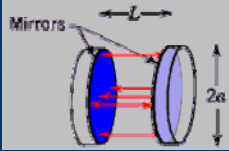
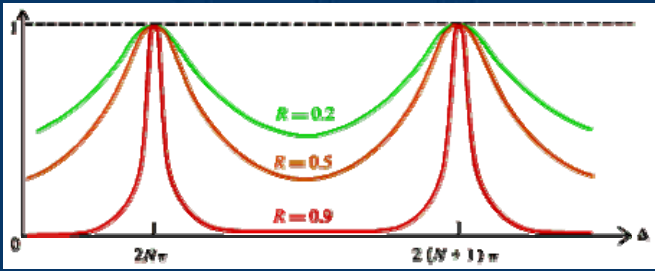


LASER RESONATOR



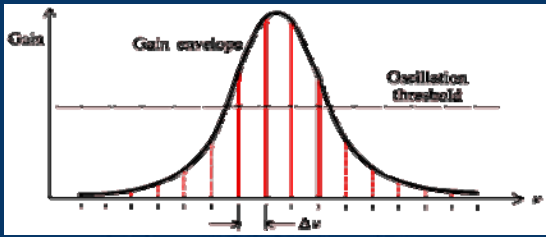
Mirrors $\leftarrow L \rightarrow$
 $2a$



$R = 0.2$
 $R = 0.5$
 $R = 0.9$

$2N\pi$ $2(N+1)\pi$

The spectral response of the cavity, where R is cavity reflectivity

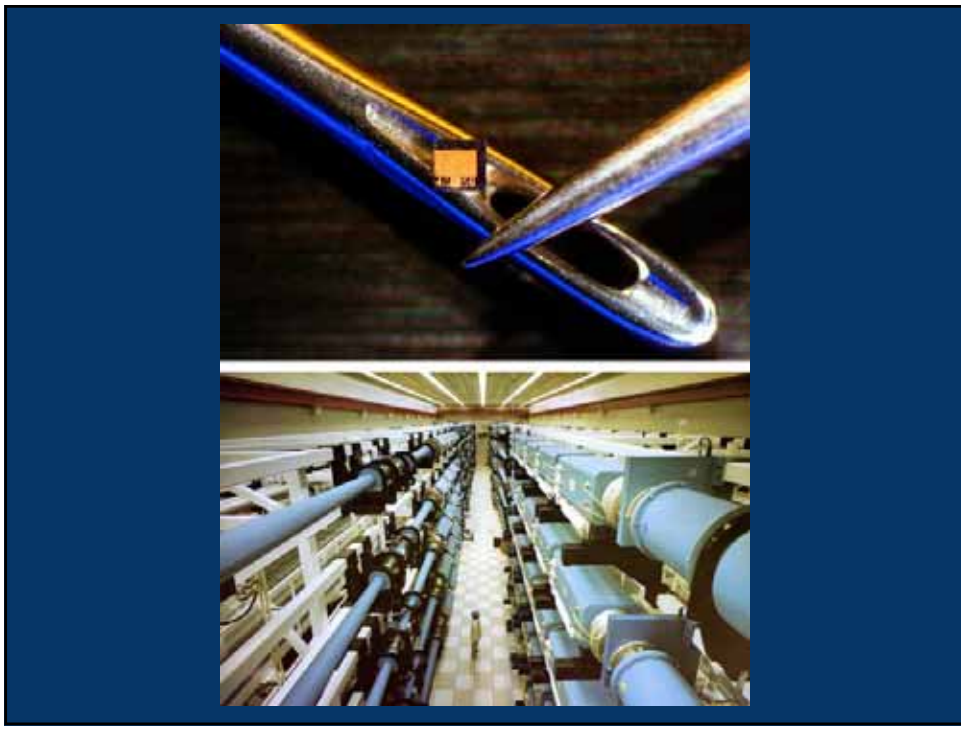
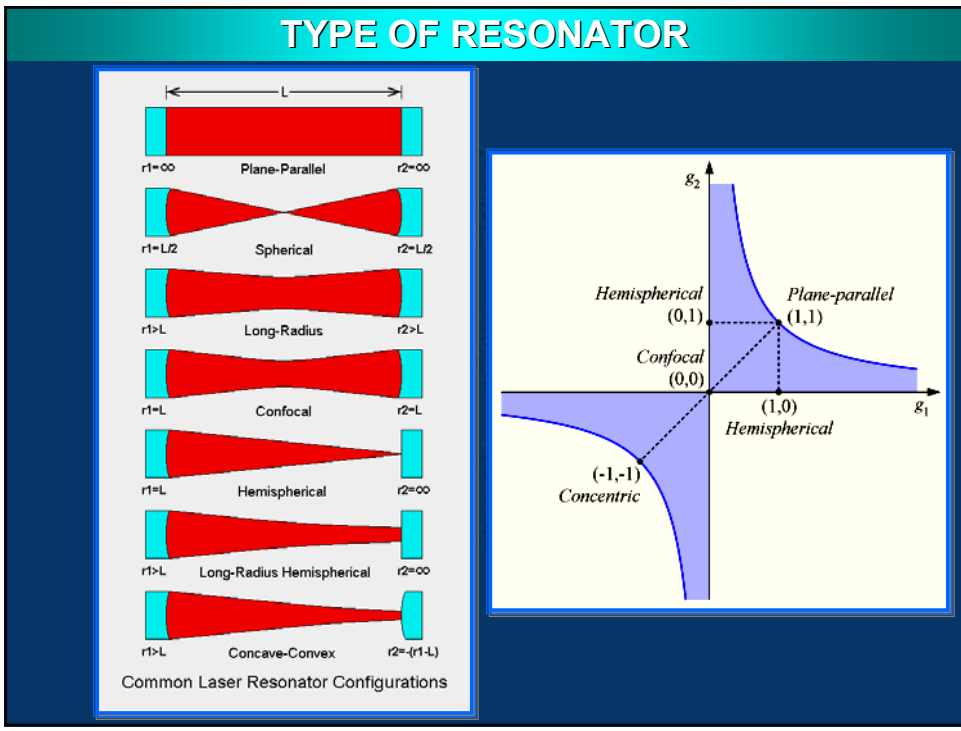


