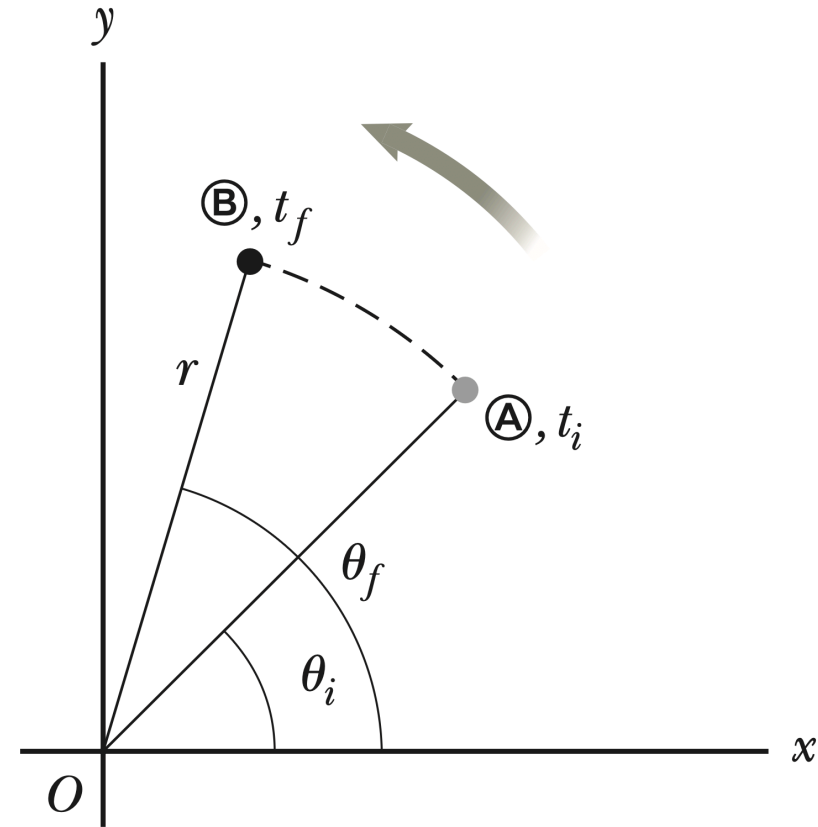
## 2.2 Angular speed

### Average Angular Speed

$$\bar{\omega} = \frac{\Delta\theta}{\Delta t}$$

### Instantaneous Angular Speed

$$\omega = \frac{d\theta}{dt}$$



## 2.3 Angular Acceleration

### Average Angular Acceleration

$$\bar{\alpha} = \frac{\Delta\omega}{\Delta t}$$

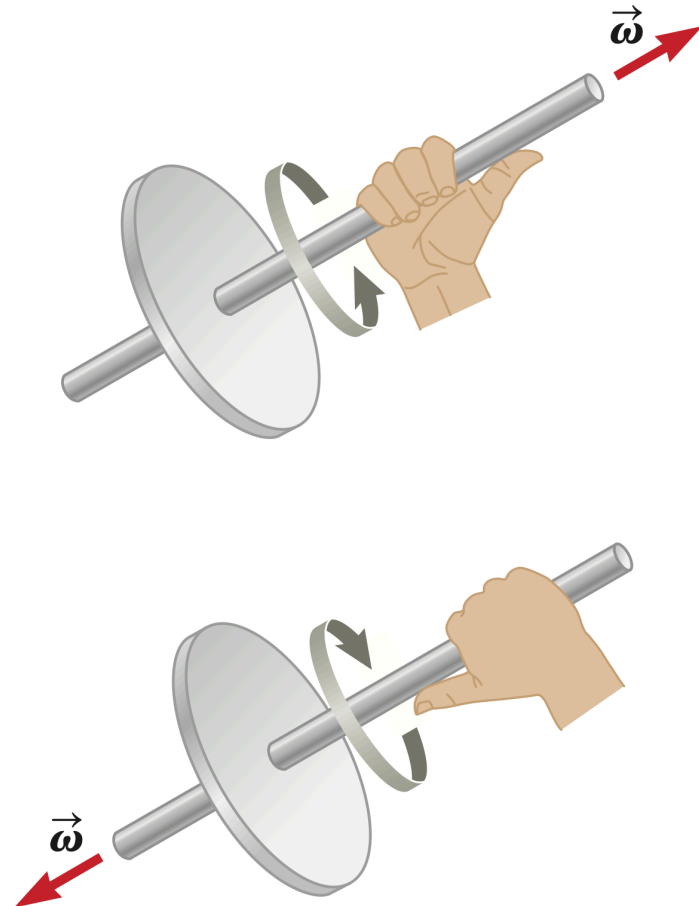
### Instantaneous Angular Acceleration

$$\alpha = \frac{d\omega}{dt}$$



## 2.4 Directions of Angular Velocity and Acceleration

- The direction of  $\vec{\omega}$  is given by the right-hand rule:
- If the fingers of the right hand curl in the direction of rotation, the thumb points in the direction of  $\vec{\omega}$ .
- Therefore, there are only two possible directions for  $\vec{\omega}$ : positive or negative along the axis of rotation.
- Similarly,  $\vec{\alpha}$  also can point in either direction along the axis of rotation, depending on whether the angular speed is increasing or decreasing.



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- 3. Rotational Motion with Constant Angular Acceleration**
4. Angular and Linear Quantities
5. Rotational Kinetic Energy
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8. Relationship Between Torque and Angular Acceleration
9. Work, Power, and Energy in Rotational Motion

### 3.1 Equations of Angular motion with constant acceleration

$$\omega_f = \omega_i + \alpha t$$

$$\theta_f = \theta_i + \omega_i t + \frac{1}{2} \alpha t^2$$

$$\theta_f = \theta_i + \frac{1}{2} (\omega_i + \omega_f) t$$

$$\omega_f^2 = \omega_i^2 + 2\alpha(\theta_f - \theta_i)$$

## 3.1 Equations of Angular motion with constant acceleration

### Example 3.1

A wheel rotates with a constant angular acceleration of  $3.5 \text{ rad/s}^2$ .

(A) If the angular speed of the wheel is  $2 \text{ rad/s}$  at  $t_i = 0$ , through what angular displacement does the wheel rotate in  $2 \text{ s}$ ?

$$\begin{aligned}\Delta\theta &= \theta_f - \theta_i = \omega_i t + \frac{1}{2}\alpha t^2 \\ &= (2 \text{ rad/s})(2\text{s}) + \frac{1}{2}(3.5 \text{ rad/s}^2)(2\text{s})^2 \\ &= 11 \text{ rad} = 11 \text{ rad} \frac{57.3^\circ}{1 \text{ rad}} \approx 630^\circ\end{aligned}$$

### 3.1 Equations of Angular motion with constant acceleration

(B) Through how many revolutions has the wheel turned during this time interval?

$$\Delta\theta = 11 \text{ rad} \frac{1 \text{ rev}}{2\pi \text{ rad}} \approx 1.75 \text{ rev}$$

(C) What is the angular speed of the wheel at  $t = 2 \text{ s}$ ?

$$\omega_f = \omega_i + \alpha t = 2 \text{ rad/s} + (3.5 \text{ rad/s}^2)(2\text{s}) = 9 \text{ rad/s}$$

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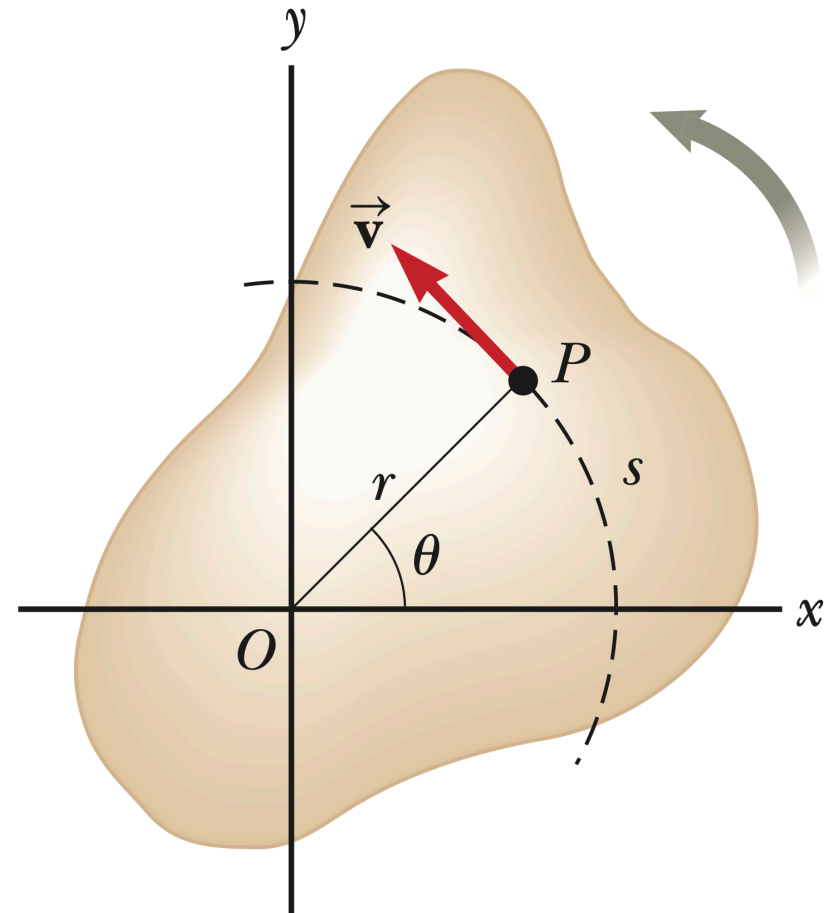
## 4.1 Velocity

- The linear velocity  $v$  of a point on a rotating rigid body is related to its angular velocity  $\omega$  by the equation:

$$v = \frac{ds}{dt} = r \frac{d\theta}{dt}$$

Therefore

$$v = r \omega$$



## 4.2 Acceleration

- The tangential acceleration  $a_t$  of a point on a rotating rigid body is related to its angular acceleration  $\alpha$  by the equation:
- Also, we can rewrite the centripetal acceleration  $a_c$ , to:

$$a_t = \frac{dv}{dt} = r \frac{d\omega}{dt}$$

Therefore,

$$a_t = r\alpha$$

$$a_c = \frac{v^2}{r} = r\omega^2$$



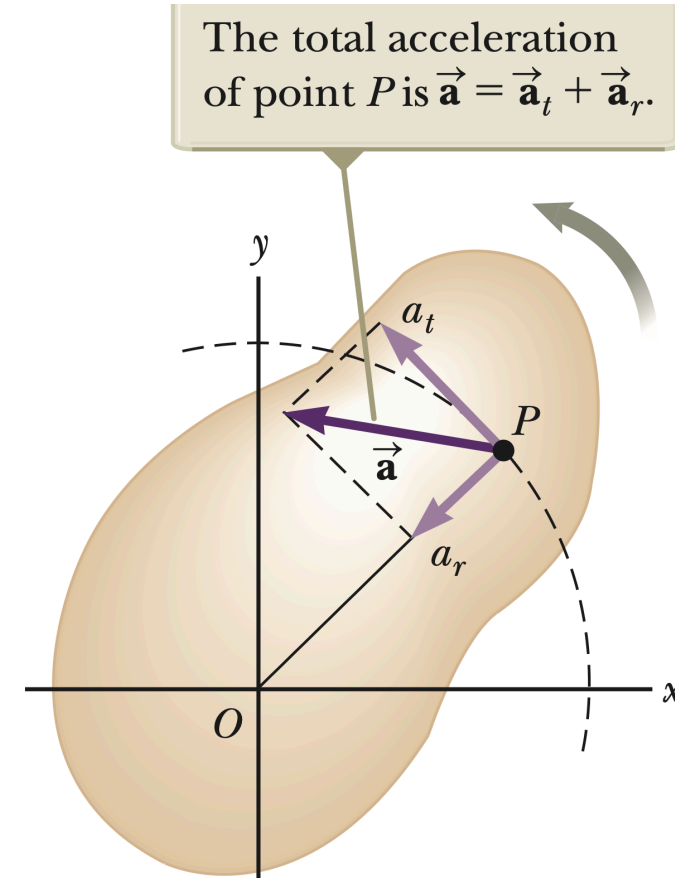
## 4.2 Acceleration

### Total acceleration

The total acceleration  $a_{\text{tot}}$  of a point on a rotating rigid body is the vector sum of its tangential acceleration  $a_t$  and centripetal acceleration  $a_c$ . Its magnitude is given by:

