

# QUANTUM MECHANICS: LECTURE 9

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## Abstract

This lecture discusses the quantum harmonic oscillator by the Ladder operator method

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## QUANTIZATION OF THE SHO HAMILTONIAN

From lecture (1) we have the classical Hamiltonian for the simple harmonic oscillator (SHO) :

$$H(p, x) = \frac{1}{2m}p^2 + \frac{1}{2}m\omega^2x^2 \quad (1)$$

Using the postulates of quantum mechanics discussed before, we obtain upon quantization - the Hamiltonian operator :

$$\hat{H} = \frac{1}{2m}\hat{p}^2 + \frac{1}{2}m\omega^2\hat{X}^2 \quad (2)$$

with:

$$[\hat{X}, \hat{P}] = i\hbar I \quad (3)$$

The Hilbert space of which  $\hat{X}$  and  $\hat{P}$  act on is

$$\mathcal{H}(0, +\infty; dx)$$

We now introduce the dimensionless Hamiltonian :

$$\hat{H}' = \frac{1}{2m\hbar\omega}\hat{p}^2 + \frac{1}{2}\frac{m\omega}{\hbar}\hat{X}^2 \quad (4)$$

This operator can be factorised and written in terms of 'creation' and 'inhalation' operators;  $a^\dagger$  and  $a$  respectively :

$$\hat{H}' = a^\dagger a + \frac{1}{2}I \quad (5)$$

with:

$$a = \sqrt{\frac{m\omega}{2\hbar}}\hat{X} + i\sqrt{\frac{1}{2m\omega\hbar}}\hat{P} \quad (6a)$$

$$a^\dagger = \sqrt{\frac{m\omega}{2\hbar}}\hat{X} - i\sqrt{\frac{1}{2m\omega\hbar}}\hat{P} \quad (6b)$$

These operators, along with  $\hat{H}'$ , satisfy a well-known commutation relations.

$$[a, a^\dagger] = I \quad (7a)$$

$$[a, H'] = a \quad (7b)$$

$$[a^\dagger, H'] = -a^\dagger \quad (7c)$$

*The operators  $a, a^\dagger$  and  $H'$  along with the commutator operation  $[\cdot, \cdot]$  satisfy the  $su(1,1)$  algebra.*

We also define the **number operator**  $N \equiv a^\dagger a$  that acts on the eigenstates  $|n\rangle$  resulting an eigenvalue of  $n$  :

$$N|n\rangle = n|n\rangle$$

as a result we may conclude that :

$$a|0\rangle = 0 \quad (8)$$

acting on the 'ground state' by the inhalation operator, kills it . Moreover :

$$a|n\rangle = \sqrt{n}|n-1\rangle \quad (9)$$

$$a^\dagger|n\rangle = \sqrt{(n+1)}|n+1\rangle \quad (10)$$

Hence, The Hamiltonian acting on these states will result ( the energy eigenvalue) :

$$\hat{H}|n\rangle = \hbar\omega\left(n + \frac{1}{2}\right)|n\rangle \quad (11)$$

Implying that the 'number states' are the excitation states for the quantum harmonic oscillator. The creation and inhalation operators excite or deceit it, and it has a descrete energy spectrum of :

$$E_n = \hbar\omega\left(n + \frac{1}{2}\right) \quad (12)$$

Even in the ground state, the quantum harmonic oscillator has a non-vanishing energy. This is a direct result for the uncertainty principle in time and energy.

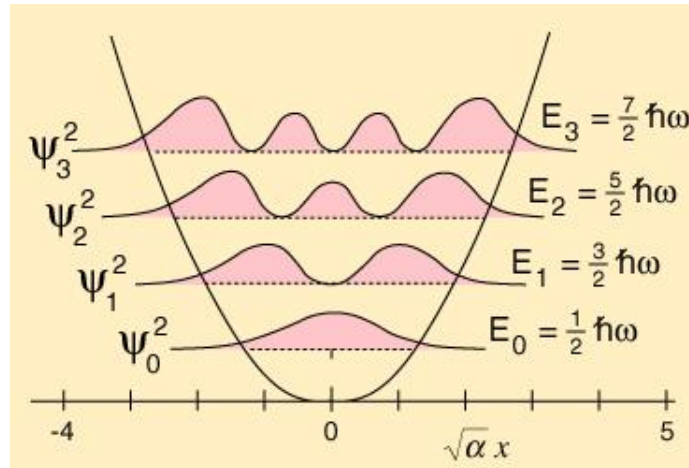


Figure 1: Energy-levels and wavefunctions of the quantum harmonic oscillator

#### REFERENCES

- [1] James Binney and David Skinner. *The physics of quantum mechanics*. Oxford University Press, 2013.
- [2] Kurt Gottfried and Tung-Mow Yan. *Quantum mechanics: fundamentals*. Springer Science & Business Media, 2013.
- [3] Sadri Hassani. *Mathematical physics: a modern introduction to its foundations*. Springer Science & Business Media, 2013.
- [4] Lev Davidovich Landau, Evgenii Mikhailovich Lifshitz, JB Sykes, John Stewart Bell, and ME Rose. Quantum mechanics, non-relativistic theory. *Physics Today*, 11:56, 1958.

- [5] John Von Neumann. *Mathematical foundations of quantum mechanics*. Number 2. Princeton university press, 1955.
- [6] Cohen Claude Tannoudji, Diu Bernard, and Laloë Franck. *Mécanique quantique*. tome i. 1973.
- [7] Angus Ellis Taylor and David C Lay. *Introduction to functional analysis*, volume 2. Wiley New York, 1958.