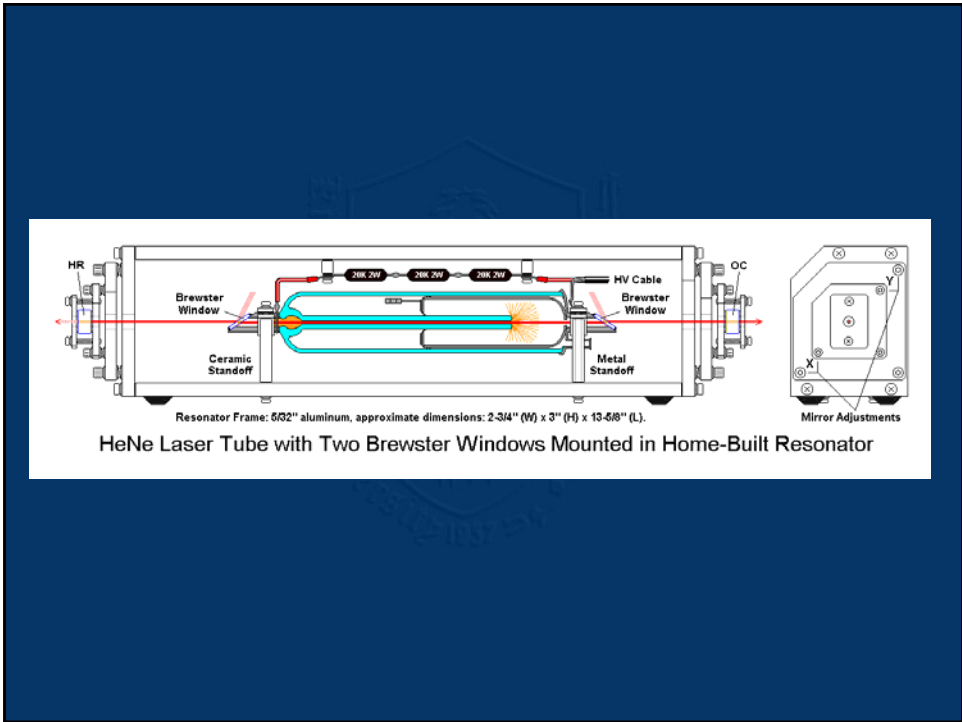
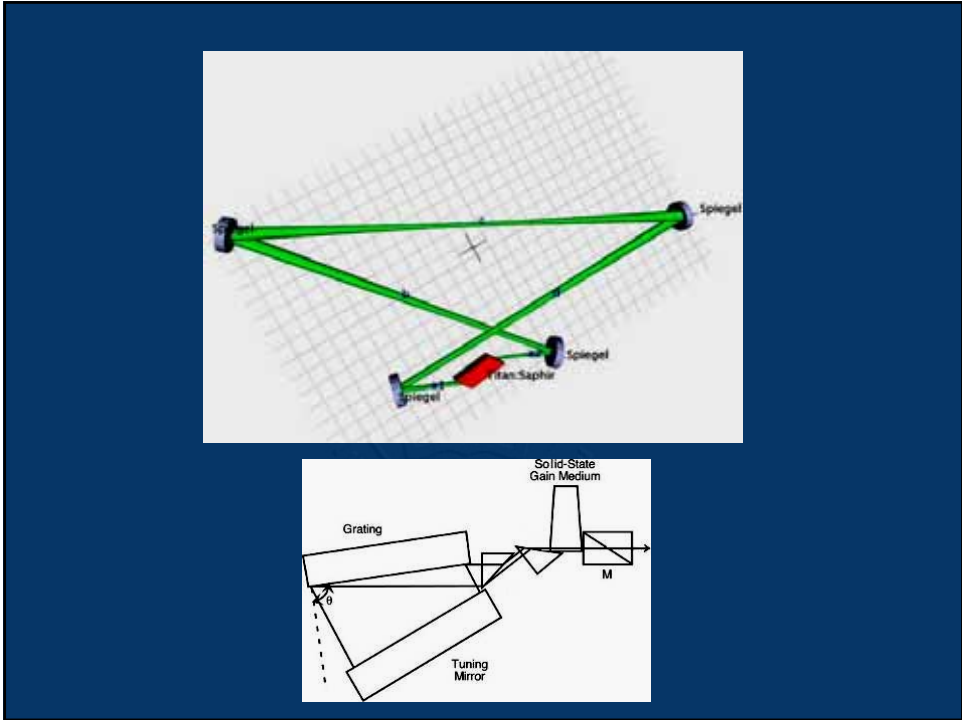
Gain
Gain envelope
Oscillation threshold
 ν
 $\Delta\nu$