$$\begin{aligned} a_{\text{tot}} &= \sqrt{a_t^2 + a_c^2} \\ &= \sqrt{(r\alpha)^2 + (r\omega^2)^2} \end{aligned}$$

$$a_{\text{tot}} = r\sqrt{\alpha^2 + \omega^4}$$

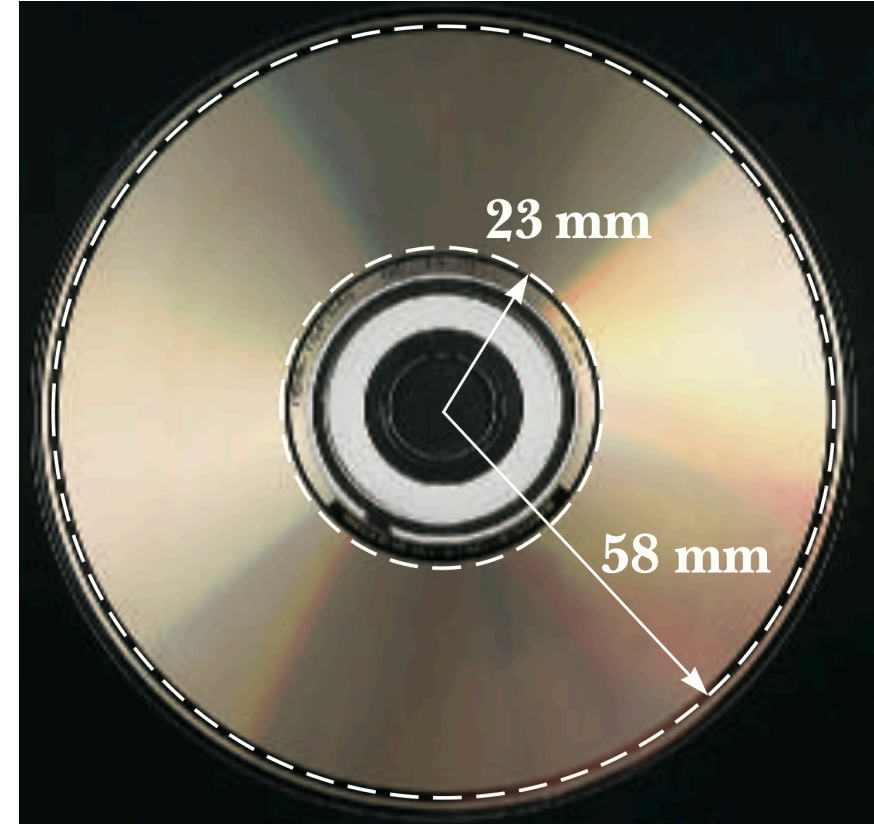


## 4.2 Acceleration

### Example 4.2

In a typical compact disc player, the constant speed of the surface at the point of the laser-lens system is 1.3 m/s.

(A) Find the angular speed of the disc in revolutions per minute when information is being read from the innermost first track ( $r = 23$  mm) and the outermost final track ( $r = 58$  mm).



## 4.2 Acceleration

### Solution 4.2

- For the innermost track:

$$\begin{aligned}\omega_i &= \frac{v}{r_i} = \frac{1.3 \text{ m/s}}{23 \times 10^{-3} \text{ m}} \\ &= 57 \text{ rad/s} \\ &= 57 \text{ rad/s} \left( \frac{1 \text{ rev}}{2\pi \text{ rad}} \right) \left( \frac{60 \text{ s}}{1 \text{ min}} \right) \\ &= 540 \text{ rev/min}\end{aligned}$$

- For the outermost track:

$$\begin{aligned}\omega_f &= \frac{v}{r_f} = \frac{1.3 \text{ m/s}}{58 \times 10^{-3} \text{ m}} \\ &= 22 \text{ rad/s} \\ &= 210 \text{ rev/min}\end{aligned}$$

## 4.2 Acceleration

(B) The maximum playing time of a standard CD is 74 min and 33 s. How many revolutions does the disc make during that time?

$$\begin{aligned}\Delta\theta &= \frac{1}{2}(\omega_i + \omega_f)t \\ &= \frac{1}{2}(57 \text{ rad/s} + 22 \text{ rad/s})(74 \text{ min} * 60 \text{ s/min} + 33 \text{ s}) \\ &= 1.8 \times 10^5 \text{ rad} \\ &= 2.8 \times 10^4 \text{ rev}\end{aligned}$$

## 4.2 Acceleration

(C) What total length of track moves past the objective lens during this time?

$$\begin{aligned}x &= vt = (1.3 \text{ m/s})(74 \text{ min} * 60 \text{ s/min} + 33 \text{ s}) \\&= 5.8 \times 10^3 \text{ m} = 5.8 \text{ km}\end{aligned}$$

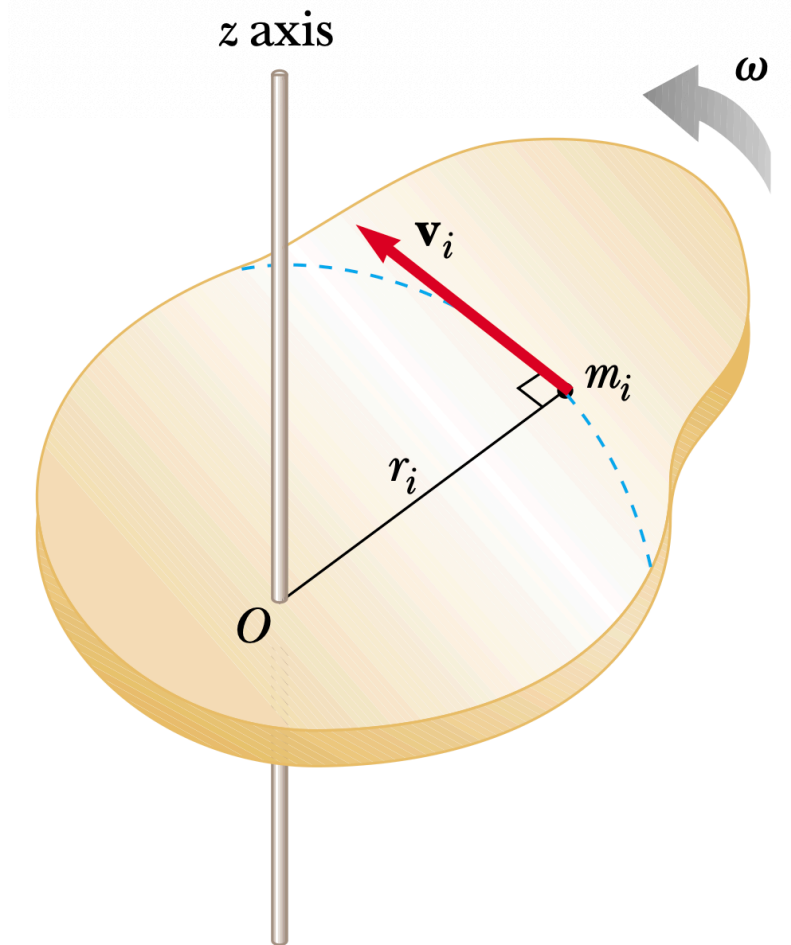
## 4.2 Acceleration

(D) What is the angular acceleration of the CD over the 4 473s time interval?  
Assume that  $\alpha$  is constant.

$$\begin{aligned}\alpha &= \frac{\omega_f - \omega_i}{t} \\ &= \frac{22 \text{ rad/s} - 57 \text{ rad/s}}{4473 \text{ s}} \\ &= -7.8 \times 10^{-3} \text{ rad/s}^2\end{aligned}$$

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## 5.1 Definition and Derivation



- A rigid object rotating around an axis with angular velocity  $\omega$ , the kinetic energy  $K_i$  of a tiny mass  $m_i$  at a distance  $r_i$  from the axis of rotation is given by:

$$\begin{aligned} K_i &= \frac{1}{2} m_i v_i^2 \\ &= \frac{1}{2} m_i (r_i \omega)^2 \\ &= \frac{1}{2} m_i r_i^2 \omega^2 \\ &= \frac{1}{2} (m_i r_i^2) \omega^2 \end{aligned}$$



## 5.1 Definition and Derivation

- The total rotational kinetic energy  $K_R$  of the rigid object is the sum of the kinetic energies of all its tiny masses:

$$\begin{aligned} K_R &= \sum_i K_i \\ &= \sum_i \frac{1}{2} (m_i r_i^2) \omega^2 \\ &= \frac{1}{2} \left( \sum_i m_i r_i^2 \right) \omega^2 \end{aligned}$$

$$K_R = \frac{1}{2} I \omega^2$$

where

$$I = \sum_i m_i r_i^2$$

Is the **moment of inertia** of the rigid object about the axis of rotation.

## 5.1 Definition and Derivation

### Example 5.3

Consider an oxygen molecule ( $\text{O}_2$ ) rotating in the  $xy$  plane about the  $z$  axis. The rotation axis passes through the center of the molecule, perpendicular to its length. The mass of each oxygen atom is  $m = 2.66 \times 10^{-26}$  kg, and at room temperature the average separation between the two atoms is  $d = 1.21 \times 10^{-10}$  m.

(A) Calculate the moment of inertia of the molecule about the  $z$  axis.

## 5.1 Definition and Derivation

### Solution 5.3

$$\begin{aligned} I &= \sum_i m_i r_i^2 = m \left( \frac{d}{2} \right)^2 + m \left( \frac{d}{2} \right)^2 \\ &= 2m \left( \frac{d}{2} \right)^2 = \frac{1}{2} m d^2 \\ &= \frac{1}{2} (2.66 \times 10^{-26} \text{ kg}) (1.21 \times 10^{-10} \text{ m})^2 \\ &= 1.95 \times 10^{-46} \text{ kg m}^2. \end{aligned}$$

## 5.1 Definition and Derivation

(B) If the angular speed of the molecule about the z-axis is  $\omega = 4.6 \times 10^{12}$  rad/s, what is its rotational kinetic energy?