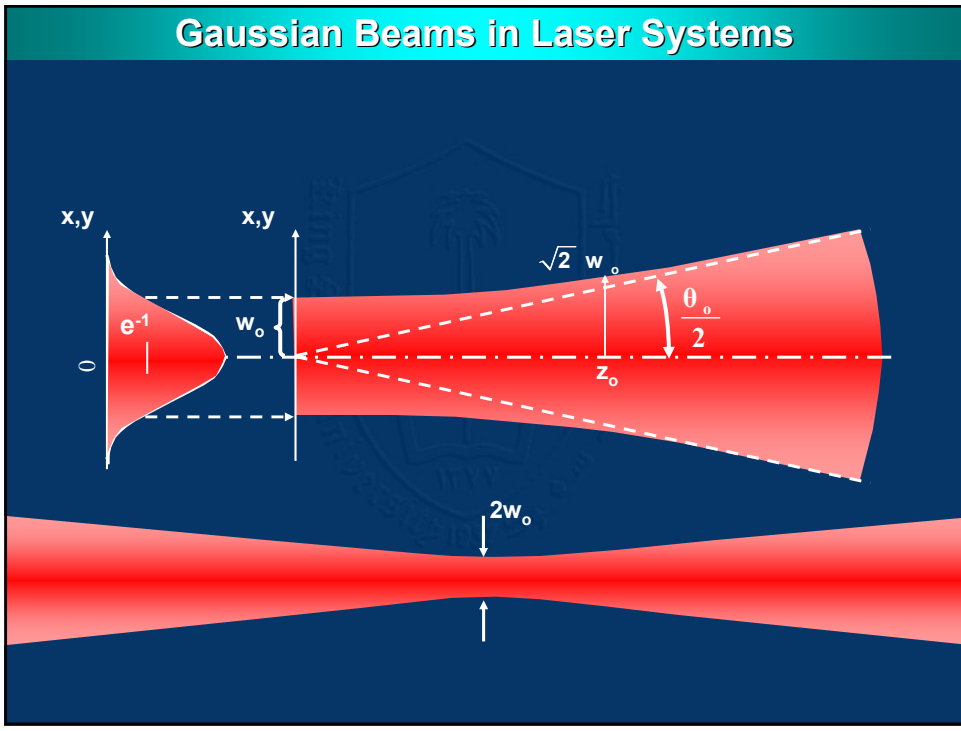
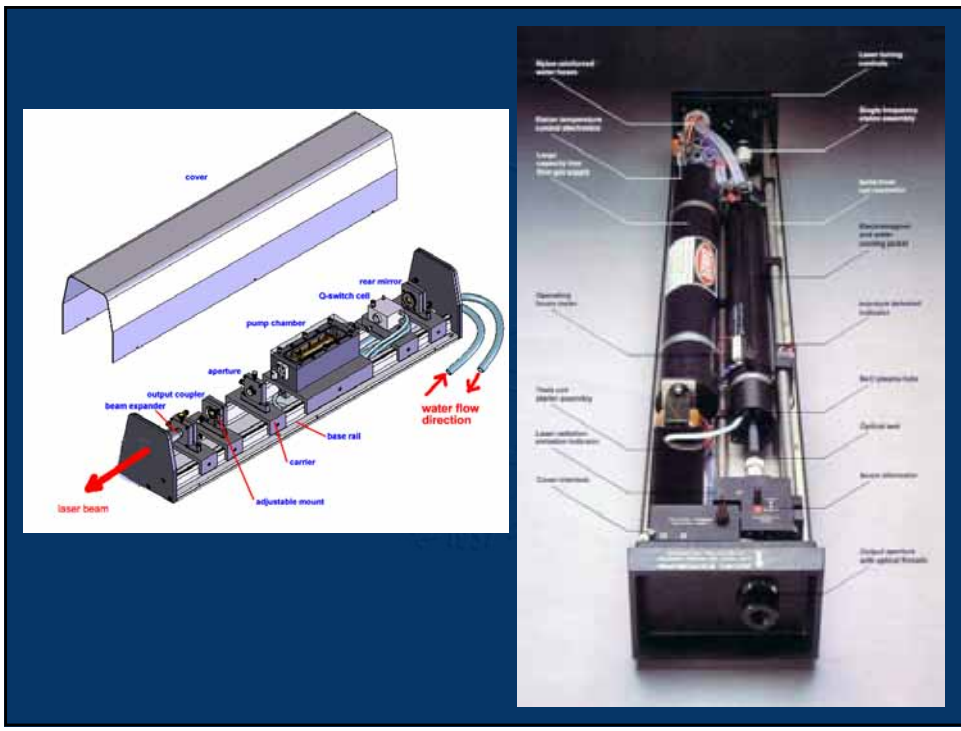
*Plot of laser gain as a function of frequency.
The natural gain bandwidth (black curve) is further restricted to discrete cavity modes (red), of which only four are above the laser oscillation threshold.*