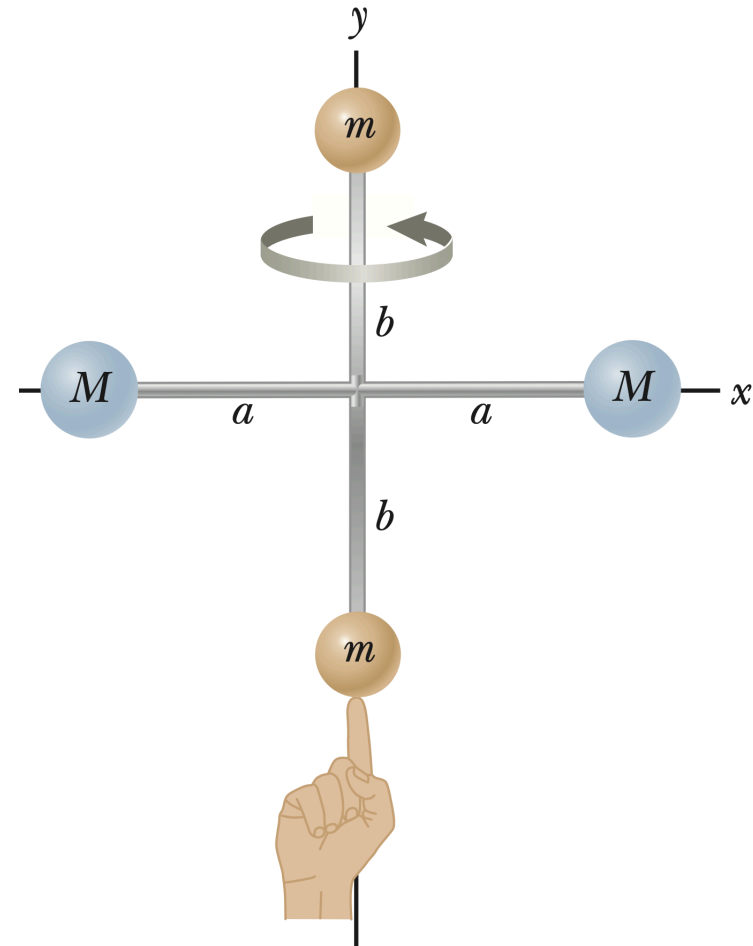
$$\begin{aligned} K_R &= \frac{1}{2} I \omega^2 \\ &= \frac{1}{2} (1.95 \times 10^{-46} \text{ kg m}^2) (4.6 \times 10^{12} \text{ rad/s})^2 \\ &= 2.1 \times 10^{-21} \text{ J.} \end{aligned}$$

## 5.1 Definition and Derivation

### Example 5.4

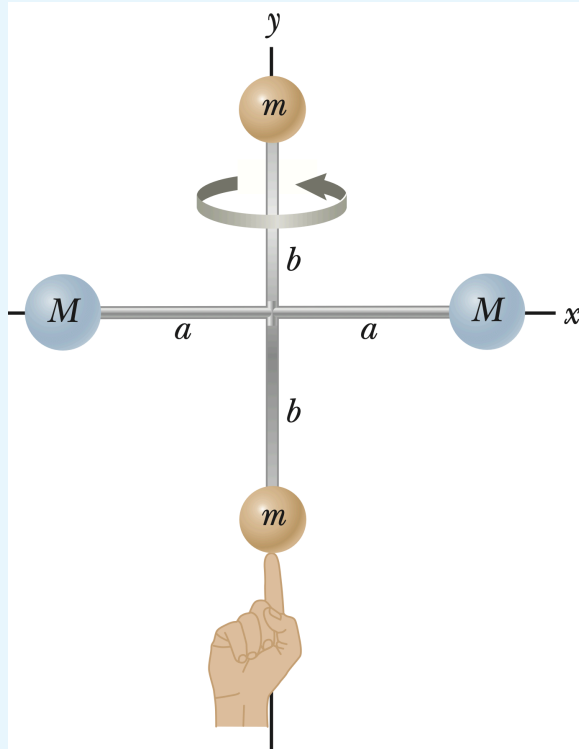
Four tiny spheres are fastened to the ends of two rods of negligible mass lying in the  $xy$  plane. Assume that the radii of the spheres are small compared with the dimensions of the rods.

(A) If the system rotates about the  $y$  axis with an angular speed  $\omega$ , find the moment of inertia and the rotational kinetic energy about this axis.



## 5.1 Definition and Derivation

### Solution 5.4



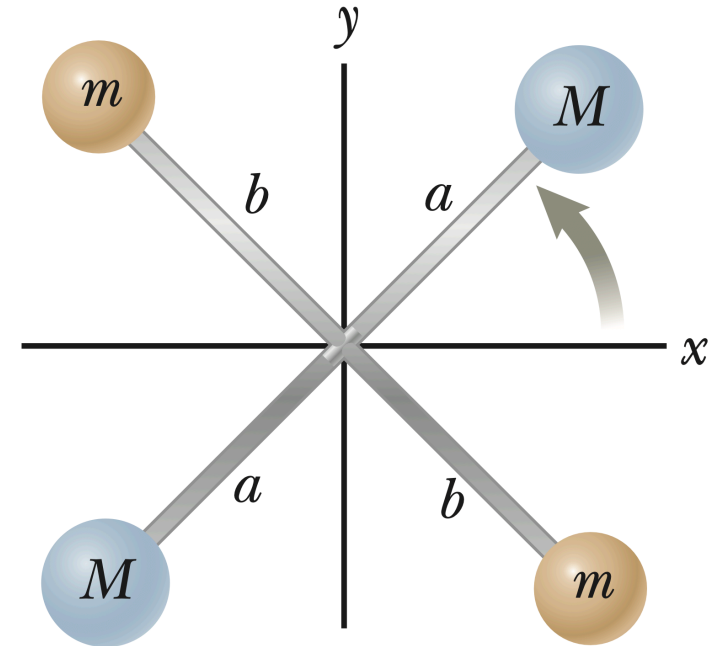
$$\begin{aligned} I_y &= m(0)^2 + Ma^2 + m(0)^2 + Ma^2 \\ &= 2Ma^2 \end{aligned}$$

Therefore, the rotational kinetic energy about the y-axis is:

$$\begin{aligned} K_R &= \frac{1}{2} I_y \omega^2 \\ &= \frac{1}{2} (2Ma^2) \omega^2 \\ &= Ma^2 \omega^2 \end{aligned}$$

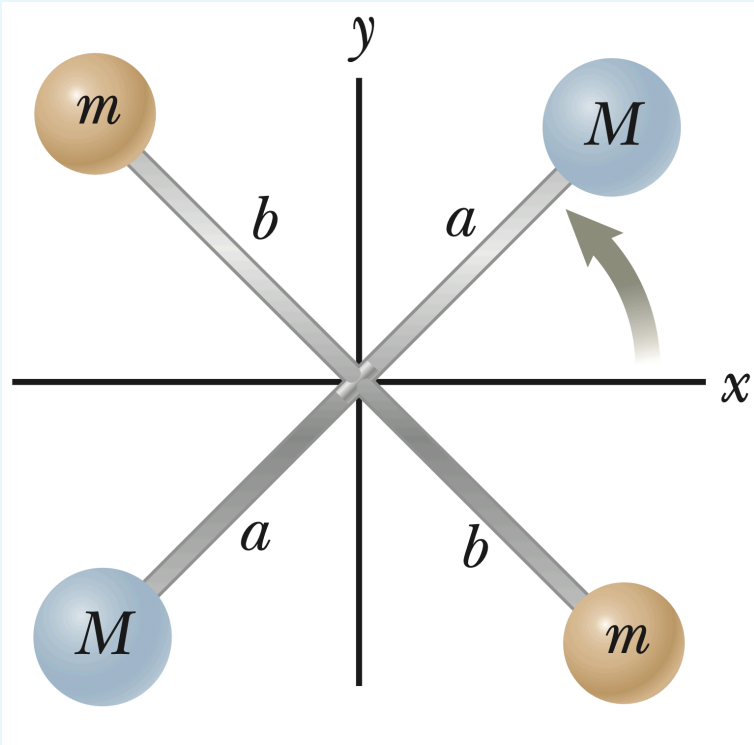
## 5.1 Definition and Derivation

(B) Suppose the system rotates in the  $xy$  plane about an axis (the  $z$  axis) through  $O$ . Calculate the moment of inertia and rotational kinetic energy about this axis.



## 5.1 Definition and Derivation

### Solution 5.4



$$\begin{aligned} I &= Ma^2 + mb^2 + Ma^2 + mb^2 \\ &= 2Ma^2 + 2mb^2 \end{aligned}$$

$$\begin{aligned} K_R &= \frac{1}{2} I \omega^2 \\ &= \frac{1}{2} (2Ma^2 + 2mb^2) \omega^2 \\ &= (Ma^2 + mb^2) \omega^2 \end{aligned}$$



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## 6.1 Integral Form of the Moment of Inertia

- For a continuous rigid body, the moment of inertia  $I$  about an axis of rotation is given by the integral:
- To evaluate this integral, we typically express  $dm$  in terms of a volume element  $dV$  and the density  $\rho = M/V$  of the material:

$$I = \int r^2 dm$$

where  $r$  is the perpendicular distance from the axis of rotation to the mass element  $dm$ .

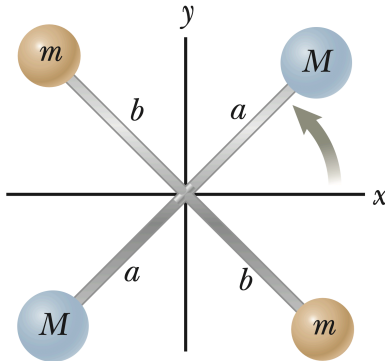
$$dm = \rho dV$$

$$I = \int \rho r^2 dV$$

## 6.2 Four Types of densities or Mass distributions

### 1. Point Masses

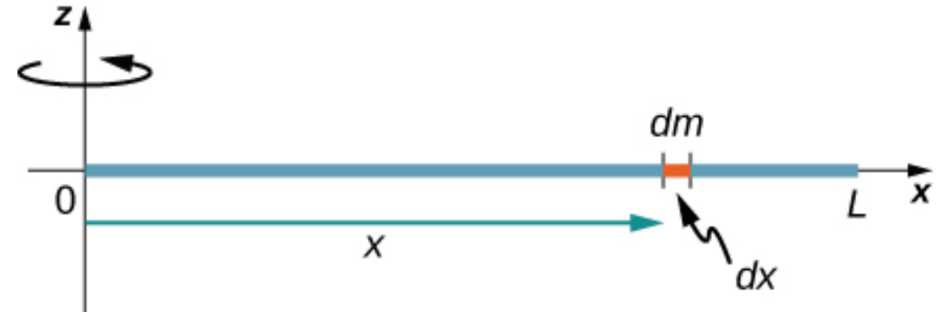
Discrete objects of individual point masses.



$$I = \sum_i m_i r_i^2$$

### 2. Linear density

For one-dimensional objects like rods.



$$\lambda = \frac{M}{L} = \rho A, \quad A: \text{cross-sectional area}$$

$$dm = \lambda dx$$

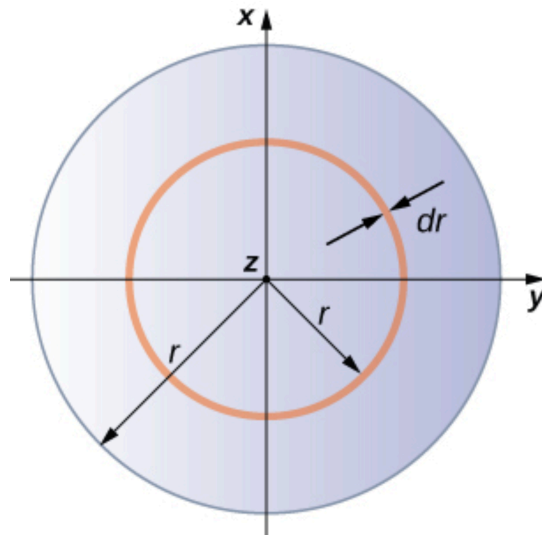
## 6.2 Four Types of densities or Mass distributions

### 3. Surface density

For 2D objects  
like plates.

$$\sigma = \frac{M}{A} = \rho t,$$

t: thickness

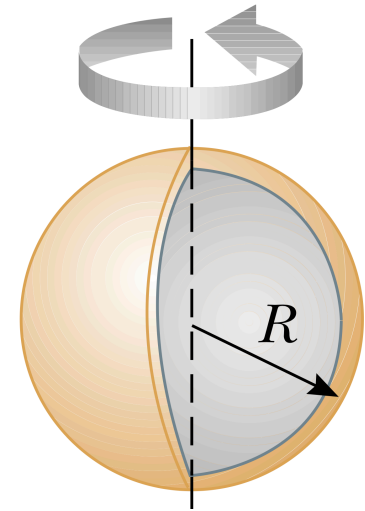


$$dm = \sigma dA$$

### 4. Volume density

For 3D objects like  
solid bodies.

$$\rho = \frac{M}{V}$$

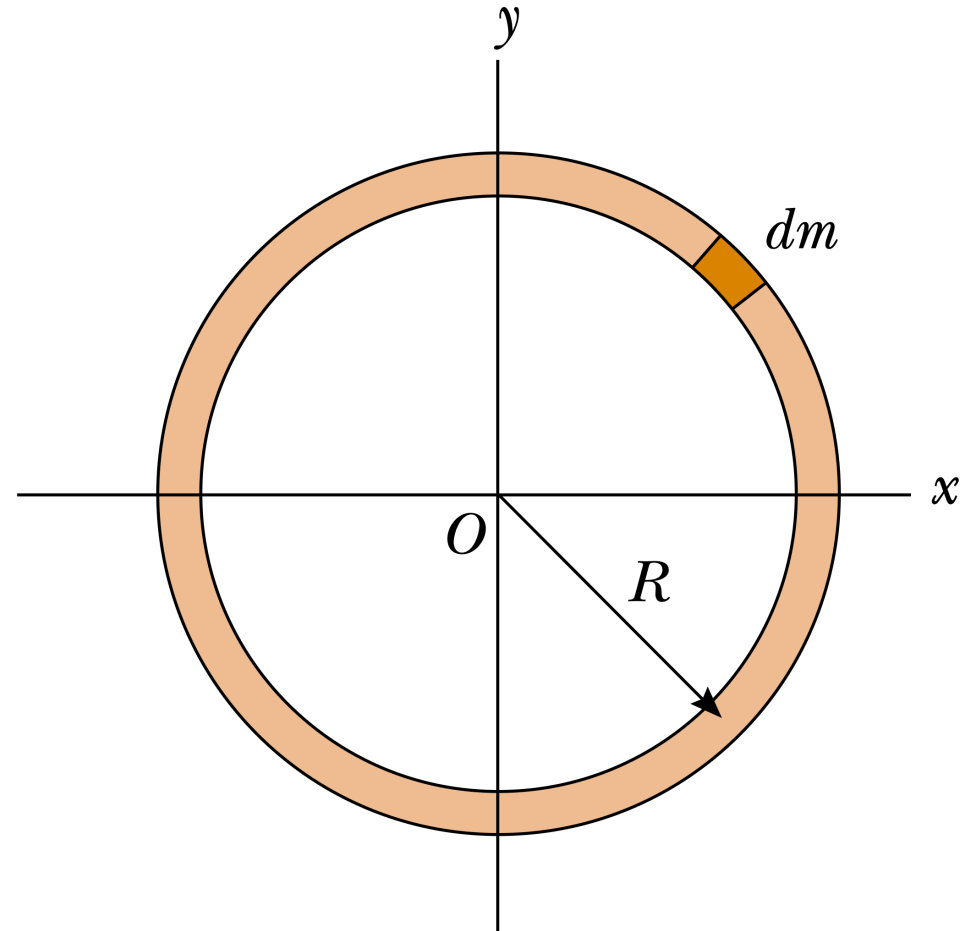


$$dm = \rho dV$$

## 6.3 Examples

### Example 6.5

Find the moment of inertia of a uniform thin hoop (طوق) of mass  $M$  and radius  $R$  about an axis perpendicular to the plane of the hoop and passing through its center.



## 6.3 Examples

### Solution 6.5

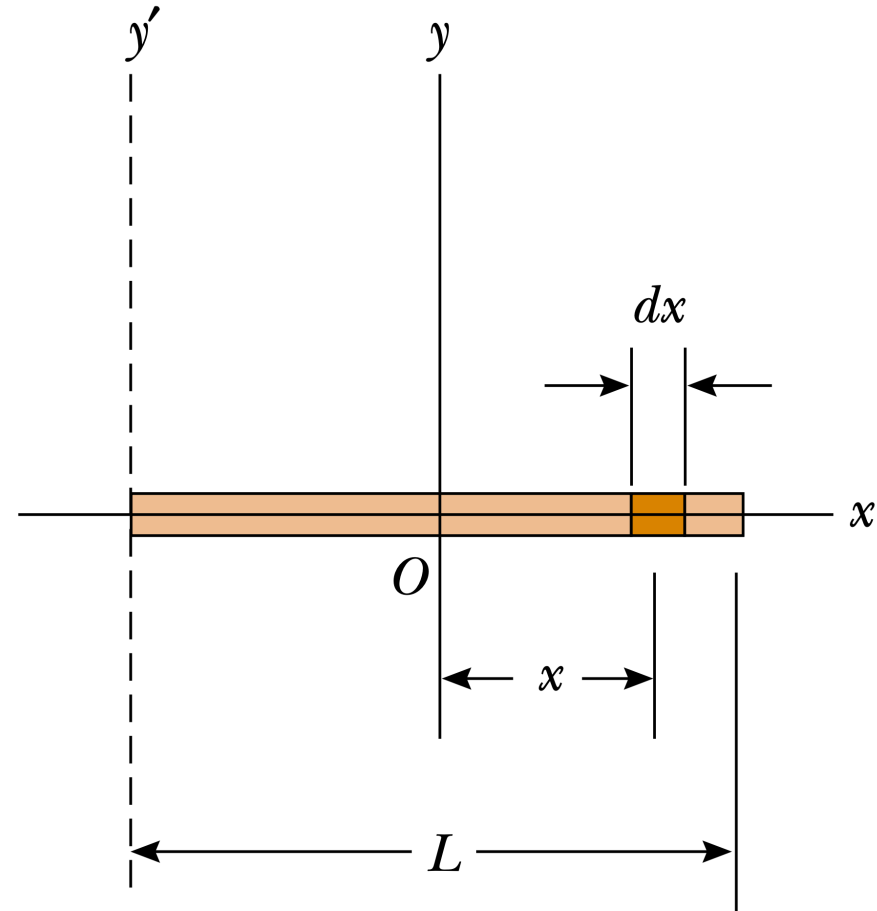
- We assume that the hoop is very thin, such that all its mass is concentrated at a distance  $R$  from the axis of rotation.
- Therefore, the moment of inertia can be calculated using:

$$I_z = \int r^2 dm = R^2 \int dm = R^2 M.$$

## 6.3 Examples

### Example 6.6

Calculate the moment of inertia of a uniform rigid rod of length  $L$  and mass  $M$  about an axis perpendicular to the rod (the  $y$  axis) and passing through its center of mass.



## 6.3 Examples

### Solution 6.6

$$dm = \lambda dx = \left( \frac{M}{L} \right) dx$$

Therefore,

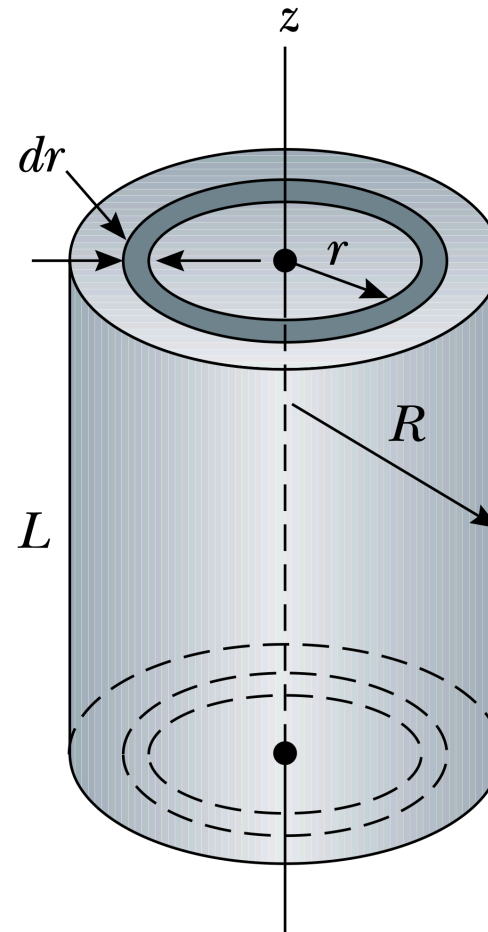
$$\begin{aligned} I_y &= \int r^2 dm = \int_{-\frac{L}{2}}^{\frac{L}{2}} x^2 \left( \frac{M}{L} \right) dx = \left( \frac{M}{L} \right) \left[ \frac{x^3}{3} \right]_{-\frac{L}{2}}^{\frac{L}{2}} \\ &= \left( \frac{M}{L} \right) \left( \frac{L^3}{24} + \frac{L^3}{24} \right) = \left( \frac{1}{12} \right) ML^2. \end{aligned}$$



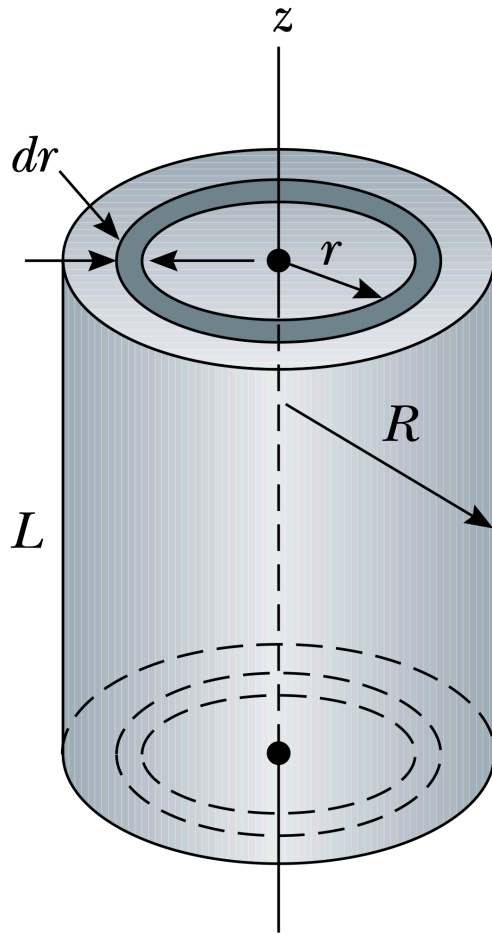
## 6.3 Examples

### Example 6.7

A uniform solid cylinder has a radius  $R$ , mass  $M$ , and length  $L$ . Calculate its moment of inertia about its central axis (the  $z$  axis).