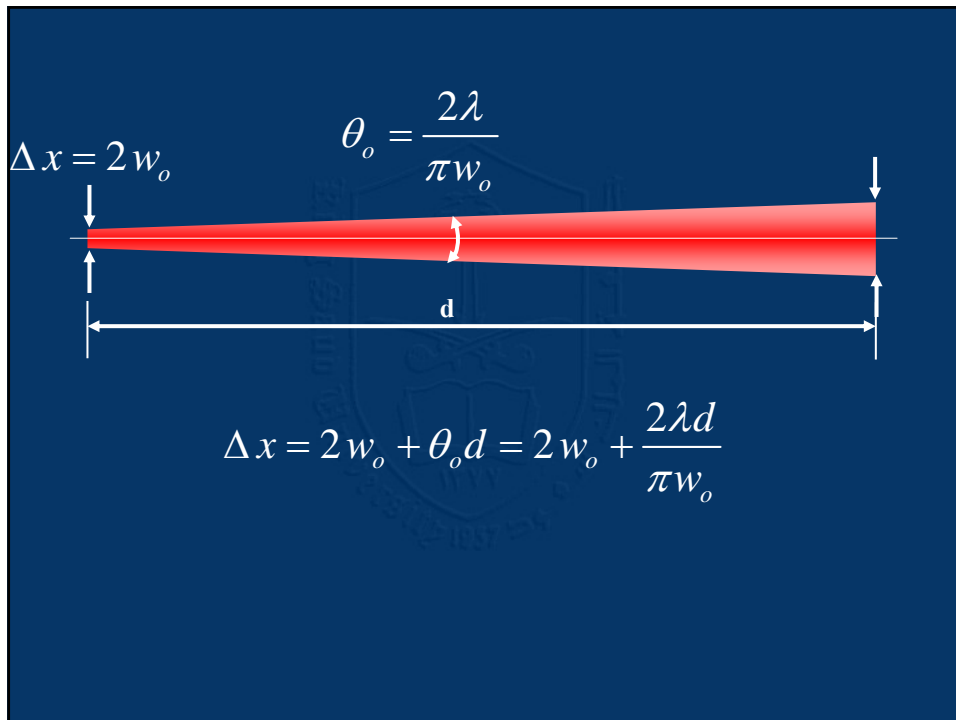
Electric field distribution of a few of the lowest order resonant modes of a cylindrical cavity. Up or down arrows (red or blue regions) indicate the phase of the electric field and arrow length indicates relative strength.