## 6.3 Examples



### Solution 6.7

$$I_z = \int r^2 dm = \int r^2 \rho dV$$

It is convenient to divide the cylinder into thin cylindrical shells of radius  $r$ , thickness  $dr$  and length  $L$ . Therefore, the volume element of each shell is given by:

$$dV = L(2\pi r) dr$$

and the mass element is:

$$dm = \rho dV = 2\pi\rho Lr dr$$

## 6.3 Examples

Substituting into the integral for  $I_z$ , we get:

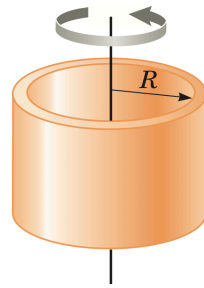
$$\begin{aligned} I_z &= \int_0^R r^2 (2\pi\rho Lr) \, dr = 2\pi\rho L \int_0^R r^3 \, dr \\ &= 2\pi\rho L \left[ \frac{r^4}{4} \right]_0^R = \frac{1}{2}\pi\rho LR^4. \end{aligned}$$

Substituting  $\rho = \frac{M}{V} = \frac{M}{\pi R^2 L}$  into the expression for  $I_z$ , we obtain:

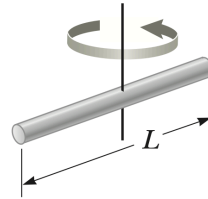
$$I_z = \frac{1}{2}\pi \left( \frac{M}{\pi R^2 L} \right) LR^4 = \left( \frac{1}{2} \right) MR^2$$

## 6.3 Examples

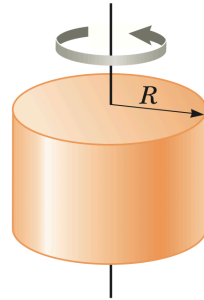
Hoop or thin  
cylindrical shell  
 $I_{\text{CM}} = MR^2$



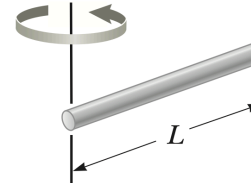
Long, thin rod  
with rotation axis  
through center  
 $I_{\text{CM}} = \frac{1}{12} ML^2$



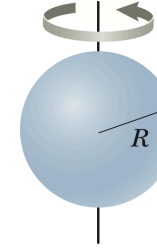
Solid cylinder  
or disk  
 $I_{\text{CM}} = \frac{1}{2} MR^2$



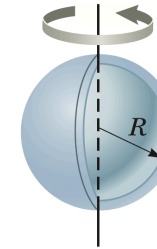
Long, thin rod  
with rotation axis  
through end  
 $I = \frac{1}{3} ML^2$



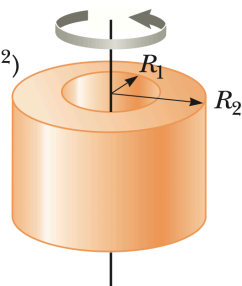
Solid sphere  
 $I_{\text{CM}} = \frac{2}{5} MR^2$



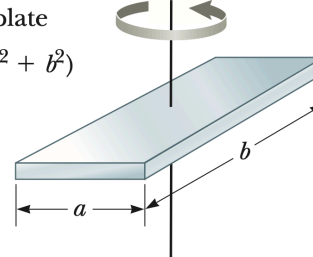
Thin spherical  
shell  
 $I_{\text{CM}} = \frac{2}{3} MR^2$



Hollow cylinder  
 $I_{\text{CM}} = \frac{1}{2} M(R_1^2 + R_2^2)$



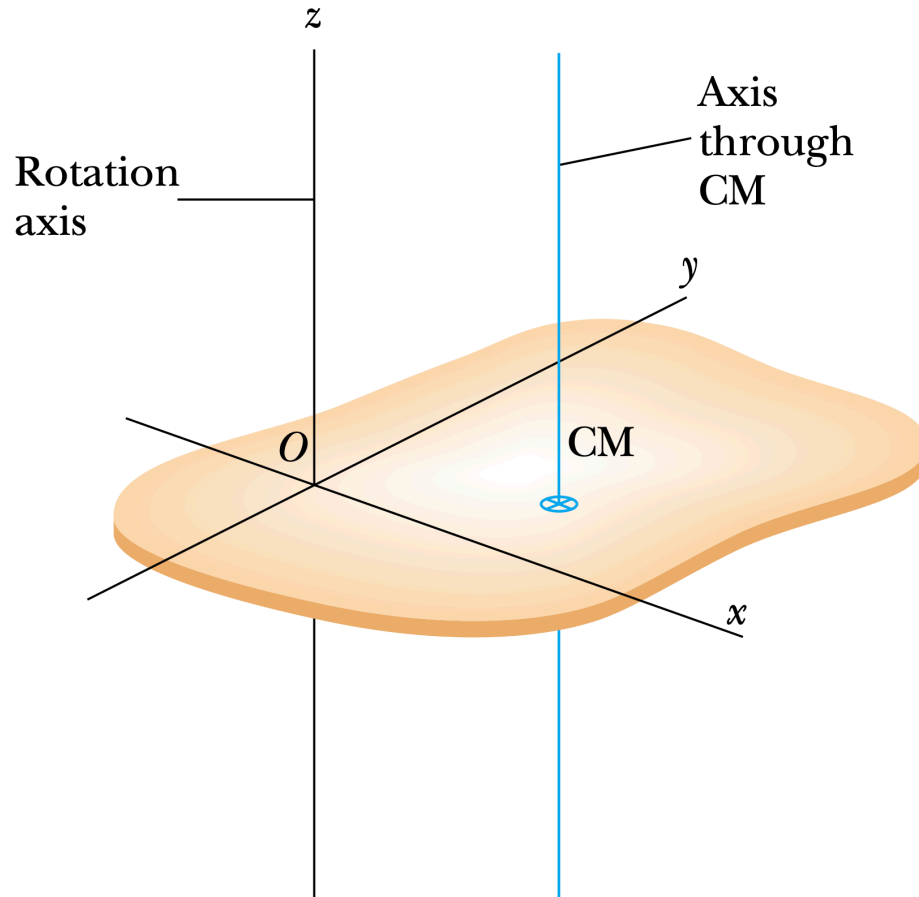
Rectangular plate  
 $I_{\text{CM}} = \frac{1}{12} M(a^2 + b^2)$



## 6.4 Center of Mass (CM)

- Each rigid body has a point called the center of mass (CM) where we can consider all the mass of the body to be concentrated for translational motion analysis to be simplified.
- For a symmetrical object with uniform density, the CM is located at its geometric center.
- When  $I$  is calculated about an axis passing through its CM, it is denoted as  $I_{CM}$ .
- Typically , calculating  $I_{CM}$  is easier than calculating  $I$  about any other axis.

## 6.5 Parallel-Axis Theorem



- The parallel-axis theorem states that the moment of inertia  $I$  of a rigid body about any axis parallel to and a distance  $D$  away from  $I_{\text{CM}}$  through its center of mass is given by:

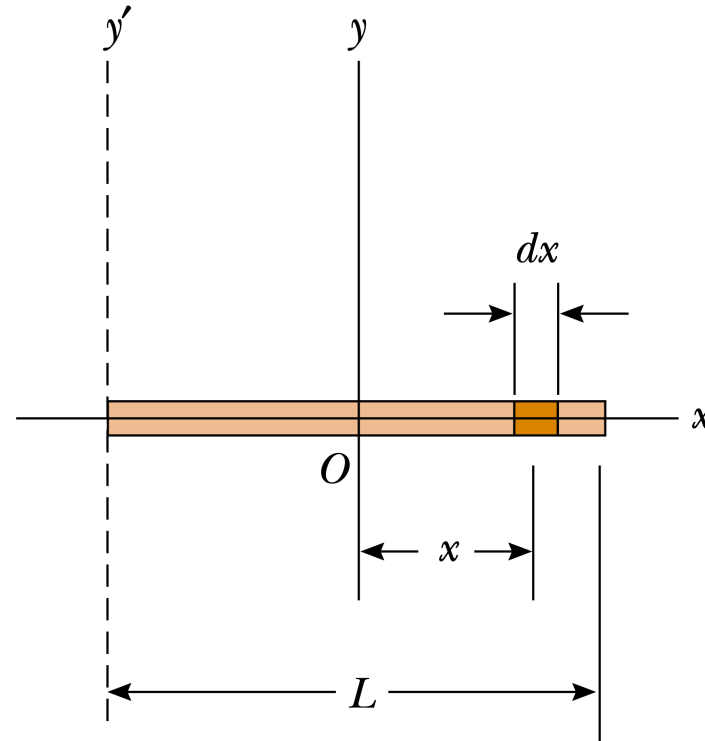
$$I = I_{\text{cm}} + MD^2$$

where  $I_{\text{cm}}$  is the moment of inertia about the center of mass axis, and  $M$  is the total mass of the body.

## 6.6 Example

### Example 6.8

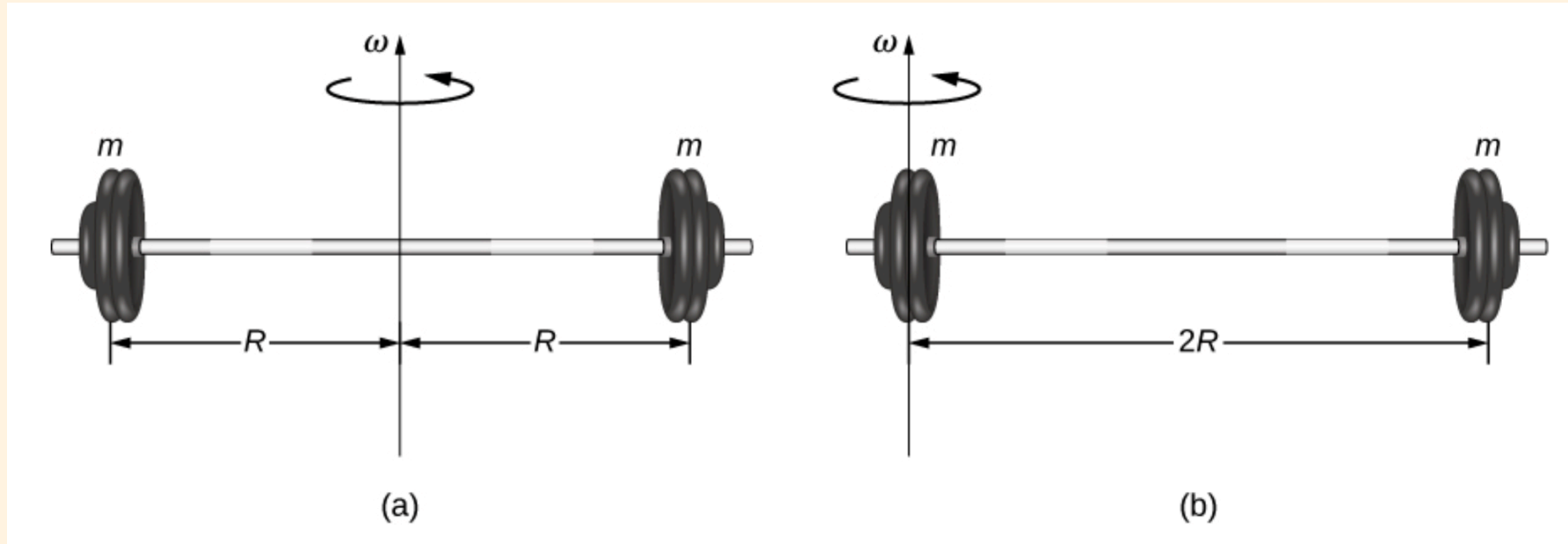
Consider once again the uniform rigid rod of mass  $M$  and length  $L$  shown in the Figure. Find the moment of inertia of the rod about an axis perpendicular to the rod through one end (the  $y'$  axis).



$$I = I_{\text{CM}} + MD^2 = \frac{1}{12}ML^2 + M\left(\frac{L}{2}\right)^2 = \frac{1}{3}ML^2$$

## 6.6 Example

### Quiz



Which one is easier to rotate?

The answer is (A) because the rotational axis passes through the center of mass.



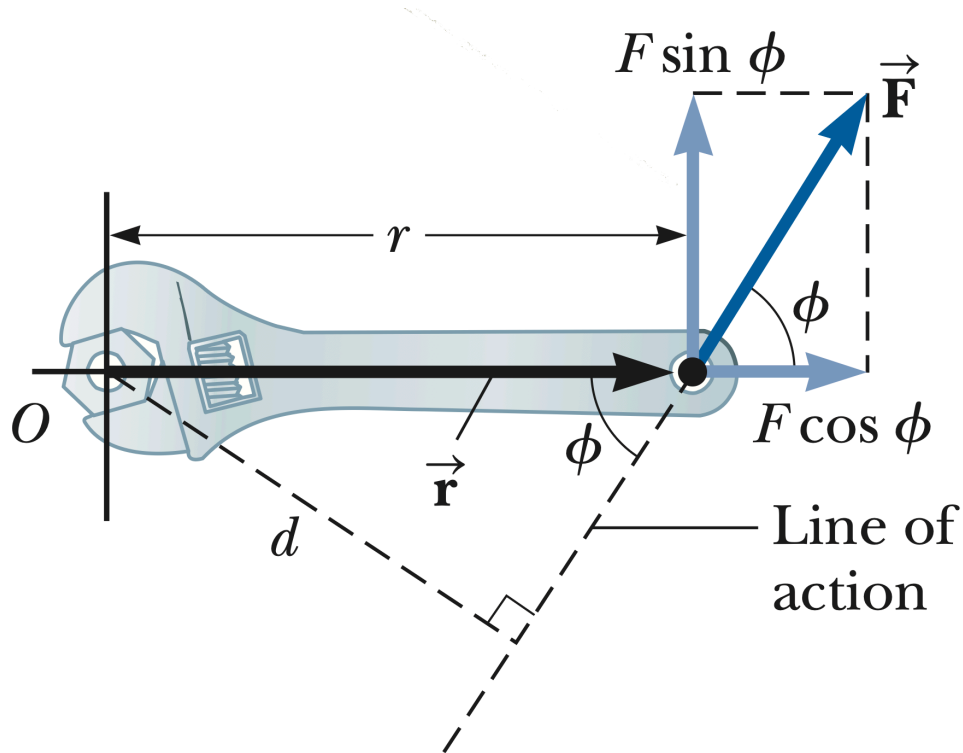
1. Rigid Body
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4. Angular and Linear Quantities
5. Rotational Kinetic Energy
6. Calculation of Moments of Inertia
- 7. Torque**
8. Relationship Between Torque and Angular Acceleration
9. Work, Power, and Energy in Rotational Motion

## 7.1 What is Torque?



**Torque**  $\tau$  is a measure of the strength of rotation or twisting effect around an axis.

## 7.1 What is Torque?

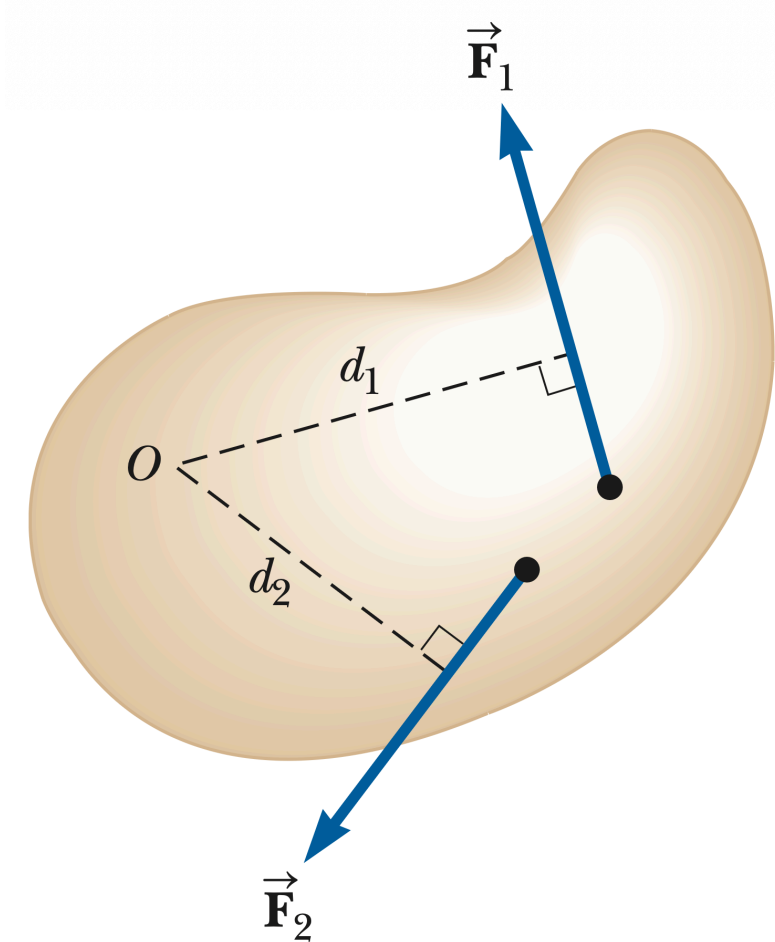


$$\tau = rF \sin \varphi = Fd$$

where the moment arm ( $d = r \sin \varphi$ ) is the perpendicular distance from the axis of rotation to the line of action of the force.

- The SI unit of torque is the newton-meter (**N·m**).
- Torque is a vector quantity, and its direction is given by the right-hand rule, similar to angular velocity  $\omega$ .

## 7.2 Multiple Forces Acting on a Rigid Body

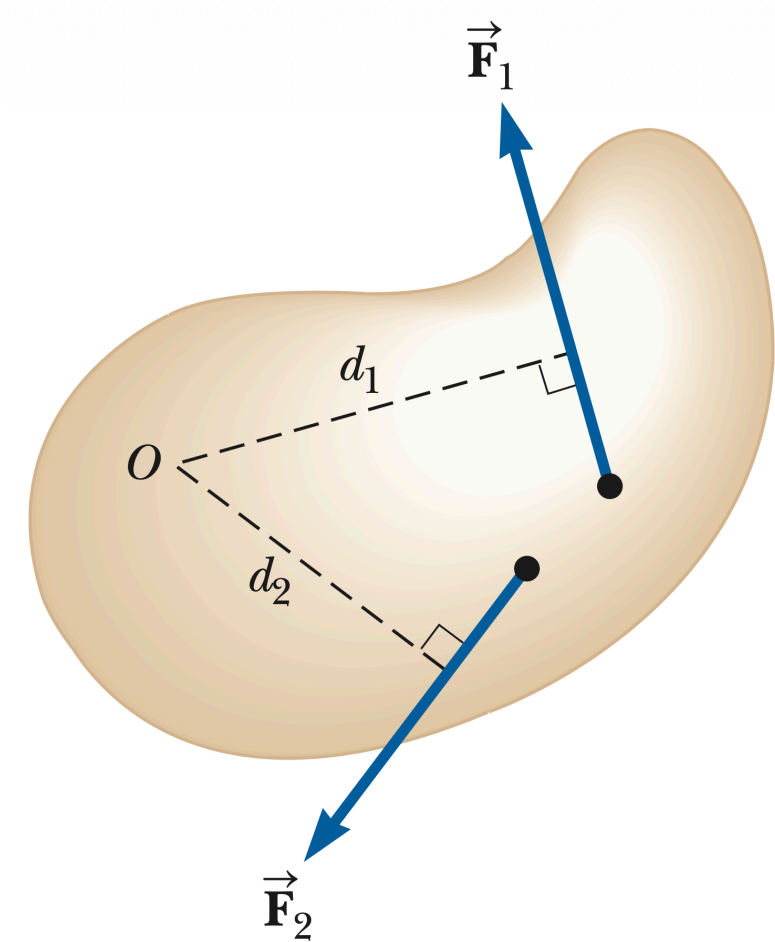


- When multiple forces act on a rigid body, the **net torque**  $\tau_{\text{net}}$  about an axis is the sum of the individual torques produced by each force:

$$\tau_{\text{net}} = \sum_i \tau_i$$

- The **sign** of each torque depends on the direction of rotation it produces: **counterclockwise** torques are **positive**, and **clockwise** torques are **negative**.

## 7.2 Multiple Forces Acting on a Rigid Body



- Therefore, the net torque for the left figure is:

$$\tau_{\text{net}} = \tau_1 - \tau_2 = F_1 d_1 - F_2 d_2$$

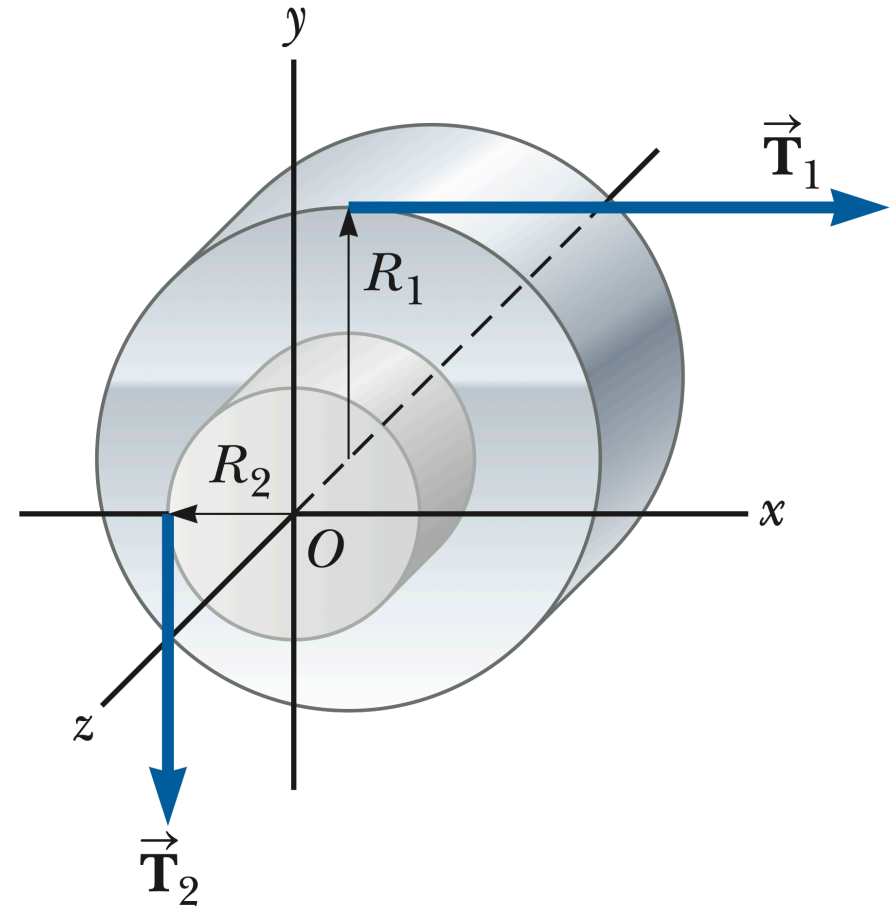
- Notice that  $F_2$  produces a clockwise rotation, so its torque is negative.

## 7.2 Multiple Forces Acting on a Rigid Body

### Example 7.9

A one-piece cylinder is free to rotate about the central axis along the  $z$ -axis. A rope wrapped around the drum, which has radius  $R_1 = 1$  m, exerts a force  $T_1 = 5$  N to the right on the cylinder. A rope wrapped around the core, which has radius  $R_2 = 0.5$  m, exerts a force  $T_2 = 15$  N downward on the cylinder.

What is the net torque acting on the cylinder about the rotation axis?