Effect of diffraction in selecting resonant modes.









Paraxial Wave Equation

$\nabla \cdot \mathbf{E} = 0$ Zero charge density

$\nabla \cdot \mathbf{E} = \nabla_t \cdot \mathbf{E}_t + \frac{\partial}{\partial z} E_z = 0$ Separate transverse and axial components

$E_z = E_0 \exp(-jkz)$ $k \sim \omega n/c = 2\pi n/\lambda_0$ Major z dependence

$\frac{\partial E_z}{\partial z} \sim -j \frac{2\pi n}{\lambda_0} E_z$ Approximate derivative of E wrt z

$\nabla_t \cdot \mathbf{E}_t \sim \frac{E_t}{D}$ Approximate transverse divergence given a finite beam diameter D

$|E_z| \approx \frac{\lambda_0}{2\pi n D} |E_t|$ z component is small compared to transverse component

Separate the **rapid variation** in z from the **slow spatial variation** in x,y and z

$$E(x, y, z) = E_o \psi(x, y, z) e^{-jkz}$$

$$\nabla^2 E + \frac{\omega^2}{c^2} n^2 E = 0 \quad \text{Substitute into the wave equation}$$

$$\nabla_t^2 E + \frac{\partial^2 E}{\partial z^2} + \frac{\omega^2}{c^2} n^2 E = 0$$

$$E = E_o \psi(x, y, z) \exp(-jkz)$$

$$k = \frac{\omega}{c} n$$

Gaussian beam equation

Derivatives needed for next step

$$\nabla_t^2 E = E_o (\nabla_t^2 \psi) \exp(-jkz)$$

$$\frac{\partial E}{\partial z} = E_o \left(-jk\psi + \frac{\partial \psi}{\partial z} \right) \exp(-jkz)$$

$$\frac{\partial^2 E}{\partial z^2} = E_o \left(-k^2\psi - j2k \frac{\partial \psi}{\partial z} + \frac{\partial^2 \psi}{\partial z^2} \right) \exp(-jkz)$$

$$\nabla_t^2 \psi - j2k \frac{\partial \psi}{\partial z} + \frac{\partial^2 \psi}{\partial z^2} = 0 \quad \text{Collect terms and cancel out the exponential factor}$$

$$\nabla_t^2 \psi - j2k \frac{\partial \psi}{\partial z} = 0 \quad (\text{paraxial wave equation})$$

Neglecting the second partial derivative

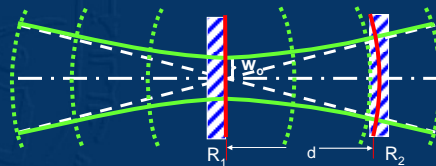
$$\nabla_{\perp}^2 \psi - j 2k \frac{\partial \psi}{\partial z} = 0$$

$$\psi_o(r) = \exp \left\{ -j \left[P(z) + \frac{k r^2}{2q(z)} \right] \right\}$$