## 7.2 Multiple Forces Acting on a Rigid Body

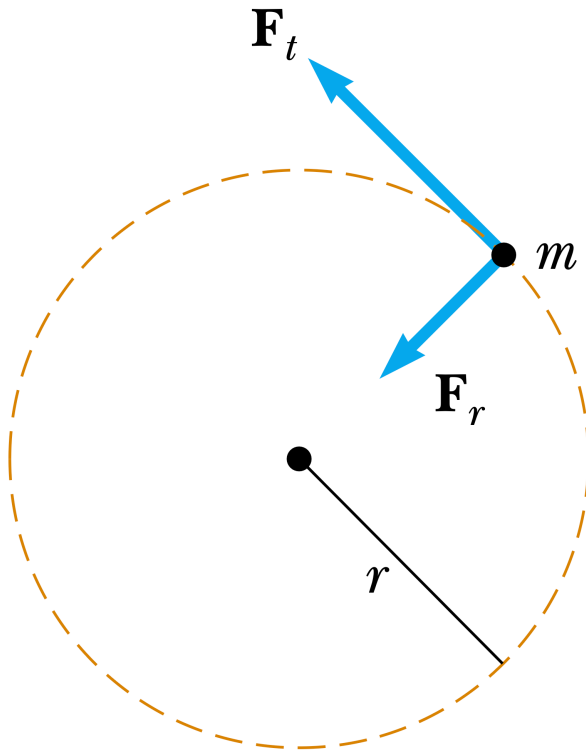
### Solution 7.9

$$t_{\text{net}} = -T_1 R_1 + T_2 R_2 = -(5 \text{ N})(1 \text{ m}) + (15 \text{ N})(0.5 \text{ m}) = 2.5 \text{ N.m}$$

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- 8. Relationship Between Torque and Angular Acceleration**
9. Work, Power, and Energy in Rotational Motion



## 8.1 Torque and Newton's Second Law for Rotation



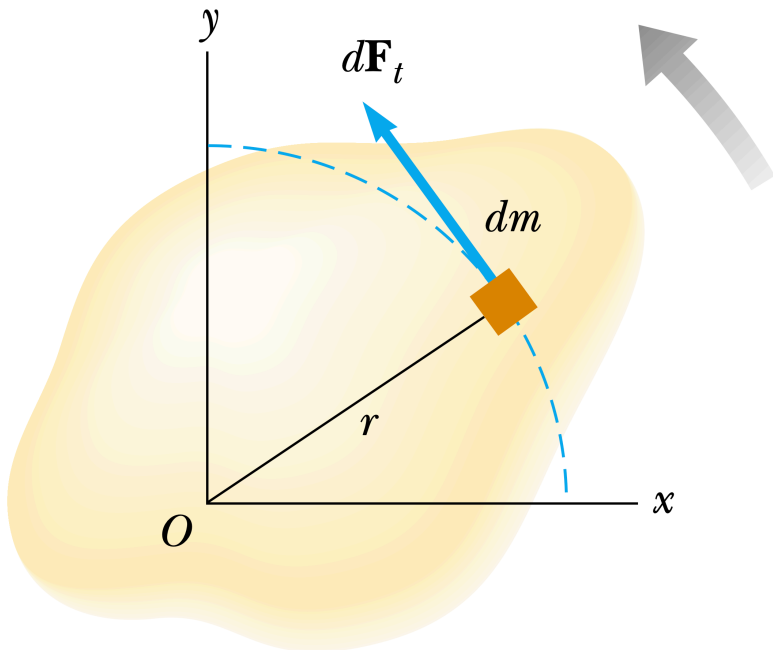
- When Forces act upon an object to rotate it about an axis, they produce a net torque  $\tau_{\text{net}}$  that causes an angular acceleration  $\alpha$  for the object.

$$\tau_{\text{net}} = F_t r = m a_t r = m(r\alpha)r = mr^2\alpha.$$

$$\tau_{\text{net}} = I\alpha$$

- The torque is proportional to its angular acceleration, and the proportionality constant is the moment of inertia.

## 8.2 Rigid Objects of Arbitrary Shape



- For a rigid object, we can divide it into tiny mass elements  $dm_i$  at distances  $r_i$  from the axis of rotation.

- Each mass element experiences a tangential force  $F_t$  that produces a torque  $\tau_i$  about the axis:

$$\begin{aligned}\tau_i &= r_i F_t = r_i (dm_i) a_t \\ &= r_i (dm_i) (r_i \alpha) = r_i^2 (dm_i) \alpha\end{aligned}$$

- The net torque on the object is

$$\tau_{\text{net}} = \sum_i r_i^2 (dm_i) \alpha = \alpha \sum_i r_i^2 (dm_i)$$

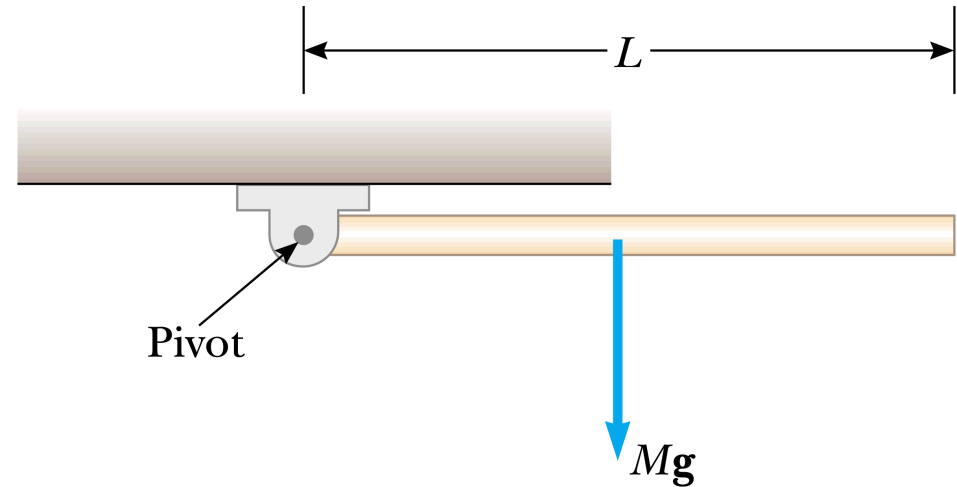
$$\tau_{\text{net}} = I \alpha$$

## 8.2 Rigid Objects of Arbitrary Shape

### Example 8.10

A uniform rod of length  $L$  and mass  $M$  is attached at one end to a frictionless pivot and is free to rotate about the pivot in the vertical plane. The rod is released from rest in the horizontal position.

What is the initial angular acceleration of the rod and the initial linear acceleration of its right end?



## 8.2 Rigid Objects of Arbitrary Shape

### Solution 8.10

- The magnitude of the torque about the pivot point due to the weight of the rod is:

$$\tau = Fd = (Mg) \left( \frac{L}{2} \right)$$

- Notice that we use  $d = L/2$  because it's the center of mass of the rod, therefore, we can assume that all its weight is concentrated at this point.

- Using  $\tau = I\alpha$ , we can find the angular acceleration  $\alpha$  of the rod:

$$\alpha = \frac{\tau}{I} = \frac{Mg \frac{L}{2}}{\frac{1}{3}ML^2} = \frac{3g}{2L}$$

- The linear acceleration  $a$  of the right end of the rod is related to its angular acceleration  $\alpha$  by the equation:

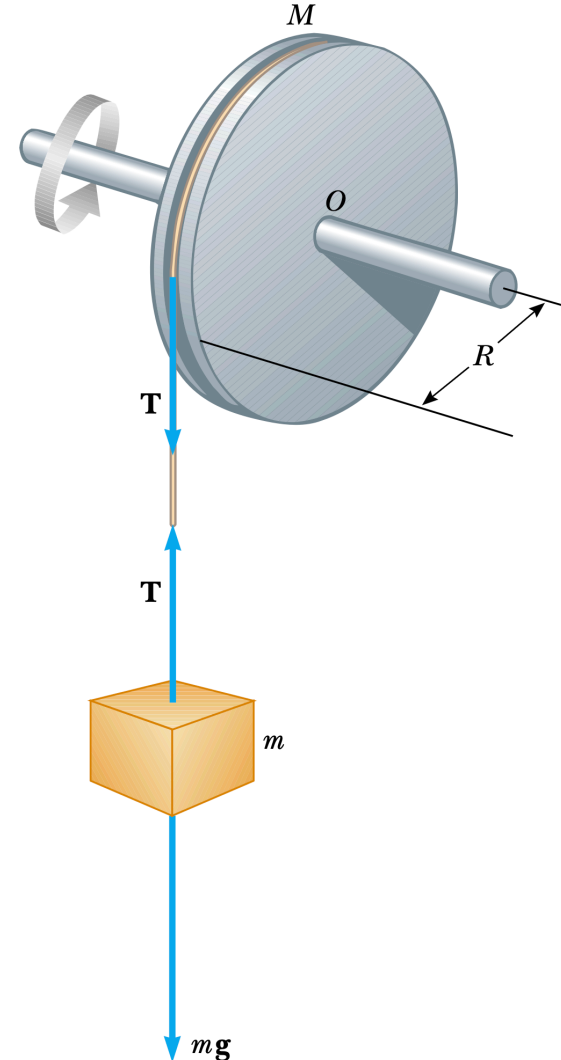
$$a = r\alpha = L \frac{3g}{2L} = \frac{3g}{2}$$

## 8.2 Rigid Objects of Arbitrary Shape

### Example 8.11

A wheel of radius  $R$ , mass  $M$ , and moment of inertia  $I$  is mounted on a frictionless horizontal axle. A light cord wrapped around the wheel supports an object of mass  $m$ .

Calculate (1) the angular acceleration of the wheel, (2) the linear acceleration of the object, and (3) the tension in the cord.



## 8.2 Rigid Objects of Arbitrary Shape

### Solution 8.11

- (1) For the wheel, the net torque about its axis is:

$$\tau_{\text{net}} = TR$$

- Using  $\tau_{\text{net}} = I\alpha$ , we have:

$$I\alpha = TR \Rightarrow \alpha = \frac{TR}{I}$$

- (2) The linear acceleration  $a$  of the falling object is related to the angular acceleration  $\alpha$  as:

$$a = R\alpha = \frac{TR^2}{I}$$

- Additionally,  $a$  can be found using Newton's second law:

$$\begin{aligned}\sum F &= mg - T = ma \\ \Rightarrow a &= \frac{mg - T}{m}\end{aligned}$$

- (3) Using the previous two equations for  $a$ , we get:

$$\frac{TR^2}{I} = \frac{mg - T}{m}$$

## 8.2 Rigid Objects of Arbitrary Shape

Rearranging this equation to solve for the tension  $T$  in the cord gives:

$$T = \frac{mg}{1 + mR^2/I}$$

We can also rewrite  $\alpha$  and  $a$  in a simpler form that does not require knowing  $T$ :

$$a = \frac{g}{1 + (I/mR^2)}$$

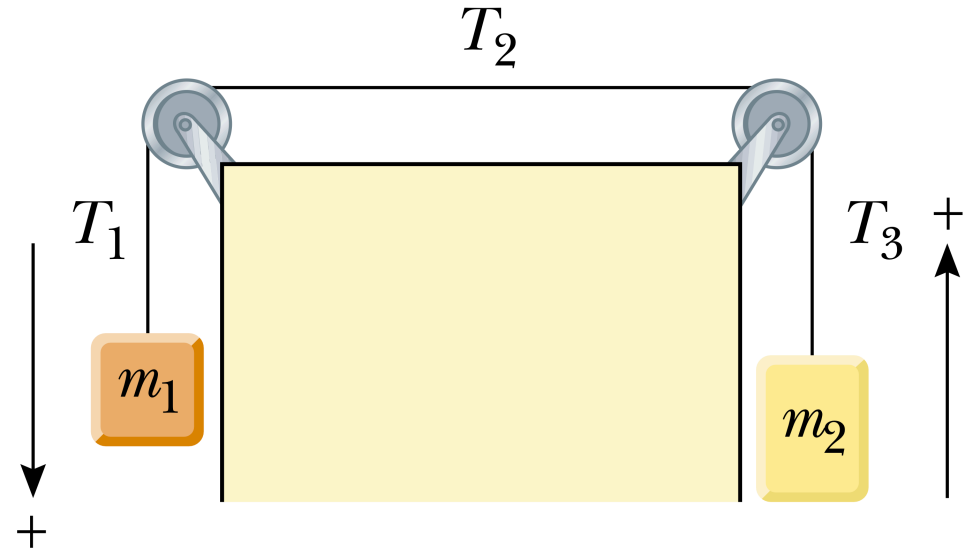
$$\alpha = \frac{a}{R} = \frac{g}{R + (I/mR)}$$

## 8.2 Rigid Objects of Arbitrary Shape

### Example 8.12

Two blocks having masses  $m_1$  and  $m_2$  are connected to each other by a light cord that passes over two identical frictionless pulleys, each having a moment of inertia  $I$  and radius  $R$ .

Find the acceleration of each block and the tensions  $T_1$ ,  $T_2$ , and  $T_3$  in the cord.