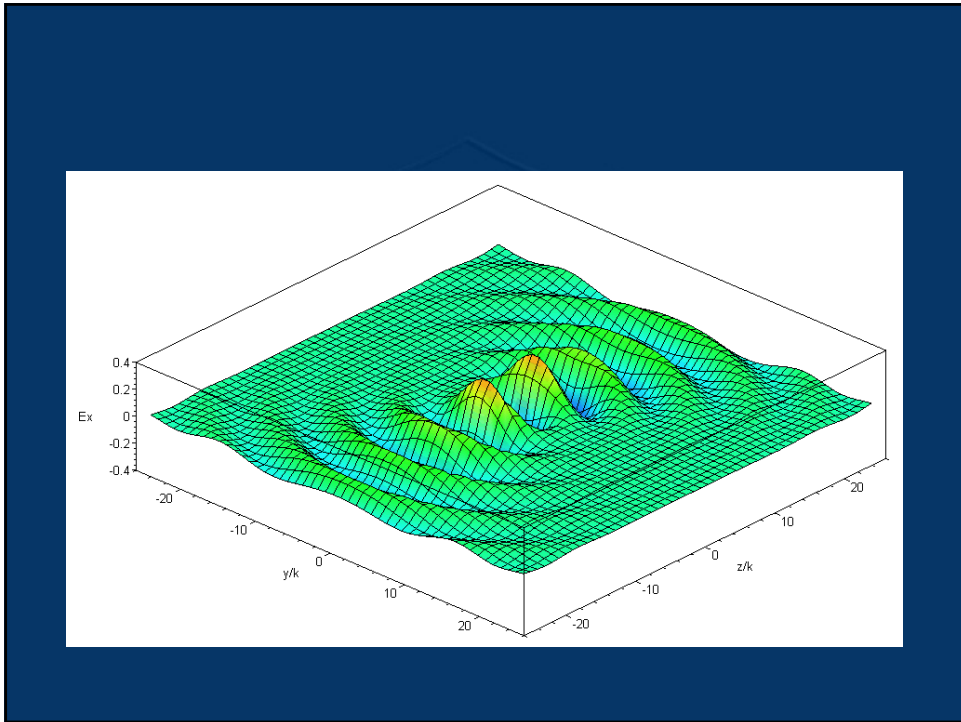
$$E(x, y, z) = \left\{ \frac{w_o}{w(z)} \exp \left[-\frac{r^2}{w^2(z)} \right] \right\} \\ \times \exp \left\{ -j \left[kz - \tan^{-1} \left(\frac{z}{z_o} \right) \right] \right\} \\ \times \exp \left[-j \frac{kr^2}{2R(z)} \right]$$

$q(z) = z + jz_o$ Complex beam parameter

$$\frac{1}{q(z)} = \frac{1}{z + jz_o} = \frac{z}{z^2 + z_o^2} - j \frac{z_o}{z^2 + z_o^2} = \frac{1}{R(z)} - j \frac{\lambda_o}{\pi n w^2}$$



$$w^2(z) = w_o^2 \left[1 + \left(\frac{z}{z_o} \right)^2 \right] \\ R(z) = \frac{z^2 + z_o^2}{z} = z \left[1 + \left(\frac{z_o}{z} \right)^2 \right] \\ z_o = \frac{\pi w_o^2}{\lambda_o}$$



Ray Tracing in an Optical System

$$r_2 = 1 \times r_1 + d \times r_1'$$

$$r_2' = 0 \times r_1 + 1 \times r_1'$$

$$\begin{bmatrix} r_2 \\ r_2' \end{bmatrix} = \begin{bmatrix} 1 & d \\ 0 & 1 \end{bmatrix} \begin{bmatrix} r_1 \\ r_1' \end{bmatrix}$$

$$\begin{bmatrix} r_{out} \\ r_{out}' \end{bmatrix} = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} r_{in} \\ r_{in}' \end{bmatrix}$$

Thin Lens Ray Matrix

$r_1 = r_2 \quad A=1 \quad B=0$

Consider $\alpha \quad C=-1/f$ Consider $\beta \quad D=1$

$$T = \begin{bmatrix} A & B \\ C & D \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{bmatrix}$$