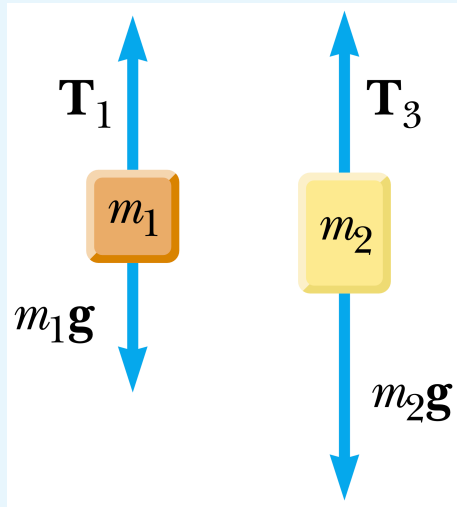




## 8.2 Rigid Objects of Arbitrary Shape

### Solution 8.12

- First, we set the **positive** direction to be counterclockwise.
- Second, we divide the two objects

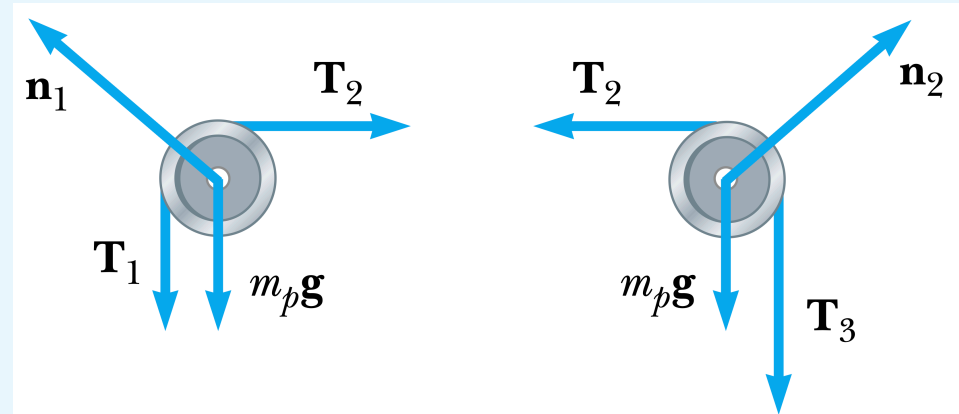


and apply Newton's second law.

$$m_1g - T_1 = m_1a \quad (1)$$

$$T_3 - m_2g = m_2a \quad (2)$$

- Next, we analyze the rotation of each pulley.



$$\tau_{\text{net}} = +T_1R - T_2R = I\alpha = I(a/R)$$

## 8.2 Rigid Objects of Arbitrary Shape

Therefore, the equations for the two pulleys are:

$$(T_1 - T_2)R = I(a/R) \quad (3)$$

$$(T_2 - T_3)R = I(a/R) \quad (4)$$

Solving these four equations simultaneously for  $a$ , we find:

$$a = \frac{(m_1 - m_2)g}{m_1 + m_2 + 2I/R^2}.$$

from equation (1), we find:

$$T_1 = m_1(g - a)$$

$$T_1 = 2m_1g \left( \frac{m_2 + I/R^2}{m_1 + m_2 + 2I/R^2} \right)$$

From equation (2), we find:

$$T_3 = m_2(g + a)$$

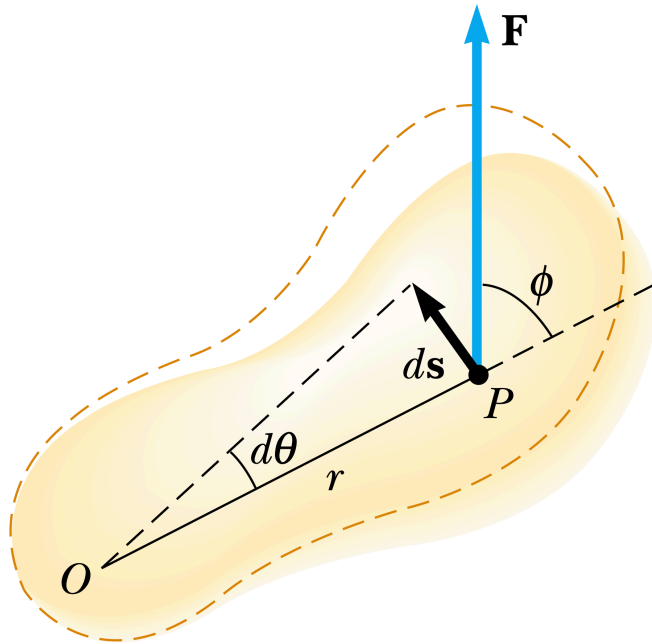
$$= 2m_2g \left( \frac{m_1 + I/R^2}{m_1 + m_2 + 2I/R^2} \right)$$

Finally, from equation (3), we find:

$$T_2 = \frac{2m_1m_2 + (m_1 + m_2)(I/R^2)}{m_1 + m_2 + 2I/R^2}g.$$

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## 9.1 Definitions



- Work**

$$dW = \vec{F} \cdot d\vec{s} = (F \sin \varphi)(r d\theta)$$

$$dW = \tau d\theta$$

- Power**

$$P = \frac{dW}{dt} = \tau\omega$$

- Work-Kinetic Energy Theorem**

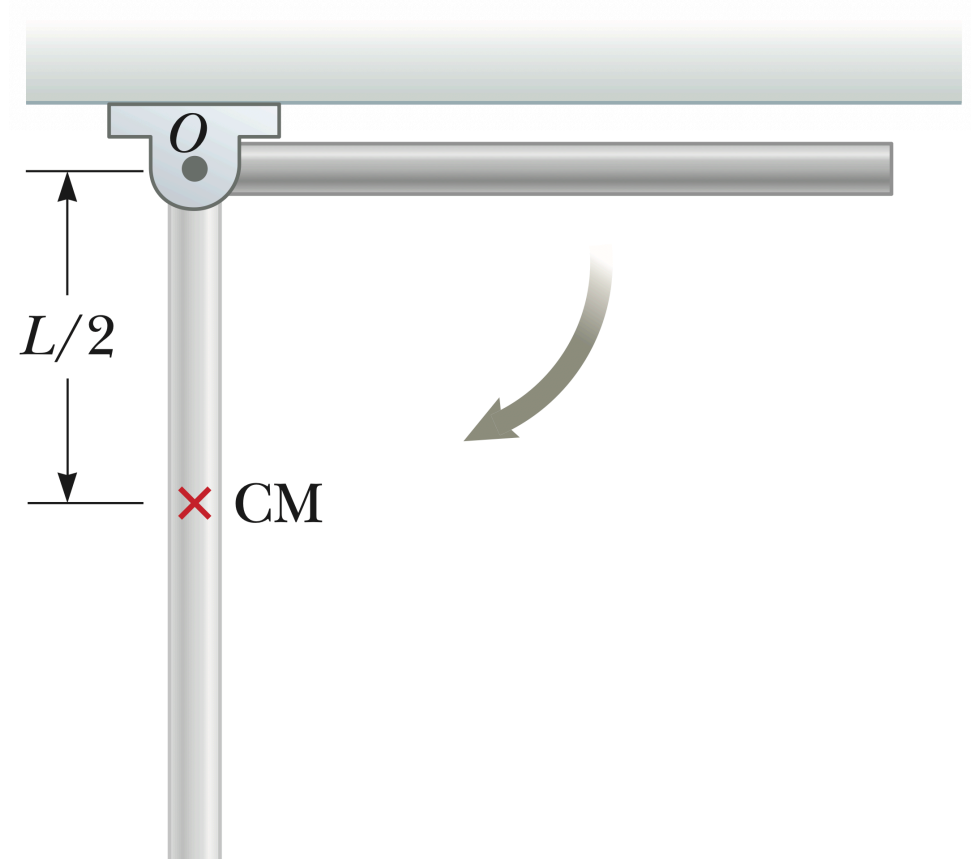
$$\begin{aligned} \sum W &= \frac{1}{2} I \omega_f^2 - \frac{1}{2} I \omega_i^2 \\ &= \Delta K \end{aligned}$$

## 9.1 Definitions

### Example 9.13

A uniform rod of length  $L$  and mass  $M$  is free to rotate on a frictionless pin passing through one end. The rod is released from rest in the horizontal position.

(A) What is its angular speed when it reaches its lowest position?



## 9.1 Definitions

### Solution 9.13

$$E_i = E_f$$

$$K_i + U_i = K_f + U_f$$

$$0 + Mgh = \frac{1}{2}I\omega_f^2 + 0$$

$$Mg\frac{L}{2} = \frac{1}{2}\left(\frac{1}{3}ML^2\right)\omega_f^2$$

Solving for  $\omega_f$ , we get:

$$\omega_f = \sqrt{\frac{3g}{L}}$$

## 9.1 Definitions

(B) Determine the tangential speed of the center of mass and the tangential speed of the lowest point on the rod when it is in the vertical position.

- The tangential speed of the center of mass is:

$$v_{\text{cm}} = r_{\text{cm}}\omega_f = \left(\frac{L}{2}\right)\sqrt{\frac{3g}{L}} = \sqrt{\frac{3gL}{4}}$$

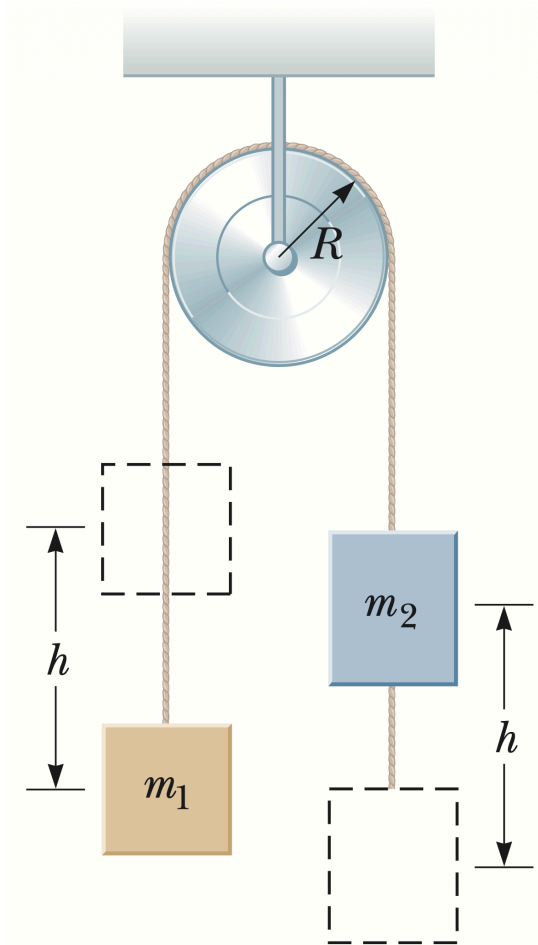
- The tangential speed of the lowest point on the rod is:

$$v_{\text{bottom}} = L\omega_f = L\sqrt{\frac{3g}{L}} = \sqrt{3gL}$$

## 9.1 Definitions

### Example 9.14

Consider two cylinders having different masses  $m_1$  and  $m_2$ , connected by a string passing over a pulley. The pulley has a radius  $R$  and moment of inertia  $I$  about its axis of rotation. The string does not slip on the pulley, and the system is released from rest. Find the linear speeds of the cylinders after cylinder 2 descends through a distance  $h$ , and the angular speed of the pulley at this time.





## 9.1 Definitions

### Solution 9.14

$$E_i = E_f$$

$$K_i + U_i = K_f + U_f$$

$$0 + m_2gh = \frac{1}{2}m_1v_f^2 + \frac{1}{2}m_2v_f^2 + \frac{1}{2}I\omega_f^2 + m_1gh$$

Replacing  $\omega_f$  with  $v_f/R$ , and solving for  $v_f$ , we get:

$$v_f = \sqrt{\frac{2gh(m_2 - m_1)}{m_1 + m_2 + I/R^2}}$$

and

$$\omega_f = \frac{v_f}{R} = \sqrt{\frac{2gh(m_2 - m_1)}{R^2(m_1 + m_2 + I/R^2)}}$$