Combination of lens and free space

$$T = \begin{bmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{bmatrix} \begin{bmatrix} 1 & d \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 1 & d \\ -\frac{1}{f} & 1 - \frac{d}{f} \end{bmatrix}$$

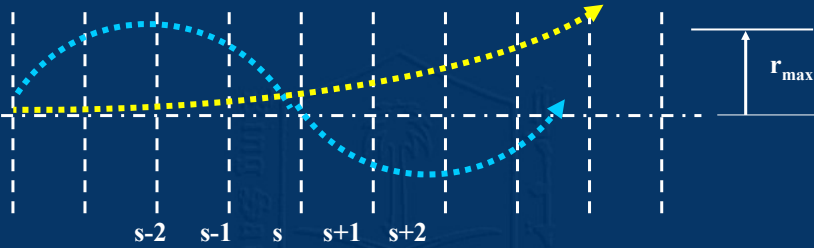
Optical Resonator

$$T = \begin{bmatrix} 1 & d \\ -\frac{1}{f_1} & 1 - \frac{d}{f_1} \end{bmatrix} \begin{bmatrix} 1 & d \\ -\frac{1}{f_2} & 1 - \frac{d}{f_2} \end{bmatrix}$$

$$T = \begin{bmatrix} 1 & d \\ -\frac{1}{f_1} & 1 - \frac{d}{f_1} \end{bmatrix} \begin{bmatrix} 1 & d \\ -\frac{1}{f_2} & 1 - \frac{d}{f_2} \end{bmatrix}$$

$$T = \begin{bmatrix} 1 - \frac{d}{f_2} & d + d \left(1 - \frac{d}{f_2} \right) \\ -\frac{1}{f_1} - \frac{1}{f_2} \left(1 - \frac{d}{f_1} \right) & \left(1 - \frac{d}{f_1} \right) \left(1 - \frac{d}{f_2} \right) - \frac{d}{f_1} \end{bmatrix}$$

Stability Conditions



$$r_{s+1} = Ar_s + Br'_s \quad \text{or} \quad r'_s = \frac{1}{B}(r_{s+1} - Ar_s)$$

$$r'_{s+1} = \frac{1}{B}(r_{s+2} - Ar_{s+1}) = Cr_s + Dr'_s$$

$$\frac{1}{B}(r_{s+2} - Ar_{s+1}) = Cr_s + \frac{D}{B}(r_{s+1} - Ar_s)$$

$$r_{s+2} - 2\left(\frac{A+D}{2}\right)r_{s+1} + r_s = 0 \quad (\text{use } AD - BC = 1)$$

assume $r_s = r_o (e^{j\theta})^s = r_o e^{js\theta}$

$$r_o \left[e^{j2\theta} - 2\left(\frac{A+D}{2}\right)e^{j\theta} + 1 \right] = 0$$

$$e^{j\theta} = \frac{A+D}{2} \pm j \left[1 + \left(\frac{A+D}{2}\right)^2 \right]^{\frac{1}{2}}$$

Form of Solution

$$r_s = r_o e^{js\theta} + r_o^* e^{-js\theta}$$

$$r_s = r_{\max} \sin(s\theta + \alpha)$$

Stability Diagram

$$-1 \leq \left(\cos \theta = \frac{A+D}{2} \right) \leq 1$$

$$0 \leq \frac{A+D+2}{4} \leq 1 \text{ stable regime}$$

$$\frac{A+D+2}{4} = \frac{1}{4} \left[1 - \frac{d}{f_2} - \frac{d}{f_1} + \left(1 - \frac{d}{f_2} \right) \left(1 - \frac{d}{f_1} \right) + 2 \right]$$

$$= 1 - \frac{d}{2f_1} - \frac{d}{2f_2} + \frac{d}{4f_1 f_2} = \left(1 - \frac{d}{2f_1} \right) \left(1 - \frac{d}{2f_2} \right)$$

since $2f_1 = R_1$ and $2f_2 = R_2$

$$0 \leq g_1 g_2 \leq 1 \text{ where } g_{1,2} = 1 - \frac{d}{R_{1,2}}$$

